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# A COMPARISON OF KINEMATIC FLOOD ROUTING METHODS

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KINEMATIC FLOOD ROUTING METHODS

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### A Thesis

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APSTRACT:

To provide a logical framework for the comparison of various methods of kinematic flood routing a general method of kinematic flood routing is developed. After presenting the general framework, the properties of the numerical model are investigated by:

- 1. Algebraic examination of the finite difference scheme.
- 2. Numerical experiments using a high speed digital computer.
- Comparison of the kinematic flood routing results with results of simulations using the complete one dimensional dynamic representation.

Particular facets of the numerical kinematic model that were studied included:

- 1. The stability of the numerical schematizations.
- 2. The degree of approximation with the finite difference system.

ii

 The applicability of kinematic methods to unsteady flow systems.

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 Methods of extending the kinematic solutions to predict attenuation as well as translation of the flood wave through the channel systems.

The results indicate that kinematic flood routing methods differ primarily in the point about which the finite difference equation is formulated, hereafter termed the nucleus, and that the general framework is capable of emulating such methods as the Muskinghum Method, other non-linear kinematic methods and reservoir routing. By varying the location of the nucleus the stability and degree of approximation is significantly altered. This results in the outflow hydrograph being sensitive to the location of the nucleus and the size of the finite difference steps.

To facilitate further research and application of the methods outlined in the thesis, a computer program was developed to enable kinematic flood routing to be performed in a natural channel with arbitrary geometry. Furthermore, the data is compatible with a program that is capable of performing a flood routing analysis using a numerical solution of the complete Saint-Venant equations. Documentation of the computer program for kinematic analysis is included with this thesis.

iii

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iv

# TABLE OF CONTENTS

Ţ

Title Page	i
Descriptive Note	ii
Acknowledgements	iv
Table of Contents	v
List of Figures	xi
List of Tables	xv

# CHAPTER 1 - INTRODUCTION 1

CHAPTER 2	2 - THE DYNAMIC SOLUTION	7
2.1 7	The Equations	8
2.2 1	Discussion of Numerical Methods	10
:	2.2.1 Characteristics Method	13
:	2.2.2 Direct Method	14
2.3 N	Numerical Method Used in This Study	15
2.4 H	Finite Difference Formulation of the Unsteady	
I	Flow Equations	19
2.5 1	Discussion of the Finite Difference Solution	24
2.6 (	Conclusions	34

CHAPTER	3 - THE KINEMATIC SOLUTION	38
3.1	Theory	39
3, 2	Application of Kinematic Wave Theory	44
3, 3	A Conceptual Model	49
	3, 3, 1 The Muskinghum Method as a Special	
	Case of Kinematic Routing	52
	3, 3, 2 Reservoir Routing as a Special Case	
	of Kinemetic Routing	55
	3, 3, 3 Further Comparisons	57
3.4	General Significance of the $ {oldsymbol a}$ and ${oldsymbol eta}$	
	Parameters	61
	3.4.1 Spatial Derivatives	61
	3.4.2 Time Derivatives	62
3, 5	Testing of the Kinematic Flood Routing	
	Techniques	63
	3.5.1 Stability	63
	3.5.2 Degree of Approximation	65
	3.5.3 Discretization Error	68
	3.5.4 Convergence	78
3.6	Verification of the General Kinematic Flood	
	Routing Method	81
3.7	Conclusions	83

------

ļ

•

•

CHAPTER 4 - COMPARISON OF COMPLETE AND	
KINEMATIC SOLUTIONS	88
4.1 Order of Magnitude Analysis	89
4.2 The Sustems Studied	97
4, 3 Comparison of Results for System 1	103
4.4 Comparison of Results for System 2	105
4.5 Further Comparisons	108
4.6 Conclusions	115
CHAPTER 5 - ATTENUATION AND KINEMATIC	•
METHODS	119
5.1 The "Molecule Effect" - Modelling Atten-	
uation by Moving the Nucleus	120
5.2 "Cascade Effect"	136
5.3 "Storage Effect"	140
5.4 Analysis of Numerical Experiments	141
5.4.1 Cascade Effect	141
5.4.2 Molecule Effect	143
5.4.3 Storage Effect	143
5.4.4 Extension to Other Systems	147
5.5 Guidelines for Calibration of Kinematic	
Routing Models	147

,

.

vii

1

•

\_\_\_\_\_

5.6	Discussion of Hydrographs Resulting from	
	Attenuated Kinematic Solutions	155
5.7 Conc	lusions	160
CHAPTER	R 6 - MODELLING ATTENUATION WITH	
	AN IMAGINARY RESERVOIR	163
6.1	The Physical System	164
6.2	The Position of the Reservoir	167
6.3	Measurement and Definition of the Response	
	Characteristics	173
6.4	System Parameters	174
6.5	Discussion of the Results	178
	6.5.1 Relation of Peak Flow Ratio, P, to	
	System Parameters	178
	6.5.2 Relation of Time of Centroid, Tc,	
	to Systems Parameters	180
	6.5.3 Relation of Skew Factor, SF to	
	System Parameters	187
6.6	Further Comparisons	184
6.7	Conclusions	188

**`** 

viii

.

,

CHAPTER 7 - A KINEMATIC FLOOD ROUTING	
MODEL	194
7.1 Objectives	195
7.? Development of Program Procedure	197
7.3 The Program	200
7.4 Use of the Program	201
7.4.1 Definition of Inflow Hydrograph	2.05
7.4.2 Performing the Routing	205
7.5 Example Applications	207
7.5.1 Application One	207
7.5.2 Application Two	213
CHAPTER 8 - CONCLUSIONS	<b>22</b> 6
BIBLIOGRAPHY	234
APPENDIX A - Stability Analysis of the General	
Kinematic Method	238
APPENDIX E - Degree of Approximation for the	
Dynamic Analysis	251

.

,

1

.

 $\mathbf{i}\mathbf{x}$ 

.

APPENDIX C - Degree of Approximation For the	
General Kinematic Routing Method	254
APPENDIX D - Documentation of Computer Routines	<b>2</b> 64
APPENDIX E - Derivation of the Lax - Wendroff Method	303
APPENDIX F - Kinematic Flood Routing Method of	
Characteristics	308
APPENDIX G - Listings of Computer Input Files,	
Routines and Output	315
APPENDIX H - Summary of Notation	369
APPENDIX I - Example of Calibration	374

· .

- ,

.

,

# LIST OF FIGURES

## FIGURE NUMBER

.

l

2.1	Space - Time Diagram	12
2,2	Finite Difference Molecule for the	
	Momentum Equation	17
2,3	Finite Difference Molecule for the	
	Continuity Equation	18
2.4	Pictorial View of System One	27
2.5	Typical Results of Finite Difference	
	Analysis	28
3.1	Relation Between Wave Velocity and	
	Stage for a Triangular Cross Section	42
3,2	General Finite Difference Scheme	46
3,3	Representation of an Elementary	
	Reach	51
3,4	Finite Difference Molecule for	
	Lax - Wendroff Method	60
3,5	Comparison of Hydrograph Shapes	76
3,6	Comparison of Hydrograph Shapes	77
3,7	Comparison of Hydrograph Shapes	79
3, 8	Results Using Characteristics Method	84
4.1	Inflow Hydrographs	100

·

I.

xi

## FIGURE NUMBER

1

4.2	Outflow Hydrographs 40,000'	
	Downstream .	104
4.3	Stage - Discharge Curves	106
4.4	Outflow Hydrographs	107
4.5	Outflow Hydrographs	109
4.6	Stage - Discharge Curves	110
5,1	Attentuation Vs. Number of Reaches	139
5.2	Attenuation per Elementary Reach	
	Vs. Stability Number	144
5.3	Attenuation per Elementary Reach	
	Due to Storage Effect	<b>1</b> 46
5.4	Attenuation Due to Molecule Effect	<b>1</b> 49
5,5	Attenuation Due to Molecule Effect	150
5.6	Outflow Hydrographs	156
5.7	Outflow Hydrographs	157
5,8	Outflow Hydrographs	158
5.9	Outflow Hydrographs	159
6.1	Schematic Channel - Reservoir	
	System	166

----

-----

•

## FIGURE NUMBER

ł

,

6.2	Typical Hydrographs Channel -	
	Reservoir	171
6.3	Definition of Response Parameters	175
6.4	Definition of System Parameters	177
6.5.	Peak Outflow Vs. Chord Slope	179
6.6	Centroidal Lag Vs, Chord Slope	181
6.7	Skew Factor Vs. Chord Slope	183
6.8	Outflow Hydrographs	185
6.9	Peak Outflow Vs. Tp	187
6.10	Typical Reservoir Ratings	189
7.1	Definition of the Channel Cross-	
	Section	199
7.2	Program Flow Chart	2.0 2
7,3	List of Commands	203
7.4	Invert Profile and Cross-Section	208
7.5	Hydrographs for Application One	209
7.6	Hydrographs for Application One	2 12
7.7	Invert Profile and Cross-Section	2 14
7.8	Hydrographs for Application Two	2 16
7.9	Hydrographs for Application Two	2 18
7.10	Hydrographs for Application Two	220

.

٠

١

.

## FIGURE NUMBER

۰

C.1	Plot of (2 <b>a</b> -1)	262
C. 2	Plot of $(1-2\beta)$	263
KINRUT <b>I</b> :	General Finite Difference Scheme	282
F. 1	Wave Velocity Vs. Flow Rate	314

.

.

~

ł

•

# LIST OF TAPLES

· · · · ·

\_

· ----

#### / TABLE NUMBER

.

2,1	Stability Tests of Explicit Method	30
2.2	Sensitivity Tests of the Explicit Method	
	and Comparison with Implicit Method	33
2,3	Comparison of Results for Flood	
	Routing Methods	35
3.1	Comparison of Finite Difference Schemes	58
3. 2	Results of Stability Analysis	66
3.3	Peak Values of the Outflow Hydrograph	70
3,4	Peak Values of the Outflow Hydrograph	71
3.5	Peak Values of the Outflow Hydrograph	72
3.6	Peak Values of the Outflow Hydrograph	74
4.1	Comparison of Systems	99
4.2	Order of Magnitude Analysis	102
4,3	Comparisons Two Hundred Miles	
	Downstream	112
5.1	Peak Values of the Outflow Hydrograph	122
5.2	Peak Values of the Outflow Hydrograph	123
5,3	Peak Values of the Outflow Hydrograph	124
5.4	Peak Values of the Outflow Hydrograph	125

# TABLE NUMPER

.

ĺ

5,5	Typical Results of Cascade Effect	
	Experiments	142
5.6	Typical Values Used to Define P Vs.	
	S <b>∆</b> X Curve	151
5.7	Typical Values Used to Define P Vs.	
	SAX Curve	152
6.1	Data From Reservoir-Channel	
	Simulation	168
6.2	Data From Channel-Reservoir	
	Simulation	<b>1</b> 69
A.1	Results of Stability Analysis	250
F,1	Flow Rate, Flow Ratio and Kinematic	
	Wave Velocity as a Function of Depth	311
F,2	Results of Kinematic Flood Routing	
	System 1	313

•

#### CHAPTER 1

#### INTRODUCTION

Flood routing is the process of calculating the deformation and position of a flood wave as it passes through a body of water with a free surface. Fluctuations in the level of this surface provide temporary changes in storage which in turn give rise to some reduction in the flood peak. In the recent past there has been a considerable amount of effort directed towards the development of efficient methods of performing these types of calculations. The main reason for the interest in this particular type of unsteady flow phenomena may be the multitude of ways that this na<sup>+</sup>ural and sometimes man made occurrence affects the life patterns of humanity.

The primary usage of flood routing techniques is frequently associated with catastrophic floods, the objective being to estimate the magnitude and/or depth of the flows that may be expected at particular locations. The utilization of flood routing calculations for flood warning purposes is however just one of a number of functions that these mathematical tools can perform. For example, the analysis of unsteady flow in river systems may be a vital part of the day to day operation of a hydro-electric scheme on a multi-purpose waterway. The use of these algorithms as an operational tool may also aid

in the efficient manipulation of flow control structures so that water is available for power generation, shipping and recreational purposes, without causing undue fluctuations in water levels or causing sudden surges in the channels.

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The use of flood routing techniques as a planning and design tool must not be overlooked. Flood routing techniques are beneficial in analyzing a river system which is being developed or modified and thus ensure that adverse side effects and environmental impact are minimized. Alternatively, the analysis of an existing watershed may be useful in the development of land zoning bylaws to prevent the construction of expensive structures in areas subject to flooding.

Currently there is increasing interest in the establishment of such flood plain maps in semi-urban areas. In such locations the incidence of highway culverts and bridges result in the formation of a chain of "reservoirs" along short, relatively steep watercourses and it is essential to have an economic and reliable tool for the analysis of such systems.

Just as there are a large number and diverse types of problems that are tackled by flood routing algorithms, there are numerous approaches to performing the actual computation. Two general classifications may be used to identify flood routing techniques. These are:

- Hydraulic techniques. The methods that fall into this classification are usually founded on the two laws that govern unsteady flow; conservation of mass and conservation of energy or momentum.
  - 2. Hydrologic techniques. These algorithms frequently do not employ the rigorous equations describing unsteady flow, but instead attempt to model the system by the use of equations that yield results similar to the observed phenomena.

The correct classification of a particular approach to flood routing may be very dependent on the system being simulated. For example, if a flood routing method is formulated from the partial differential equations which describe unsteady flow, but with several of the time variant terms ignored, the approach could be classified as hydraulic if the exclusion of the terms is justifiable. However, if a system were encountered where it was not realistic to neglect some terms, the algorithm would fall into the hydrologic classification. Kinematic flood routing methods are a particular example of techniques which may be classified as either hydraulic or hydrologic.

Kinematic flood routing is a generic term which identifies a broad class of numerical methods used to route flood waves through a channel or waterway. Flood waves which behaved in a kinematic fashion were observed on the Mississippi River by Seddon (1900). .3

The term kinematic was applied to the particular phenomena by Lighthill and Whitham (1955) for the reason that the equations are described in terms of velocities rather than forces.

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Prior to Lighthill and Whitham's presentation, storagerouting methods, which are special cases of kinematic routing, were employed by engineers and researchers. This was necessitated by the prohibitively large computational load imposed by a solution of the equations which describe the conservation of energy and mass. Even the simplified approaches were oftimes fraught with problems associated with excessive amounts of calculations. This led to the use of large time and distance steps in the numerical methods which in turn created stability problems in the numerical calculations.

With the advent of the high speed digital computer, it became feasible to tackle unsteady flow problems using a hydraulic approach. At the same time, attention was directed towards programming kinematic algorithms so that solutions could be obtained using the computer. This has resulted in a variety of techniques for approaching the problems. Because the hydraulic methods of analysis are usually founded on the momentum and continuity equations, these methods have a relatively unified basis on which they may be compared.

It seems that kinematic flood routing algorithms are not as clearly set in a logical framework for comparison, even though they

are all founded on the continuity equation. Thus, the primary purpose of this thesis is to develop a general framework that may be used to compare the various kinematic algorithms.

The usefulness of a numerical technique depends not only on how well it models the physical phenomena; but also on the familiarity of the user with the characteristics of the algorithm. This aids the user by providing information that will allow the user to identify the limitations of the model and particular properties of the model that may be used to an advantage. In addition, a thorough understanding of the computational tool being employed is necessary to differentiate between physical phenomena that are predicted and numerical phenomena, the result of instability, numerical error or poor convergence characteristics, which do not truthfully represent the actual physical system. Thus a second objective of this thesis, after developing the general framework, is to investigate the numerical characteristics of the algorithm.

To provide a verification of the theoretical studies of the general kinematic flood routing method, a series of numerical experiments are presented for purposes of demonstration and comparison. These computer simulations are also utilized to demonstrate the accuracy of kinematic simulations when representing dynamic physical systems. These numerical experiments provide a comparison to determine the

limits of kinematic algorithms and this investigation of the practical limitations of the kinematic methods is another of the thesis objectives. An explicit finite difference representation of the momentum and continuity equation is utilized to provide a data base for comparison purposes.

The application of a kinematic method to a particular problem may require the use of specific types of kinematic flood routing, for example, lag and route, Muskinghum flood routing, or reservoir routing. After studying the general method, several chapters are devoted to studying non-linear cases of these problems and a comparison of the results with dynamic simulations. The primary objective is to determine not only the useful ness of the various methods; but also to identify the effect of varying the step size in finite difference analyses and the effects of non-linearity.

The final objective of the study is to provide an efficient and versatile computer program that will enable a user to perform the various kinematic flood routing computations. This has been accomplished and a chapter is devoted to describing the development and use of the computer program. Designed to be used in a timeshared mode, the program is compatible with the Civil Engineering Program Library to allow easy access to the completed program.

#### CHAPTER 2

#### THE DYNAMIC SOLUTION

To provide a precise data base against which alternative flood routing methods may be compared it is necessary to have a means of generating the time history of flows resulting from inflow of a specified hydrograph to a known channel system. The use of natural channel systems to provide a data base can be immediately ruled out due to the complexity of both system and input, the inability of controlling flow parameters and the expense of monitoring flows at downstream sections A laboratory facility of sufficient scope and flexibility was not available and it was therefore decided to substitute a numerical model capable of generating the required data base with reasonable accuracy for systems of simple geometry.

The mathematical analysis of unsteady flow phenomena is founded upon the partial differential equations based on the laws of conservation of mass and linear momentum. Methods of solving these equations to yield values of flow depth and quantity at any desired point in space and time are numerical in nature (as opposed to analytical) and consist of the solution of finite difference formulations

of the partial differential equations. This chapter is concerned with the discussion of the method adopted in this study to obtain such a solution.

A brief description of the literature and alternate schemes for analyzing unsteady flow is provided prior to presenting the method used in this study. Following the description of the finite difference formulation, the sensitivity tests used to verify the numerical model are outlined. The conclusions summarize the findings that the algorithm was capable of providing the data necessary for comparing alternative methods of flood routing.

#### 2.1 THE EQUATIONS

The equations describing unsteady flow may be written as follows:

The Momentum Equation:

$$\frac{\delta h}{\delta x} + \frac{\delta}{\delta x} \left( \frac{Q^2}{2gA^2} \right) + \frac{1}{gA} \frac{\delta Q}{\delta t} + Sf = 0$$
 (2.1)

The Continuity Equation:

$$\frac{\delta Q}{\delta x} + \frac{T_w \delta h}{\delta t} = \bar{q}$$
 (2.2)

Where: x = distance

t = time

- h = water surface elevation
- Q = flow rate
- A = area of cross section
- g = acceleration of gravity
- $T_{W}$  = surface width
- $\vec{q}$  = rate of lateral inflow
- Sf = slope of the friction line

The derivation of these equations, often called De Saint Venant's equations, may be found in numerous books. Several of these are Stoker (1957), Chow (1959), and Henderson (1965). The first two terms in equation 2.1 describe effects caused by nonuniform flow. Results of unsteady conditions are reflected in the third term which describes temporal acceleration. The effect of friction is modelled by the slope of the friction line. Terms in the continuity equation (equation 2.2) describe the change in flow rate along the channel, the rate of change of storage, and the amount of lateral inflow.

Because no analytic solution has been found for the partial differential equations which describe unsteady flow conditions, finite difference techniques must be used. Previous to the advent of the high speed digital computer it was not practically feasible to obtain solutions of these equations. In special cases, the method of characteristics was employed. However, the amount of calculations required for this relatively rapid method of solution limited its use. With the increased availability of high speed digital computers, it has become possible to solve the partial differential equations using numerical techniques. Still, the use of these methods may be rather difficult and expensive due to problems of stability with some formulations and the expense involved in programming and operating the computer.

Numerical solutions of De Saint Venant's equations do have several advantages. These equations provide an accurate description of the one dimensional flow system and enable the user to obtain very detailed information about the wave shape and its position during the time in which the wave is being propogated along the channel. It is beyond the scope of this chapter to provide an exhaustive comparison of all the various techniques used to solve the equations which describe unsteady flow. However, a short section giving a general description of several techniques follows.

#### 2.2 DISCUSSION OF NUMERICAL METHODS

There exist two general methods of obtaining solutions to the partial differential equations. These are direct methods and characteristics methods. The partial differential equations describing unsteady flow may be rearranged in the form of four ordinary

differential equations. Two of these equations define characteristic lines, paths of energy transfer, while the two other equations define energy change along the characteristic lines. Applying finite difference techniques to these equations is termed the characteristics method. Alternately, finite difference formulations of the partial differential equations may be used to provide solutions. This is known as the direct method.

The finite difference techniques used to solve either characteristics or direct methods may in turn be classified as explicit or implicit. An explicit method provides for a specific solution for an unknown quantity while an implicit technique requires the solution of several similtaneous equations to provide the values of a number of quantities. In both methods, the object is to obtain values of flow rate (or velocity) and water surface elevation (or water depth) at discrete points on a space-time diagram. Figure 2.1 shows a spacetime diagram with a staggered rectangular grid. The solution progresses from the known initial conditions through successive increments in time. Thus, conditions along the channel at time "t" are used in conjunction with the boundary conditions to find the solution at time "t +  $\Delta$ t".

The type of boundary conditions encountered depends on the physical system being simulated. If flow is subcritical there is a

# FIGURE 2.1

# SPACE-TIME DIAGRAM

FLOW RATE TO BE CALCULATED

MOMENTUM EQUATION APPLIED HERE

O SURFACE ELEVATION TO BE CALCULATED

CONTINUITY EQUATION APPLIED HERE



boundary condition at the upstream and downstream limits. Typically, the upstream boundary would be an inflow hydrograph while the downstream condition would be water surface elevation as a function of time and/or flow rate. When supercritical flow is encountered the boundary conditions are found only at the upstream limit.

The type of grid used on the space-time diagram is related to the type of solutions used. A characteristics solution, which results from applying finite difference methods to characteristic equations may be used with an irregular grid defined by intersections of characteristic lines or may be applied to a rectangular grid. The direct methods are usually used with a regular rectangular grid; the staggered rectangular grid being used primarily with explicit formulations of the direct method.

The attributes and drawbacks of various methods are briefly outlined in the following paragraphs.

#### 2.2.1 Characteristics Method.

This method is believed to solve a given space-time diagram in the least time. More accuracy is claimed, especially where the flow varies quite rapidly as the solution progresses along the paths of energy transfer (characteristic lines). The answers are not provided at fixed points in time or space, which is a disadvantage when informa-

tion at a particular time or location is required. This difficulty can be overcome by using a method of characteristics which solves for fixed points on a space-time diagram. Several other factors that favour the use of characteristics methods are:

- 1. The solution is more stable when flow conditions are supercritical.
- The case of a flood wave propagating down an initially dry stream bed is more correctly modelled.
- Characteristics methods are the most accurate methods of modelling rapidly varied flow as the characteristics lines are closer in regions of rapid variation.

References are: Amein (1966), Henderson (1965), Woolhiser and Liggett (1967), and Yevjevich and Barnes (1970).

#### 2.2.2 Direct Method.

This method is widely used due to the relative ease of algebraically expressing the various equations and the subsequent reduction of programming difficulties. Answers are provided at fixed points in time and space which is convenient for interpretation of results. The disadvantages of direct finite difference techniques are that explicit versions are subject to stability problems especially when there are rapid variations of flow or if supercritical conditions are encountered. Stability problems have been overcome by using implicit algorithms.

References are: Woolhiser and Liggett (1967), Smith (1968), Amein (1968), Amein and Fang (1970), Yevjevich and Barnes (1970), and Walden (1973).

#### 2.3 NUMERICAL METHOD USED IN THIS STUDY

In choosing a numerical method to use as a base for comparison with approximate methods, preference was given to an algorithm which would provide the necessary accuracy with a minimal amount of computer programming. As initial tests were going to be made using rectangular channels, stability would not be as difficult a problem as would be encountered with a natural non-prismatic channel. Thus, an explicit method which used a staggered mesh on the timespace diagram was employed. This scheme had been successfully used in a similar situation where the channels were very nearly prismatic. James and Horne (1969), Smith (1968). In addition experience had been obtained with this method in conjunction with classroom studies. Thus, a small computer program was available which could be easily adapted to the present study.

The staggered mesh used by this method is shown in figure 2.1. An inflow hydrograph was used as the upstream boundary condition

while the downstream condition for subcritical flow was assumed to be uniform flow. Initial conditions consist of a horizontal row of known flow rates at  $t = -\Delta T/2$  and a horizontal row of water surface elevations specified at t = 0. The row of flow rates is displaced  $\Delta X/2$  upstream from the known water surface elevations.

For time  $\Delta T/2$ , the momentum equation is applied to a point under the first unknown flow rate downstream of the upstream limit. The unknown value is calculated using the initial conditions and the upstream boundary. This calculation is repeated as the process moves in the downstream direction. The previously unknown value is treated as the upstream boundary condition in calculating the next unknown flow rate. Figure 2.2 shows the way in which the dynamic equation is applied. When the downstream limit is reached, time is incremented by  $\Delta T/2$ , the downstream boundary value is obtained and a series of calculations is begun in the upstream direction. This time, the continuity equation is applied to a point below the first unknown water surface elevation upstream of the downstream limit. After the unknown elevation is determined, the calculation is repeated at the next upstream location. Again the previously unknown value is used as the downstream boundary condition in predicting the next unknown. Figure 2.3 portrays the application of the equation of continuity. When the upstream limit is reached, time is incremented by  $\Delta T/2$ 

# FIGURE 2.2

# FINITE DIFFERENCE MOLECULE FOR THE MOMENTUM EQUATION



# FIGURE 2.3

# FINITE DIFFERENCE MOLECULE FOR THE CONTINUITY EQUATION



a new value for the upstream boundary is obtained, and the cycle begins again with the application of the momentum equation in successive steps moving downstream.

By repeating the previously described cycle, a time history of flow conditions along the channel may be obtained. The computation is stopped when a defined time is reached or when a nearly steady state is reached after a flood wave has passed through the channel. The cost of computing is related to the number of iterations required to fill the time space diagram and the amount of calculations in each cycle. When only the outflow hydrograph is desired, this method may seem to be quite wasteful due to the amount of unnecessary data which must be generated. However, it will be shown later that some methods of finite difference solution of unsteady flow can be relatively inexpensive.

# 2.4 FINITE DIFFERENCE FORMULATION OF THE UNSTEADY FLOW EQUATIONS

The first step in the solution of a problem requiring the solving of partial differential equations by a finite difference method is to express the partial differentials as finite differences. This was done in the following manner. The momentum equation is:

$$\frac{\delta h}{\delta x} + \frac{\delta}{\delta x} \left( \frac{Q^2}{2gA^2} \right) + \frac{1}{gA} \frac{\delta Q}{\delta t} + Sf = 0$$
 (2.3)
Where:

: |

$$\frac{\delta}{\delta x} \left( \frac{Q^2}{2 g A^2} \right) = -\frac{Q^2}{g A^3} \frac{\delta A}{\delta x} + \frac{Q}{g A^2} \frac{\delta Q}{\delta x}$$
(2.4)

and

$$\frac{\delta A}{\delta x}\Big|_{t=\text{constant}}^{2} = \frac{T_{w} \delta h}{\delta x} + \frac{\delta A}{\delta x}\Big|_{h=\text{constant}}$$
(2.5)

Substituting equation 2.5 into equation 2.4

 $\frac{\delta}{\delta x} \left( \frac{Q^2}{2gA^2} \right) = \frac{Q^2 T_w \delta h}{gA^3 \delta x} - \frac{Q^2}{gA^3 \delta x} \left|_{h=\text{constant}}^{+} \frac{Q}{gA^2 \delta x} \right|_{h=\text{constant}}^{+} (2.6)$ 

From the continuity equation, assuming lateral inflow is equal to zero.

$$\frac{\delta Q}{\delta x} + T_w \frac{\delta h}{\delta t} = 0 \qquad (2.7)$$

$$\frac{\delta Q}{\delta x} = -T_w \frac{\delta h}{\delta x}$$
(2.8)

Thus

$$\frac{Q}{gA^{2}\delta x} = -\frac{QT_{w}}{gA^{2}}\frac{\delta h}{\delta t}$$
(2.9)

The dynamic equation may now be rewritten in the following manner

$$\frac{\delta h}{\delta x} - \frac{Q^2 T_w \delta h}{g A^3 \delta x} - \frac{Q^2 \delta A}{g A^3 \delta x} \Big|_{h=\text{constant}} \frac{Q T_w \delta h}{g A^2 \delta t} + \frac{1}{g A} \frac{\delta Q}{\delta t} + \text{Sf} = 0 \quad (2.10)$$

Figure 2.2 shows an enlarged portion of the space-time diagram. The momentum equation is applied at point "A". Q(K, L) is unknown and all the data is known for points on rows below Q(K, L). In addition, information is available for points on the L row to the left of Q(K, L) thus the following approximations can be made

$$\frac{\delta h}{\delta x} = \frac{H(I,J) - H(I-I,J)}{\Delta X}$$
(2.11)

$$\frac{\delta h}{\delta t} = \frac{H(I-I,J) - H(I-I,J-I) + H(I,J) - H(I,J-I)}{2\Delta T}$$
(2.12)  
$$\frac{\delta Q}{\delta t} = \frac{Q(K,L) - Q(K,L-I)}{\Delta T}$$
(2.13)

As this portion of the study is limited to rectangular channels with uniform slopes, the following approximations are appropriate.

$$A = T_{x} (H(I,J) - B(I) + H(I-I,J) - B(I-I)) / 2$$
(2.14)

$$\frac{\delta A}{\delta x} = T_{w} \frac{B(I) - B(I-I)}{\Delta x}$$
(2.15)

Where:

B = invert elevation at the section.

Describing the slope of the friction line by Mannings equation results in the following expression.

$$Sf = \frac{Q(K,L) \times |Q(K,L-1)| \times n^2 \times P^{4/3}}{2.21 \times A^{10/3}}$$
(2.16)

Where: n = Mannings roughness coefficient

$$p = T_w + (H(I,J) + H(I-I,J) - B(I) - B(I-I))$$

Using the absolute value of the known flow rate gives the energy slope term the same sign as the unknown flow rate. Thus, flow reversals which do not occur in a rapid fashion can be modelled with this scheme.

Rewriting the momentum equation in terms of finite differences yields the following:  $\begin{pmatrix} 1 - \frac{Q(K,L)Q(K,L-1)T_w}{gA^3} \end{pmatrix} \times \frac{H(I,J) - H(I-I,J)}{\Delta X}$   $- \frac{Q(K,L)Q(K,L-1)}{gA^3} \times \frac{B(I) - B(I-I)}{\Delta X}$   $- \frac{Q(K,L) + Q(K,L-1)}{2} \times \frac{T_w}{gA^2} \times \frac{H(I,J) - H(I,J-I) + H(I-I,J) - H(I-I,J-I)}{2\Delta T}$   $+ \frac{I}{gA} \times \frac{Q(K,L) - Q(K,L-I)}{\Delta T}$ 

$$+ \frac{Q(K,L) \times |Q(K,L-1)| \times n^2 \times P^{4/3}}{2.21 \times A^{10/3}} = 0$$
 (2.17)

Solving for the unknown flow rate produces the following expression.

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$$Q(K,L) = \left(\frac{H(I,J) - H(I-I,J)}{\Delta X} - Q(K,L-I)\frac{1}{gA\Delta T} + \frac{Q(K,L-I)T_{w}}{2gA^{2}} \times \frac{H(I-I,J) - H(I-I,J-I) + H(I,J) - H(I,J-I)}{2\Delta T}\right) / \left(\frac{Q(K,L-I)T_{w}}{gA^{3}} \times \frac{H(I,J) - H(I-I,J)}{\Delta X} + \frac{Q(K,L-I)}{gA^{3}} \times \frac{B(I) - B(I-I)}{\Delta X} + \frac{T_{w}}{2gA^{2}} \times \frac{H(I-I,J) - H(I-I,J-I) + H(I,J) - H(I,J-I)}{2\Delta T} - \frac{1}{gA\Delta T} - \frac{IQ(K,L-I)I \times n^{2} \times P^{4/3}}{2.2I \times A^{10/3}}\right)$$
(2.18)

The continuity equation is applied in a similar fashion. Assuming no lateral inflow, the continuity equation is:

$$\frac{\delta Q}{\delta x} + T_{w} \frac{\delta h}{\delta t} = 0$$
 (2.19)

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Applying this equation at point "B" on the space-time diagram of figure 2.3 allows the following approximations to be made:

$$\frac{\delta Q}{\delta x} = \frac{Q(K,L) - Q(K-I,L)}{\Delta X}$$
(2.20)  
$$\frac{\delta h}{\delta t} = \frac{H(I-I,J) - H(I-I,J-I)}{\Delta T}$$
(2.21)

Thus the continuity equation in finite difference form is:

$$\frac{Q(K,L) - Q(K-I,L)}{\Delta X} + T_{w} \frac{H(I-I,J) - H(I-I,J-I)}{\Delta I} = 0$$
 (2.22)

The unknown water level is given by the expression:

$$H(I-I,J) = H(I-I,J-I) - \frac{\Delta T(Q(K,L) - Q(K-I,L))}{T_{W} \Delta X}$$
(2.23)

#### 2.5 DISCUSSION OF THE FINITE DIFFERENCE SOLUTION

After a finite difference scheme has been developed, it is necessary to determine the limitations of the method. Numerical stability, convergence properties, the degree of approximation and the discretization errors are factors which influence the way in which the algorithm may be used. Numerical stability is a property of the numerical method which keeps errors from concealing the true solution. Convergence is a measure of the accuracy with which a finite difference equation will represent the partial differential equation as  $\Delta X$  and  $\Delta T$  approach zero. Another measure of the accuracy of the numerical solution is the degree of approximation. The discretization errors represent errors caused by replacing a derivative.(a tangent) with a finite difference (a chord).

The finite difference algorithm must first be numerically stable to be useful as a tool. An unstable formulation will allow small errors to grow unbounded which in turn will mask the true solution. Some algorithms are unstable, others are conditionally stable. Stability in explicit formulations of finite differences is largely dependent on the size of time step used in the calculation. This has been demonstrated by the Courant Condition which is:

$$(V+C)\frac{\Delta T}{\Delta X} \leq I$$
 (2.24)

Garrison et. al. (1969) reported that another condition must be satisfied for a particular explicit scheme to be stable. The additional constraint is:

$$(V + C)\frac{\Delta T}{\Delta X} = 1 - \frac{gn^2 I V I \Delta T}{2.21 R^{2/3}}$$
 (2.25)

For the present study, a pragmatic approach to determining the stability properties of the numerical scheme was used. A short description of the problem used in the tests is outlined in the following paragraphs.

To facilitate easy computation, a rectangular channel was used as a prototype. Figure 2.4 shows a picture of the channel similar to the one used in this study. Channel properties were as follows:

> Length = 50,000 ft. Width = 100 ft. Depth = 20 ft. Slope = varied for various executions

n = varied for various executions

The upstream boundary condition was a symmetrical triangular hydrograph. Figure 2.5 shows the characteristics of this hydrograph. Downstream control was assumed to be uniform flow depth for the flow rate of the previous time step.

Equal increments of  $\Delta X$  were used in this analysis. By entering the number of subreaches into the program,  $\Delta X$  was computed by dividing the total length by the number of subreaches. Further documentation of the computer program is provided in Appendix "D".





During the execution of the program,  $\Delta T$  was held fixed at a value determined by the following relationship:

$$(V+C)\frac{\Delta T}{\Delta X} \leq Z$$
 (2.26)

Where: v = full bank velocity

- c = full bank celerity
- Z = an arbitrarily chosen constant between 0 and 1

The constant, Z, is known as the Courant Number. It is defined as the time step used for the computation divided by the time step which satisfies the Courant Condition. The value of the Courant Number was reduced until no instabilities were detected. Results of several runs are shown in table 2.1. With Z = 0.67 the first test showed signs of instability on the falling limb of the hydrograph. Putting Z = 0.5 resulted in a hydrograph which showed no signs of instability. Increasing the slope to 0.001 and executing the program with Z=0.5 resulted in an unstable solution which terminated the job. The problem was successfully tackled with Z=0.25. This does not appear to agree with the stability condition as reported by Garrison. As the slope increases, the size of the Z value should decrease, this is in agreement with the condition of Garrison. However, their formula predicts that as the roughness coefficient n increases, the value of Z should also decrease. These tests indicate that as n increases

# TABLE 2.1

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# STABILITY TESTS OF EXPLICIT METHOD

Slope	n	Z	Comments	
0.0002	0.0149	0.67	slight signs of instability	
0.0002	0.0149	0.50	stable	
0.0010	0.0149	0,50	unstable	
0.0010	0.0149	0.25	stable	
0.0010	0,0322	0.50	stable	

the value of Z should also increase. It appears that:

$$Z = f\left(\frac{\sqrt{S}}{n}\right)$$
 (2.27)

The sensitivity of the solution to changes in  $\Delta X$  and  $\Delta T$  is related to the way that the finite difference equation approximates the differential equation and to the discretization errors.

The degree of approximation is obtained by substituting a Taylor's Series expansion into the finite difference solution and observing how well it represents the partial differential equation. Appendix "B" contains the calculations which show that the momentum equation is represented by the finite difference equation in the following manner:

$$\frac{\delta h}{\delta x} + \frac{\delta}{\delta x} \left( \frac{Q^2}{2gA^2} \right) + \frac{1}{gA} \frac{\delta Q}{\delta t} + Sf + O(\Delta X^2, \Delta T^2) = 0 \qquad (2.28)$$

Similarly the representation of the continuity equation is:

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} \div O(\Delta X^2, \Delta T^2) = \overline{q}$$
 (2.29)

Several executions of the computer program were performed to determine the effects of approximation errors and discretization errors as  $\Delta X$  and  $\Delta T$  vary.

Table 2.2 contains vital data from several executions which demonstrate the changes caused by varying  $\Delta X$  and  $\Delta T$ . Based on a comparison of peak values, it appears that the solution is sensitive only to the value of  $\Delta X$  used. The lack of sensitivity to changes in  $\Delta T$  can be attributed to the fact that stability criteria are a more stringent constraint than are convergence requirements. The stability criteria also makes  $\Delta T$  a function of  $\Delta X$ with the result that convergence appears dependent on  $\Delta X$ .

Figure 2.5 shows the inflow hydrograph and a typical outflow hydrograph obtained using the finite difference analyses.

A check was provided by comparing the results with an implicit method developed by Walden (1973) with those provided by the explicit method described in the report. These results are presented in table 2.2.

Further checks were provided by using an example proposed by Thomas (1934). Amein (1967) shows the results of routing a flood through a very wide channel using characteristics, a direct explicit method and a direct implicit technique. The channel had a slope of one foot per mile and a Mannings n approximately equal to 0.03. The inflow hydrograph was sinusoidal with an initial flow of 50 cfs/ft width, a peak flow of 200 cfs/ft. width, and a time base of 96 hours.

## TABLE 2.2

<b>Δ</b> X (Ft)	Z	∆T ∙(Secs.)	Time of Maximum Flow (hrs.)	Maximum Flow (cfs)	Time of Maximum Depth (hrs.)	Maximum Depth (ft.)
1000'	0.50	14.8	2.274	12537	2.620	15.858
2000	0.50	29.7	2.290	12577 _	2.636	15.841
5000	0.50	74.2	2.348	12619	2.636	15.850
2000	0.25	14.8	2.295	12551	2.636	15.843
2000	0.50	29.7	2.290	12577	2.636	15.841 .
2000	0.67	39.7	2.296	12602	2,628	<u>1</u> 5.845
*2000	10.2	600	2.330	12565	2.50	15.14

## SENSITIVITY TESTS OF THE EXPLICIT METHOD AND COMPARISON WITH IMPLICIT METHOD

Slope = 0.0002n = 0.0149

All data from a station located 40,000' from the start.

\*Implicit method

ί Ω These results are presented in table 2.3 along with results obtained from the direct explicit method utilized in this report. Based on these results it can be seen that the method used in this thesis are less sensitive to change in  $\Delta X$  and  $\Delta T$  than the other methods reported by Amein (1968). This may be due to the size of binary word used by the computer. A large binary word will reduce errors due to truncation in the numerical calculation. Other differences may be explained as follows: Using a staggered mesh, which displaces the inflow hydrograph  $\Delta X/2$  upstream, may explain why the peak is predicted at a slightly later time. Reducing the size of  $\Delta X$  causes the peak to occur at an earlier time. Varying the time step has only a small effect on the time of the peak.

#### 2.6 CONCLUSIONS

After studying the various numerical methods of solving the partial differential equations which describe unsteady flow, a direct explicit scheme was utilized to provide a precise data base for comparison of alternative flood routing methods. From tests to determine the stability and convergence properties of the finite difference formulation the following observations have been drawn:

1. To insure stability of the numerical solution, the size of the time increment must be reduced to a value which is smaller than

#### TABLE 2.3

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# COMPARISON OF RESULTS FOR FLOOD ROUTING METHODS

# DEPTH AND TIME OF ARRIVAL OF PEAK FLOW AT 200 M1 STATION

Method	<b>∆</b> X Mi	<b>∆</b> T Hrs	Depth Ft	Time of Arrival Hrs
Explicit (As per Amein)	5 5 10 10	0.05 0.10 0.15 0.20	28.6 28.6 26.9 26.0	76.0 76.0 76.0 76.0 76.0
Characteristics		0.2, 0.4 0.8, 1.2 2.0 2.5 3.3	29.0 28.6 26.9 26.0	76.0 76.0 76.0 76.0
Implicit	5 5 5 10	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	29.11 29.13 29.20 29.11	, 76.0 76.0 76.0 76.0 76.0
Explicit (As used in this report)	5 10 10	0.10 0.15 0.20	29.18 29.18 29.18	76.70 76.95 77.00

that required to satisfy the Courant Condition. The Courant Number was used to define the size of the time step. This number, which is the size of the time step used divided by the time step implied by the Courant Condition appears to be inversely proportional to  $\sqrt{S}/n$ . As n increases, the Courant Number may also increase and as  $\sqrt{S}$ increases, the Courant Number must decrease to insure stability.

While a smaller size of time step increased computation costs, it was still economically feasible to perform the required computations. Using a CDC 6400, only 36 seconds of central processor time were required to route a flood down the rectangular channel when it was divided into 25 sections each 2000 feet long and with Z = 0, 5.

2. Sensitivity tests and theoretical analysis show that the solution of the equation is sensitive to the size of  $\Delta X$  and  $\Delta T$ . However, the variation of the solution was deemed sufficiently accurate for the purposes of this study. The variation in the peak value was less than 1%. Also, the variations between alternate flood routing techniques are on a larger order of magnitude.

3. Comparisons with other finite difference schemes. which have been successfully employed, show that the method used in this study compares very favourably in representing unsteady flow in a system with simple geometry.

Thus it was concluded that the numerical model will provide the data base which is required in the study of alternative flood routing methods.

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#### CHAPTER 3

#### THE KINEMATIC SOLUTION

Kinematic flood routing is a generic term that identifies a class of methods for calculating the deformation of a flood wave as it passes through an open channel such as a river reach or man made conduit. These methods are based primarily upon the continuity equation and the assumption that there exists a single valued function relating flow rate to the physical properties of the channel. The popularity of these techniques can be attributed to the relative simplicity and low cost of obtaining solutions, especially before the widespread availability of the high speed digital computer lightened the computational load imposed by a complete solution of the momentum and continuity equations.

A literature review revealed numerous approaches to solving the continuity equation using either direct finite difference methods or characteristics techniques. This chapter attempts to provide a standard basis for comparison of the various algorithms by developing a general framework against which the different kinematic flood routing techniques may be compared. The first portion provides a review of the theoretical background of waves and shows the development of a general direct numerical method of kinematic flood routing. After discussing several special cases of kinematic flood routing and showing how they fit into the general method, a section is devoted to sensitivity tests of the numerical algorithm to determine limitations imposed by convergence properties, numerical stability, the degree of approximation and discretization errors.

Studies were carried out, by the writer, of the practical limitations of kinematic techniques in modelling physical systems and the validity of assuming single valued rating curves. These, however, are not discussed in this chapter but are reported later in the thesis. The chapter ends with several suggestions for further study and a summary of the results of various theoretical considerations of the general kinematic method.

#### 3.1 THEORY

Seddon (1900) was one of the first to report on what are now termed kinematic waves. His observations of flood waves on the Mississippi River formed the basis of his classical report. The term "kinematic wave" was applied by Lighthill and Whitham (1955) in a paper which provided a thorough discussion of the theoretical background for the phenomena. This paper outlined the basic assumptions used in their study of kinematic waves.

Briefly these assumptions are as follows:

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- The channel flows may be assumed to be incompressible and one-dimensional, i.e., zero flow component normal to the flow direction.
- 2. For each point in the channel there exists a single valued relationship between the flow rate and the cross section area

The continuity equation which describes the phenomena is written in the following manner:

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} = \overline{q}$$
(3.1)

Where: Q = flow rate A = Area of cross section  $\bar{q} = rate of lateral inflow$  x = distancet = time

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For the case of no lateral inflow, equation 3.1 can be written as:

$$\frac{dQ}{dy}\frac{\delta y}{\delta x} + T_{w}\frac{\delta y}{\delta t} = 0$$
 (3.2)

Where:  $T_{w}$  = top width

y = depth of water

To an observer moving downstream with a speed equal to the velocity of propagation of the wave, both depth, y and discharge, Q will appear to remain constant. Thus the total derivative of depth dy/dt is zero. Therefore:

$$\frac{dy}{dt} = 0 = \frac{dx}{dt}\frac{\delta y}{\delta x} + \frac{\delta y}{\delta t}$$
(3.3)

Substituting for dy/dt in equation 3.2 yields:

$$\frac{dx}{dt} = \frac{1}{T_w} \frac{dQ}{dy} = \frac{dQ}{dA}$$
(3.4)

That is, the wave velocity is equal to the slope of the curve relating flow rate and cross section area. (See Figure 3.1 for example.)

In Chapter 2, characteristic lines were defined as the paths in a space-time co-ordinate system along which changes in flow parameters may be described by ordinary differential equations. For flow situations described by the dynamic equation it can be shown that two sets of characteristic lines exist, being projected in the upstream and downstream direction respectively.

# FIGURE 3.1

RELATION BETWEEN WAVE VELOCITY AND STAGE FOR A TRIANGULAR CROSS SECTION





STAGE Y

$$\left(\frac{dx}{dt}\right)^+ = C^+ = V + \sqrt{\frac{gA}{T_w}}$$

$$\left(\frac{dx}{dt}\right)^{-} = C^{-} = V - \sqrt{\frac{gA}{T_w}}$$
(3.6)

In the kinematic wave situation only one set of characteristic lines exists--those projected in the downstream direction along paths defined by

$$\left(\frac{dx}{dt}\right)^{+} = \frac{dQ}{dA}$$
(3.7)

From equation 3.4 it may be seen that wave velocity is a function of the depth of flow alone and therefore kinematic waves are nondispersive and do not attenuate. However, they do change shape as a result of a variation of wave velocity with depth. The slope of the stage-discharge curves dQ/dA is usually steeper with higher flowrates. Thus higher flow rates (stages) move downstream faster than low flow rates. This variation in wave velocity can result in the intersection of characteristic lines from low flows and high flows. When this occurs a kinematic shock wave is formed. More discussion of kinematic shock waves is provided by Lighthill and Whitham (1955) and by Henderson (1966).

(3.5)

Lighthill and Whiteham also reported an additional criterion that must be satisfied in order for kinematic wave theory to apply. To prevent the formation of a hydraulic bore, the dynamic wave must attenuate. This attenuation will occur if the rate of change is given by:

$$\frac{dy}{dt} = \frac{gy_0 S_0 (2 - Fr) (I + Fr)}{3 V_0}$$
(3.8)

Where:

So = Bed slope

yo = Depth  
Vo = Velocity  
$$Fr = Froude number = Vo/\sqrt{gyo}$$

This imposes an upper limit on the rate of change of depth in the rising limb of the flood wave.

#### 3.2 APPLICATION OF KINEMATIC WAVE THEORY

The application of kinematic wave theory can be carried out by using either (i) a characteristics solution or (ii) by the direct application of finite difference methods to the continuity equation. Characteristic solutions have been proposed by several authors. Lighthill and Whitham (1955); Henderson and Wooding (1964). These algorithms are very simple and efficient for systems with simple geometry and constant lateral inflow; the main advantage is that they may be solved using analytic techniques. The extension of characteristic methods to allow for variation of lateral inflow in time and space could be accomplished, but this would prohibit the use of analytic techniques for obtaining solutions.

Numerous direct finite difference methods have been proposed and implemented. Several of these are reported by Kibler and Woolhiser (1970) and Brakensiek (1967). Direct methods offer a very flexible algorithm which provides the best approach to the development of a general kinematic method suitable for sections of arbitrary geometry.

The continuity equation expressed earlier is:

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} = \bar{q} \qquad (3.1)$$

and may be approximated by the finite difference expression.

$$\frac{\Delta Q}{\Delta X} + \frac{\Delta A}{\Delta T} = \bar{q}$$
(3.9)

Figure 3.2 shows a portion of a time-space diagram which is typical of those used in direct solutions of the continuity equation. In reviewing the approaches used by various workers, the basic difference appeared to be in the definition of the finite difference molecule. With the exception of the Lax-Wendroff method, all the techniques use

# FIGURE 3.2 GENERAL FINITE DIFFERENCE SCHEME

GRAPHICAL REPRESENTATION

NUMERICAL APPROXIMATIONS

1.



$$\frac{\delta A}{\delta T} = \frac{(I-\alpha)(A_4-A_3) + \alpha(A_2-A_1)}{\Delta T}$$

$$\frac{\delta Q}{\delta X} = \frac{\beta(Q_4 - Q_2) + (1 - \beta)(Q_3 - Q_1)}{\Delta X}$$

- MOLECULE THE SPACE-TIME ELEMENT BOUNDED BY POINTS 1243
- NUCLEUS THE POINT "P" ABOUT WHICH THE FINITE DIFFERENCE EQUATION IS APPLIED

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a rectangular "molecule" defined by the  $\Delta X$  and  $\Delta T$  steps of the space-time grid. The equation is applied at a point somewhere within or on the boundary of the rectangle. A molecule was, therefore, chosen which allows the point of application of the continuity equation to be varied and defined by two parameters  $\alpha$  and  $\beta$  as shown in figure 3.2. The algorithm based on this molecule provides a general framework which allows comparison with the other methods by the simple device of adjusting the values of the  $\alpha$  and  $\beta$  parameters.

The general method used to describe points in the space-time co-ordinate system usually employs a double subscript notation as indicated in figure 3.2.

Thus the upstream points of the molecule are located at points I, J and I, J+1 while the downstream points are at I + I, J and I + I, J + I. As a convenient short hand notation, the points are also numbered 1 through 4 as defined in figure 3.2 The unknown quantities are at location I + I, J + I or in short hand form, point 4. Point "P" defines the location about which the continuity equation is applied, hereafter termed the nucleus.

Rewriting the continuity equation in terms of the approximations shown in figure 3.2 yields the following equation:

$$\frac{\beta(Q_4 - Q_2) + (1 - \beta)(Q_3 - Q_1)}{\Delta X} + \frac{(1 - \alpha)(A_4 - A_3) + \alpha(A_2 - A_1)}{\Delta T} = \bar{q} (3.10)$$

Multiplying by  $\triangle X$  and collecting unknownsleads to:

$$\beta Q_{4} + (1-\alpha)A_{4} \frac{\Delta X}{\Delta T} = (\beta - 1)(Q_{3} - Q_{1}) + \alpha (A_{1} - A_{2}) \frac{\Delta X}{\Delta T} + (1-\alpha)A_{3} \frac{\Delta X}{\Delta T} + \beta Q_{2} + \tilde{q} \Delta X$$
(3.11)

Further rearranging gives:

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$$\beta Q_4 + (1-\alpha) A_4 \frac{\Delta X}{\Delta T} = \beta Q_3 + (1-\alpha) A_3 \frac{\Delta X}{\Delta T} - Q_3$$
$$-\beta Q_1 + \alpha A_1 \frac{\Delta X}{\Delta T} + \beta Q_2 - \alpha A_2 \frac{\Delta X}{\Delta T}$$
$$+ Q_1 + \bar{q} \Delta X \qquad (3.12)$$

In this equation the two unknowns  $Q_4$  and  $A_4$  are defined implicitly. Another equation relating flow rate and cross section area at point 4 allows a solution to be obtained for  $Q_4$  and  $A_4$  by iteration or by means of a technique which employs functional relationships and interpolation.

These functional relationships are written as:

$$f(Q) = \beta Q - \alpha A \frac{\Delta X}{\Delta T}$$
(3.13)

for the upstream points and

$$g(Q) = \beta Q + (1 - \alpha) A \frac{\Delta X}{\Delta T}$$
(3.14)

for the downstream locations.

Thus equation 3.12 can be rewritten as:

$$g(Q_4) = g(Q_3) - Q_3 + f(Q_2)$$
  
-  $f(Q_1) + Q_1 + \bar{q} \Delta X$  (3.15)

The solution of the unknown is obtained by computing g(Q) from the known quantities on the right hand side of the equation 3.15 and obtaining  $Q_4$  by interpolation from a curve relating  $Q_4$  and  $g(Q_4)$ . The solution may be continued either to the next molecule in the downstream direction or to the succeeding time increment. Thus the method may be used to define the conditions throughout the system at one time step or a complete time history can be determined for each elementary reach.

# 3.3 A CONCEPTUAL MODEL

In comparing special cases of kinematic flood routing methods, a conceptual model is helpful in visualizing the relationship of one algor-

ithm to another. The elementary reach which served as an example is shown in figure 3.3. Traditionally, storage in a channel is assumed to consist of two components, prism storage and wedge storage. Prism storage is considered to be a function of outflow whereas wedge storage is related to the amount by which inflow and outflow differ. However, in this study storage was thought of as two prisms. The storage in the downstream section is related to outflow, with storage in the upstream section a function of inflow. Storage in the upstream section is shown in figure 3.3 between sections "a" and "b" and is labelled  $ST_I$  . The volume of water stored in the reach between section "b" and "c" is labelled  $ST_0$ , and is a function of the outflow. Examining equation 3.12 reveals that it can be rewritten with the following substitutions.

$\alpha A_{I} \Delta X = ST_{II}$	(3.16)
$\alpha A_2 \Delta X = ST_{12}$	(3.17)
$(I-\alpha) A_3 \Delta X = ST_{03}$	(3.18)
$(I-\alpha) A_4 \Delta X = ST_{04}$	(3.19)



Thus equation 3, 12 can be written as:

$$\beta Q_4 + \frac{ST_0 4}{\Delta T} = \beta Q_3 + \frac{ST_0 3}{\Delta T} - Q_3 - \beta Q_1 + \frac{ST_{II}}{\Delta T} + Q_1$$

$$+ \beta Q_2 - \frac{ST_{I2}}{\Delta T} + \bar{q} \Delta X \qquad (3.20)$$

This shows the portions of the continuity equation which accounts for storage in the system and is helpful in considering two commonly used flood routing techniques which are special cases of the general kinematic method.

3. 3. 1 THE MUSKINGHUM METHOD AS A SPECIAL CASE OF KINEMATIC ROUTING.

The derivation of the Muskinghum method begins with the continuity equation expressed in the following manner:

$$I = O + \frac{\Delta ST}{\Delta T}$$
(3.21)

employing the notation conventionally applied to the method,

$$I = Inflow$$
  
O = Outflow

# ST = StorageT = Time

Another equation relating the amount of water stored in an elementary reach to the inflow and outflow is defined in the following manner:

$$ST = KO + Kx(I-O)$$
 (3.22)

The parameter K is a constant which relates storage to flow rate, while x is a factor which determines how much of the storage is related to outflow and how much is related to inflow. Prism storage is determined by the first term in the above equation and wedge storage is accounted for by the second term. This equation can be rewritten as:

ST = (1 - x)KO + xKI (3.23)

Setting the parameter x equal to zero makes storage a function of outflow alone. Storage is equally dependent on inflow and outflow when x = 0.5. Expressing the change of storage  $\Delta$ ST as a function of the flows defined on the space-time diagram of figure 3.2 yields the following:

$$\Delta ST = (I - x) KQ_4 + x KQ_2 - (I - x) KQ_3 - x KQ_1$$
 (3.24)

which may be expressed using expressions similar to those defined in equation 3. 16 - 3. 19.

$$\Delta ST = ST_{04} + ST_{12} - ST_{03} - ST_{11}$$
 (3.25)

Substituting equation 3.25 into equation 3.21 and defining inflow and outflow using the notation defined on figure 3.2 leads to:

$$\frac{Q_2 + Q_1}{2} + \bar{q}\Delta X = \frac{Q_4 + Q_3}{2} + \frac{ST_{04} + ST_{12} - ST_{03} - ST_{11}}{\Delta T}$$
(3.26)

From which the following is obtained:

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$$0.5Q_{4} + \frac{ST_{0}^{4}}{\Delta T} = 0.5Q_{3} + \frac{ST_{0}^{3}}{\Delta T} - Q_{3} - 0.5Q_{1} + \frac{ST_{11}}{\Delta T} + Q_{1} + 0.5Q_{1} - \frac{ST_{12}}{\Delta T} + \bar{q}\Delta X \qquad (3.27)$$

Comparing equation 3.27 with equation 3.20 reveals that they have a similar form. This shows that the Muskinghum method in a special case of the general method with the cross section area and flow rate related by a linear function which results in dQ/dA being constant for all values of Q.

The parameter x has the same meaning as the  $\alpha$  parameter in the general method. Values of the parameters are:

$$\beta = 0.5$$
 (3.28)

$$0.0 \le \alpha \le 0.5$$
 (3.29)

# 3.3.2 RESERVOIR ROUTING AS A SPECIAL CASE OF KINE-

The derivation of the numerical methods for reservoir routing begins with the same equation as is used for Muskinghum routing. That is:

$$I = O + \frac{\Delta ST}{\Delta T}$$

(3.21)
The storage in a reservoir is usually a function of the outflow alone, but for completeness storage will be expressed as a function of inflow and outflow using the following equation:

$$ST = (I - \alpha)ST_0 + \alpha ST_I$$
(3.30)

In order for storage to be a function of only outflow, **Q** must equal zero.

$$ST = ST_0 \tag{3.31}$$

and

.

$$\Delta ST = ST_{04} - ST_{03}$$
 (3.32)

Substituting equation 3.32 into equation 3.21 and expressing the terms as defined on the space-time diagram of figure 3.2 yields:

$$\frac{Q_2 + Q_1}{2} + \bar{q}\Delta X = \frac{Q_4 + Q_3}{2} + \frac{ST_{04} - ST_{03}}{\Delta T}$$
(3.33)

which can be expressed as:

$$0.5Q_4 + \frac{ST_04}{\Delta T} = 0.5Q_3 + \frac{ST_{03}}{\Delta T} - 0.5Q_1 + Q_1 + 0.5Q_2 + \bar{q}\Delta X - Q_3 \qquad (3.34)$$

The above equation is also similar in form to equation 3.20

with  $\boldsymbol{\beta} = 0.5$  and the storage independent of inflow. (i.e.  $\boldsymbol{\alpha} = 0.0.$ ) Setting  $\boldsymbol{\alpha} = 0.0$ , the functional relationships defined earlier (equations 3.13 and 3.14) become:

$$f(Q) = \beta Q \tag{3.35}$$

and

$$g(Q) = \beta Q + \frac{ST}{\Delta T}$$
(3.36)

Thus the general kinematic method may be used to route a flood through a reservoir where the characteristics of the storage element can be defined as a function of outflow. Traditionally the value of  $\beta$ chosen for this type of analysis is 0.5.

#### 3.3.3 FURTHER COMPARISONS.

Table 3. 1 provides further comparisons of the general kinematic method with several other methods in addition to the two cases discussed earlier. The first of these methods (1, 2 and 3) were proposed by Brakensiek (1967) in a paper which reported the results of numerical experiments to determine the properties of the three formulations.

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## COMPARISON OF FINITE DIFFERENCE SCHEMES

SCHEME	GENERAL	METHOD I	METHOD 2	METHOD 3	MUSKINGHUM	RESERVOIR
MOLECULE				$\frac{2}{1}$		
а	0.0± a ±1.0	0.5	0.5	0.5	0.0≤ (1 ≤ 0.5 (DEFINED AS X )	0.0
β.	0.0 <i>±</i> β ±l.0	0.5	1.0	0.0	0.5	0.5
$\frac{\delta Q}{\delta x}$	$\frac{\beta^{(Q_4-Q_2)+(1-\beta)(Q_3-Q_1)}}{\Delta^{\times}}$	<u>φ<sub>4</sub>-φ<sub>2</sub>+φ<sub>3</sub>-φ<sub>1</sub></u> 2Δ×	$\frac{Q_4 - Q_2}{\Delta X}$	$\frac{Q_3 - Q_1}{\Delta X}$	SAME AS METHODI	SAME AS METHOD I
<u>δA</u> δt	$\frac{(1-\alpha)(A_4-A_3)+\alpha(A_2-A_1)}{\Delta \times}$	$\frac{A_4 A_3 + A_2 - A_1}{2\Delta X}$	SAME AS METHOD I	SAME AS METHOD I	SAME AS GENERAL METHOD	$\frac{A_4 - A_3}{\Delta x}$

He found that the location of the nucleus has a marked affect on the behaviour of the algorithm and the solutions provided by the analysis. The three approaches can be modelled using the general kinematic method by defining  $\boldsymbol{\alpha}$  and  $\boldsymbol{\beta}$  as follows:

> Method 1 a = 0.5  $\beta = 0.5$ Method 2 a = 0.5  $\beta = 1.0$ Method 3 a = 0.5  $\beta = 0.5$  $\beta = 0.0$

Another approach to formulating a finite difference approximation to the continuity equation is known as the Lax-Wendroff method. This method employs four of the six nodal points defined by a double molecule and it is therefore not practical to provide a direct comparison with the general kinematic method. The arrangement of the molecule is shown in figure 3.4 and the finite difference equation is derived in Appendix E.

## FIGURE 3.4

FINITE DIFFERENCE MOLECULE FOR LAX-WENDROFF METHOD



- 60

## 3.4 GENERAL SIGNIFICANCE OF THE $\alpha$ and $\beta$ parameters.

By detailed analytic comparison with reservoir and Muskinghum methods and by graphical representation and comparison with other methods as in table 3. 1, it has been shown that most methods previously documented are special cases of the general method derived herein. By simple variation of the parameters  $\alpha$  and  $\beta$  one is able to emulate the numerical behaviour of these other methods. It is, therefore, instructive to attempt to gain some insight as to the physical and numerical significance of varying  $\alpha$  and  $\beta$  between their extreme values.

## 3.4.1 SPATIAL DERIVATIVES.

 $\delta Q/\delta x$  is the gradient of Q with respect to distance and due to the wave passing along the channel, must clearly be different at time t = T and time t = T +  $\Delta T$ .

The choice of  $\beta$  determines the manner in which the partial derivative  $\delta Q / \delta x$  is approximated in terms of the nodal values. With  $\beta = 0.0$ , the gradient  $\delta Q / \delta x$  is approximated using the conditions at time t = T. When  $\beta = 1.0$  conditions at time t = T +  $\Delta T$ are used to describe  $\delta Q / \delta x$  and with  $\beta = 0.5$  averages of the two previously mentioned values are used.

#### 3.4.2 TIME DERIVATIVES.

The second term in the continuity equation  $\delta A/\delta t$  describes the rate of change of cross section areas with respect to time. The way in which  $\delta A/\delta t$  is specified by the physical conditions is determined by the value of the parameter  $\alpha$ .

Using a value of  $\mathbf{a} = 0.0$  the gradient  $\delta A/\delta t$  is determined by conditions which exist at a position  $\mathbf{x} = \mathbf{X} + \mathbf{\Delta} \mathbf{X}$  while setting  $\mathbf{a} = 1.0$ will define  $\delta A/\delta t$  by conditions at  $\mathbf{x} = \mathbf{X}$ . Using  $\mathbf{a} = 0.5$  provides a value for  $\delta A/\delta t$  which is a simple average of the conditions at the ends of the increment being considered.

The physical significance of describing  $\delta A/\delta t$  at different locations within the molecule has been discussed previously in connection with the comparison of the Muskinghum method. Briefly, however, defining  $\delta A/\delta t$  at the downstream location in the molecule ( $\alpha = 0.0 \text{ or } x = X + \Delta X$ ) describes the storage in the reach of length  $\Delta X$  as a function of outflow alone. Similarly with  $\alpha = 1.0$  storage is dependent only on the inflow. When  $\alpha = 0.5$ ,  $\delta A/\delta t$  is obtained as a simple average of the values at each end of the molecule and storage is equally dependent on inflow and outflow.

### 3.5 TESTING OF THE KINEMATIC FLOOD ROUTING TECHNIQUES

There are four properties which determine the suitability of a finite difference scheme in representing a differential equation.

These are: 1. Stability

- 2. Degree of approximation
- 3. Discretization errors
- 4. Convergence

Before implementing a finite difference algorithm, it is necessary to check the suitability of the method in each of these four respects. Each of these questions must be answered through engineering judgment based on mathematical analysis and experience. The next portion of this chapter provides some discussion of analytic investigations and numerical experiments which demonstrate the applicability of and differences between the alternative systems.

A portion of the following text is devoted to a discussion of each of the properties which describe the performance of a finite difference scheme.

#### 3.5.1 STABILITY

A numerically stable procedure is a method in which small errors introduced into the calculations at a particular point in the x-t plane are not amplified as the solution is advanced through space and time. Errors caused by an unstable finite difference scheme will destroy the usefulness of the solution, thus the stability of the general kinematic method was the first item investigated. An analytic method of determining the stability of a finite difference scheme was successfully employed by several researchers, Walden (1973), Strelkoff (1970), Kibler et al (1970). This procedure, known as Van Neumann analysis, involves the investigation of a locally linearized version of the finite difference scheme on the assumption that the more complex non-linear system will behave in a similar manner to the linearized model.

Appendix A contains a detailed description of the various steps involved in determining a stability criterion. Briefly these steps are as follows. The finite difference scheme is first expressed in terms of the errors, which in turn are expressed as a Fourier Series. Because the system is linear the principle of superposition is applicable and only one component of the system need by examined at a time. If none of the harmonics of the Fourier Series are amplified in the succeeding computations, stability is attained.

The condition for stability may be expressed as follows:

$$\left|\frac{\tilde{A}_4}{\bar{A}_3}\right| = 1.0$$
 (3.37)

Where:	$\mathbf{\tilde{A}}_4$ = Error of the unknown
	$\mathbf{\overline{A}}_{3}$ = Error of the value preceeding the un-
	known .

The results of this analysis are tabulated in table 3.2 along with results by Kibler et al (1970) which provides the data for the Lax-Wendroff Scheme. These results indicate a trend toward de= creasing stability as the nucleus moves upstream and toward earlier time levels. The worst condition was found to occur when the equation is applied so that the finite difference approximations are obtained by backward differences, that is when  $\mathbf{a} = 1.0$  and  $\mathbf{\beta} = 0.0$ . The solution is unstable at this point regardless of the size of time and distance step used for computation.

Other points were located that provided solutions which were unconditionally stable as well as conditionally stable as shown in table 3.2.

#### 3.5.2 DEGREE OF APPROXIMATION

The degree of approximation provides a measure of how well the finite difference scheme represents the differential equation. If the differential equation is first order and finite difference schemes involve errors which are on the order of  $\Delta X$  and  $\Delta T$ , the finite differential equation may not converge to the differential equation as

 $\Delta X$  and  $\Delta T$  approach zero.

## **RESULTS OF STABILITY ANALYSIS**

$$KN = C \frac{\Delta T}{\Delta X}$$

GENERAL METHOD

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β	1.00	0.75	0.50	0.25	0.00
1.00	∞≥KN≥I	∞≐KN≐ <u>3</u> 4	∞ <sup>≥</sup> KN <sup>≥</sup> <sup>1</sup> 2	∞≥KN≥ <u> </u>	∞ <b>≚ KN ≥</b> 0
0.75	UNŞTABLE	KN = 1	$2 \ge KN \ge \frac{2}{3}$	3≥KN≥ <u> </u> 3	4≥KN≥0
0.50	UNSTABLE	UNSTABLE	KN = 1	<u>3</u> ≥KN≥ <u>1</u> 2	2 <b>≥</b> K N ≥ O
0.25	UNSTABLE	UNSTABLE	UNSTABLE	KN =	4 <u>3</u> ≥KN≥O
0.00	UNSTABLE	UNSTABLE	UNSTABLE	UNSTABLE	I ≥ KN≥ 0

LAX-WENDROFF: KN = I

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Analyzing a finite difference scheme involves expanding the finite difference equation using a Taylor series and determining what terms are truncated. Appendix C shows the process by which the order of approximation of the general kinematic method was determined.

The investigation showed that the degree of approximation involves first and higher order terms in  $\Delta X$  and  $\Delta T$  as shown by the error terms in the modified continuity equation (3.37).

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} + (2\alpha - i)O(\Delta X) + (i - 2\beta)O(\Delta T) + O(\Delta X^2, \Delta T^2) = \bar{q} \qquad (3.37)$$

The finite difference scheme provides the best approximation to the differential equation when  $\mathbf{a} = 0.5$  and  $\mathbf{\beta} = 0.5$ . When  $\mathbf{a} = 0.5$  and  $\mathbf{\beta} = 0.5$ , the equation for the finite difference representation of the differential equation becomes:

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} + O(\Delta X^2, \Delta T^2) = \bar{q}$$
 (3.38)

As the parameters  $\boldsymbol{\alpha}$  and  $\boldsymbol{\beta}$  are varied and the nucleus moves off the centre point of the molecule, errors on the order of

 $\Delta X$  and  $\Delta T$  are introduced. The absolute value of errors on the order of  $\Delta X$  and  $\Delta T$  introduced by varying these parameters increases linearly as the nucleus moves away from the centre of the molecule. Kibler et al (1970) reported that the Lax-Wendroff method represented the continuity equation in the following manner:

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} + O(\Delta X^2, \Delta T^2) = \bar{q}$$
 (3.39)

Thus the Lax-Wendroff method can be expected to behave in a manner similar to the general method with the equations applied to the centre of the molecule. However, the Lax-Wendroff molecule cannot be applied to obtain the downstream points on the time-space diagram because of the inverted "T" shape of the molecule. One of the other methods must be used for this portion of the grid.

#### 3.5.3 DISCRETIZATION ERROR

Discretization errors are a result of replacing a tangent (differential) with a chord (finite difference). These discrepancies are often analyzed using several numerical experiments to determine the sensitivity of the solution to the size of increments used in the computation. As the size of the finite difference increases the accuracy of the aforementioned approximation decreases.

This type of study is closely related to the analysis of the degree of approximation. Thus, based on the results of the previous section it is reasonable to expect the computer solutions of the system to be sensitive to the position of the nucleus as well as to the size of time and distance steps used.

Chapter 2 contains a description of two physical systems which were used to test the finite difference scheme for solving the dynamic equations. The results reported in this section are based on the first system described in the previous chapter. The problem can be described as a wide, rectangular channel which is subject to an inflow defined by a trapezoidal hydrograph. The particular values used in the simulation were:

The inflow hydrograph is shown in figure 2.5.

A computer program was written which provided a solution using the general kinematic flood routing method. For each execution of this program, fixed values of  $\Delta X$  and  $\Delta T$  were specified, which the routine used to provide twenty-five solutions with various values of the parameters  $\alpha$  and  $\beta$ . Several executions of this program provided the data necessary to carry out the sensitivity tests which demonstrated the effect of varying key parameters.

Tables 3. 3, 3. 4, and 3. 5 summarize the results of three computer runs. The elements of the matrix define positions of the nucleus

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## PEAK VALUES OF THE OUTFLOW HYDROGRAPH

#### SYSTEM 1

40,000 FT DOWNSTREAM

 $\Delta T = 200 \text{ SECONDS}$ 

 $\Delta X = 2500 \text{ FEET}$ 

βα	1.0	0.75	0.50	0,25	0.00
1,0	unstable	0.927	0.894	0.867	0.843
0.75	unstable	unstable	0.925	0.892	0.865
0.50	unstable	unstable	0.988	0.924	0.890
0.25	unstable	unstable	unstable	0.984	0.922
0.00	unstable	unstable	unstable	unstable	0.981

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## PEAK VALUES OF THE OUTFLOW HYDROGRAPH

#### SYSTEM 1

## 40,000 FT DOWNSTREAM

 $\Delta T = 200 \text{ SECONDS}$ 

 $\Delta X = 5000 FEET$ 

βα	1.00	0.75	0.50	0.25	0.00
1.00	unstable	0.969	0,894	0.844	0.804
0.75	unstable	unstable	0.924	0.865	0.821
0.50	unstable	unstable	0.965	0.890	0.840
0.25	unstable	unstable	unstable	0.920	0.861
0,00	unstable	unstable	unstable	0.961	0.886

### PEAK VALUES OF THE OUTFLOW HYDROGRAPH

### SYSTEM 1

40,000 FT DOWNSTREAM

 $\Delta T = 200 \text{ SECONDS}$ 

 $\Delta X = 10,000 \text{ FEET}$ 

βα	1,00	0.75	0.50	0,25	0.00
1.00	unstable	unstable	0.888	0.806	0.748
0.75	unstable	unstable	0.914	0.822	0.579
0.50	unstable	unstable	0.941	0.840	0.771
0.25	unstable	unstable	0973	0.859	0.784
0.00	unstable	unstable	unstable	0.881	0.799

and each is characterized by a parameter equal to the peak flow divided by full bank flow. Only stable solutions are reported.

Examination of the results show that the peak value decreased as the value of  $\boldsymbol{\alpha}$  was reduced and/or as  $\boldsymbol{\beta}$  increased. The rate at which the peak decreased or increased was related to size of the increments used in the computation.

This agrees with the results predicted by the analysis of the degree of approximation. As the size of the distance step increased, the peak value provided by the general method with the equation applied at a point where  $\boldsymbol{\alpha}$  was less than 0.5, decreased.

Similarly, for a fixed value of  $\Delta X$ , the peak value provided by the solution decreased as  $\mathcal{A}$  was varied from 0.5 towards 0.0. The decrease was not linear as predicted by the previous analysis. However, this prediction was based on a linear system. It is encouraging to note that the non-linear system does behave in a manner similar to the linear system.

Table 3.6 shows the ratios of peak outflow 40,000' downstream of the point of inflow obtained using the following fixed parameters:

 $\boldsymbol{\alpha} = 0.5$  $\boldsymbol{\beta} = 0.5$  $\boldsymbol{\Delta}T = 200 \text{ seconds}$ 

Values of  $\Delta X$  used were 2, 500, 5,000 and 10,000 feet.

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## PEAK VALUES OF THE OUTFLOW HYDROGRAPH

SYSTEM 1

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40,000 FEET DOWNSTREAM

 $\Delta T = 200 \text{ SECONDS}$ 

**α** = 0.5

**B** = 0.5

$\Delta$ X (FT)	PEAK FLOW RATIO
2500	0.988
5000	0.965
10000	0.941

Comparing the peak flow rates at the outflow with the peak flow rate modified to account for the error in interpreting the hydrograph reveals that discretization errors and other types of inaccuracies have not caused any error of the peak flow ratio at the location where the outflow was recorded with  $\Delta X = 2,500$ '. However, when the size of  $\Delta X$  was doubled to 5,000', a 2.4% reduction was made to the peak value. Furthermore, increasing the distance step to 10,000' introduced another 2.4% error to which is a total error of 4.8% from the value predicted by kinematic wave theory.

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Thus, based on comparisons of the peak value of the outflow hydrographs, using a solution with the continuity equation centred in the middle of the molecule and time and distance steps of:

> $\Delta T$  = 200 seconds  $\Delta X$  = 2, 500 feet.

the solution appears to be free of discretization errors.

It may be argued that considering only the peak value of the outflow hydrograph does not give a complete picture of the behaviour of the solution in reproducing a particular shape of hydrograph. To compare the shape of the hydrographs, several plots were made of hydrographs which had peaks that were approximately equal.

These hydrographs are shown in figures 3.5 and 3.6. Only one hydrograph was plotted on each graph because the results from the

# FIGURE 3.5 COMPARISON OF HYDROGRAPH SHAPES



# FIGURE 3.6 COMPARISON OF HYDROGRAPH SHAPES



TIME (HOURS)

different simulations compared so closely. The values of the parameters for these different simulations are shown on the plots.

As a result of these comparisons, it was concluded that for the purposes of this study, comparing the peak value of the hydrograph was a reasonable method of quickly determining the sensitivity of the solution to changes in the size of increments used in the solution and to the location of the point where the continuity equation was applied to the finite difference molecule. To provide a visual method of correlating the peak value of the shape of the hydrograph, figure 3.7 is provided. This graph contains outflow hydrographs which have a range of peak values from 0.748 to 0.988, and shows the variety of shapes obtained for the outflow hydrograph. Generally, as the peak value reduced the rising limb of the hydrograph started earlier and the falling limb dropped less rapidly

#### 3.5.4 CONVERGENCE

Convergence is a measure of how well a finite difference equation approximates the differential equation as the size of  $\Delta X$  and

 $\Delta$ T approach zero. Finite difference schemes may be classified as convergent, conditionally convergent or non-convergent.

If the finite difference scheme converges to the differential equation as  $\Delta X$  and  $\Delta T$  approach zero, the scheme is termed

# FIGURE 3.7 COMPARISON OF HYDROGRAPH SHAPES



TIME (HOURS)

convergent. A conditional convergent algorithm is one which only converges for particular values of  $\Delta X$  and  $\Delta T$ , while a non-convergent scheme will not accurately represent the differential equation regardless of the size of the finite difference steps.

Examination of the results of the degree of approximation analysis shows that the terms of the continuity equation are approx-

$$\frac{\delta Q}{\delta x} = \frac{\delta Q}{\delta x} + \frac{1}{2} (2\alpha - 1) \Delta x \frac{\delta^2 Q}{\delta x^2} + \frac{1}{6} (1 - 3\alpha + 3\alpha^2) \Delta x^2 \frac{\delta^3 Q}{\delta x^3}$$

$$+ \frac{1}{2} (\beta - \beta^2) \Delta T^2 \frac{\delta^3 Q}{\delta x \delta t^2} \cdots \qquad (3.40)$$

$$\frac{\delta A}{\delta t} = \frac{\delta A}{\delta t} + \frac{1}{2} (1 - 2\beta) \Delta T \frac{\delta^2 A}{\delta t^2} + \frac{1}{6} (1 - 3\beta + 3\beta^2) \Delta T^2 \frac{\delta^3 A}{\delta t^3}$$

$$+ \frac{1}{2} (\alpha - \alpha^2) \Delta x^2 \frac{\delta^3 A}{\delta x^2 \delta t} \cdots \qquad (3.41)$$

These equations show that as  $\Delta X$  and  $\Delta T$  approach zero, the finite difference equation will converge to differential equation. Furthermore, convergence of general kinematic flood routing method is independent of the values of the parameters  $\alpha$  and  $\beta$  and the relative sizes of  $\Delta X$  and  $\Delta T$ .

The fact that convergence is independent of  $\alpha$  and  $\beta$  verifies a conclusion made by Dooge (1959) that a series of reservoirs can be used to model a kinematic channel; further discussion of this will be provided in a later chapter.

## 3.6 VERIFICATION OF THE GENERAL KINEMATIC FLOOD ROUT -ING METHOD

The previous sections have each dealt with a particular aspect of finite difference techniques which are used to gauge the performance of a num erical solution of a differential equation. This section provides a comparison of a solution obtained from a finite difference algorithm with the results of an analytic solution as verification of the numerical method.

The analytic solution was obtained in the following manner: Kinematic waves have one set of characteristics which travel downstream with a velocity.

$$\frac{dx}{dt} = \frac{dQ}{dA}$$
(3.4)

Along a characteristic line, the flow rate Q is constant. Thus, associated with each flow rate there is a particular wave velocity. Knowing the wave velocity and the length between the point of inflow and the location where the outflow is being measured, the time between inflow and outflow for a particular value of Q can be obtained. Using Mannings equation, the flow rate can be expressed in the following manner:

$$Q = \frac{1.49}{n} \frac{A^{5/3}}{P^{2/3}} \sqrt{S}$$
 (3.42)

Where: n = roughness coefficient A = cross section area P = wetted perimeter S = slope

For the rectangular channel being studied:

$$A = T_{wx} Y$$
 (3.43)

$$P = T_w + 2Y$$
 (3.44)

Where:  $T_{W} = top width$ Y = depth of flow

From which:

$$\frac{dQ}{dA} = \frac{1.49\sqrt{S}}{nT_{w}} \frac{d}{dY} \left( (T_{w}Y)^{5/3} (T_{w} + 2Y)^{2/3} \right)$$

$$= \frac{1.49\sqrt{S}}{nT_{w}} \left( \frac{5}{3} \frac{T_{w}^{5/3}Y^{2/3}}{(T_{w} + 2Y)^{2/3}} + (T_{w}Y)^{5/3} (-\frac{4}{3}) (T_{w} + 2Y)^{-5/3} \right)$$
(3.46)

$$=\frac{Q}{T_w}\left(\frac{5}{3Y}-\frac{4}{3P}\right)$$
(3.47)

A computer program was utilized to solve for flow rate and kinematic wave velocity as a function of depth. Using the results of this program, a curve of wave velocity vs. flow rate/full bank flow rate was plotted. This is shown in Appendix F along with a listing of the computer program and results.

The inflow time of a particular flow rate was determined, and the wave velocity of the flow rate obtained from the graph. Multiplying the wave velocity by the length of the channel and adding the time of inflow, provided the time at which outflow of the particular flow rate would occur. By repeating this operation for a number of values in the range of flows specified by the inflow hydrograph, the outflow hydrograph was obtained.

Figure 3.8 shows a comparison of results obtained using both the analytic technique and the general kinematic flood routing method. These results compare very favourably and verify the useful ness of the general kinematic flood routing method in simulating the particular system under consideration when  $\mathbf{a} = 0.5$  and  $\mathbf{b} = 0.5$ .

#### 3.7 CONCLUSIONS

This chapter has provided a discussion of kinematic wave theory and has shown the development of a general method of routing a flood using kinematic wave theory.



TIME (THOUSAND SECONDS)

<sub>4</sub>84

With the framework provided by the general method of kinematic flood routing, it has been possible to compare several commonly used methods of flood routing. These techniques include reservoir routing, Muskinghum routing and several methods proposed by Brakensiek (1967).

The main feature of the general kinematic method, which allows flexibility in the choice of simulation method, is the ability to move the nucleus to any location on or within the boundaries of the finite difference molecule. This point is specified by two parameters  $\boldsymbol{\alpha}$ and  $\boldsymbol{\beta}$ . The value of  $\boldsymbol{\alpha}$  specifies the position of the nucleus in the space domain while  $\boldsymbol{\beta}$  locates the point in time. Most routing techniques utilize a finite difference molecule with  $\boldsymbol{\beta} = 0.5$ . The parameter  $\boldsymbol{\alpha}$  is set equal to zero when a reservoir routing technique is used. When the widely known Muskinghum method is simulated,

 $\boldsymbol{\beta}$  is set equal to 0.5 and  $\boldsymbol{\alpha}$  is given the value of the Muskinghum weight factor x.

Analytic studies and numerical experiments have been employed to determine the numerical behaviour of the general kinematic method. The two parameters  $\boldsymbol{\alpha}$  and  $\boldsymbol{\beta}$  have a very pronounced effect on the solutions provided with this method. Numerical errors and apparent attenuation may result depending on the values chosen for the

 $\alpha$  and  $\beta$  parameters. Amplification of the wave, as predicted by theoretical studies, was not reproduced by the numerical experi-

ments due to a peculiarity of the computer program which constrained the wave to full bank flow conditions. Further study should be undertaken utilizing a routine free from this constraint so that the theoretical predictions of amplification may be verified.

Stability also has been found to be dependent on the choice of

 $\boldsymbol{\alpha}$  and  $\boldsymbol{\beta}$ . For example, with the molecule defined to route a flood using the reservoir technique, the scheme is conditionally stable. However, when the space and time derivatives are evaluated using a backwards differencing method ( $\boldsymbol{\alpha} = 1.0$ ,  $\boldsymbol{\beta} = 0.0$ ), the algorithm is unstable regardless of the size of  $\Delta X$  or  $\Delta T$  used in the simulation. A typical location of the nucleus, which defines a conditionally stable routine, is found at the centre of the molecule

$$(\alpha = 0.5, \beta = 0.5).$$

With other properties being affected by the choice of  $\alpha$  and  $\beta$ , it is reasonable to expect the degree of approximation of the finite difference scheme to be sensitive to these two parameters. This has been verified through analysis of the degree of approximation. Introductions of errors on the order of  $\Delta X$  and  $\Delta T$  appear to be the cause of the pseudo-attenuation and amplification mentioned earlier. The amount of attenuation is related not only to the choice of  $\alpha$  and

 $\boldsymbol{\beta}$  , but also to the size of the finite space and time steps used in

the simulation. It has been shown that the best representation of a kinematic wave in an open channel is obtained by locating the nucleus in the centre of the finite difference molecule. With  $\mathbf{a} = 0.5$  and

 $\boldsymbol{\beta}$  = 0.5, the degree of approximation contains only second order and higher terms, no first order errors being introduced into the computations. This is comparable with the accuracy obtained using the Lax-Wendroff method.

Having studied the numerical properties of the general kinematic flood routing method, the next direction for further regard may be a comparison of the kinematic solutions with the more rigorous dynamic method.

From these comparisons, it may be possible to determine guidelines for choosing appropriate values of  $\boldsymbol{\alpha}$  and  $\boldsymbol{\beta}$  so that the properties of the general method may be more efficiently utilized in the simulation of dynamic flow systems. In addition, information may be obtained to determine the practical limitations of kinematic techniques.

#### CHAPTER 4

#### COMPARISON OF COMPLETE AND KINEMATIC SOLUTIONS

The two previous chapters have each dealt with a particular method of flood routing. The solution of the momentum and continuity equations describing unsteady flow was discussed in Chapter 2, and tests showed that it provided reasonable answers when modelling unsteady flow in open channels with simple geometry. Chapter 3 presented the theory and applications of kinematic methods used to simulate unsteady flow phenomena. The question now arises, "Which method is appropriate for the analysis of a specific physical system?"

In an attempt to answer this question, two different physical systems were modelled using both dynamic and kinematic analysis techniques. This chapter provides a comparison of these results and a discussion of the similarities and discrepancies of the results.

It has been pointed out in an earlier chapter that the dynamic method in general provides a more accurate answer than the kinematic solution. However, certain classes of problems may be identified in which the solutions obtained by both methods are not significantly different and it is not clear that any advantage results from the use of

more rigorous dynamic methods. One of the objectives in this chapter will be to investigate criteria which may aid in the classification of such problems.

The approach taken to achieve these objectives may be outlined in the following manner. First, a review of the momentum equation and the various terms of that equation is made, and a method for performing an order of magnitude analysis of the terms is presented. Secondly, the two physical systems that were utilized in the study are introduced and the importance of the terms in the momentum equation are discussed. Thirdly, the results of the simulations are presented and discussed. Also, a critique of current discussion in the literature is provided along with the framework developed in this chapter.

In conclusion, a brief review of the chapter is provided and suggestions for further study are made.

#### 4.1 ORDER OF MAGNITUDE ANALYSIS

Computation of flowrate in an open channel may be obtained typically by means of Manning's equation: i.e.

$$Q = A \frac{1.49}{n} \left(\frac{A}{P}\right)^{\frac{2}{3}} \sqrt{S_f}$$
(4.1)

Where: Q = flowrate

A = Cross section area

P = Wetted perimeter

 $S_f$  = Slope of the friction line

n = roughness coefficient

If a steady flow regime in a prismatic channel is being considered and the flow is assumed to be uniform, then the slope of the friction line must be the same as the bed slope. However, when an unsteady flow system is being considered, the friction slope is dependent on spatial and temporal acceleration, the variation of flow depth along the channel as well as the bed slope.

This can be seen by examining the dynamic equation expressed earlier in this thesis. It is:

$$\frac{\delta h}{\delta x} + \frac{\delta}{\delta x} \left( \frac{Q^2}{2gA^2} \right) + \frac{I}{gA} \frac{\delta Q}{\delta A} + S_f = 0$$
 (2.1)

Where: x = distance t = time h = water surface elevation Q = flow rate A = cross section area g = acceleration of gravity  $S_f$  = slope of the friction line

Expanding this equation leads to:

$$\frac{\delta y}{\delta x} - S_o + \frac{Q}{gA^2} \frac{\delta Q}{\delta x} - \frac{\delta y}{\delta x} \frac{Q^2 T_w}{gA^3} + \frac{1}{gA} \frac{\delta Q}{\delta t} + S_f = 0 \qquad (4.2)$$

Where: y = depth of flowSo = bed slope  $T_w = \text{surface width}$ 

The other variables have the same meaning as previously defined. Viewing each of the terms individually in the order they appear, they may be described as: (i) the rate of change of depth along the channel, (ii) the bed slope, (iii) the spatial acceleration, (iv) the rate of change of velocity head, (v) the temporal acceleration and (vi) the friction slope.

Solving for friction slope yields:

$$S_{r} = S_{o} - \frac{\delta y}{\delta x} + \frac{\delta y}{\delta x} \frac{Q^{2} T_{w}}{g A^{3}} - \frac{Q}{g A^{2}} \frac{\delta Q}{\delta x} - \frac{I}{g A} \frac{\delta Q}{\delta t}$$
(4.3)

This may also be expressed in the form

$$S_{f} = S_{o} - \frac{\delta y}{\delta x} \left(1 - F_{R}^{2}\right) - \frac{Q}{gA^{2}} \frac{\delta Q}{\delta x} - \frac{1}{gA} \frac{\delta Q}{\delta t}$$
(4.4)
Where:

$$F_R = Froude number = \frac{Q}{\sqrt{g\frac{A^3}{T_w}}}$$

Substituting equation 4.4 into equation 4.1 gives

$$Q = A \frac{1.49}{n} \left(\frac{A}{P}\right)^{\frac{2}{3}} \sqrt{S_o - \frac{\delta y}{\delta x} (1 - F_R^2) - \frac{Q}{g A^2} \frac{\delta Q}{\delta x} - \frac{1}{g A} \frac{\delta Q}{\delta t}}$$
(4.5)

Equation 4.5 shows the terms mentioned earlier in this section which describe the slope of the friction line which in turn is used to compute the flowrate for an unsteady flow system.

When a kinematic flood routing technique is employed, flowrate is assumed to be a function only of y and x, that is no time variation of the stage discharge relation occurs. Expressing this in a slightly different manner; flow rate at any section of the reach is assumed to be dependent only on the stage. Thus, for a single value stage-discharge curve to adequately describe a physical system, so that a kinematic solution is appropriate, it appears that terms other than the bed slope must be relatively small.

Spatial acceleration  $\frac{Q}{gA^2}\frac{\delta Q}{\delta x}$  and temporal acceleration  $\frac{1}{gA}\frac{\delta Q}{\delta t}$  are functions of time because they vary as the flood wave passes through a section. However, variation of the spatia! acceleration may also occur due to lateral inflow or outflow which may or may not be time dependent.

Similarly, the term describing the rate of change of specific energy due to a variation in depth with respect to distance will be a function of time as the flood passes, but it is also a description of the non-uniformity of flow under steady conditions. Thus, it is not necessary to assume that the single valued stage-discharge must be described by uniform flow conditions (i. e.  $S_f = S_0$ ). (Henderson, 1966.)

For the present discussion, all the terms describing the slope of the friction line will be considered time dependent with the exception of the bed slope.

Henderson (1963) reported an analysis that was made to compare the relative order of magnitude of the four slope terms in an equation similar to equation 4.5 using the kinematic wave as an approximation of the wave under consideration. The ensuing discussion follows the analysis presented by Henderson.

Defining kinematic wave velocity as c, equation 3.3 may be written as:

$$\frac{dy}{dt} = \frac{\delta y}{\delta t} + C \frac{\delta y}{\delta x} = 0$$
 (4.6)

This can be rewritten as:

$$-\frac{\delta y}{\delta x} = \frac{1}{C} \frac{\delta y}{\delta t}$$
(4.7)

For a system with a given inflow hydrograph  $Q = Q(\uparrow)$ , equation 4.7 can in turn be redefined in the following manner:

$$\frac{\delta y}{\delta x} = \frac{1}{C^2} \frac{\delta y}{\delta t} \frac{\delta Q}{T_y \delta y} = \frac{1}{C^2} \frac{\delta Q}{T_y \delta t}$$
(4.8)

Henderson's analysis continues by considering a wide rectangular channel in which the flow resistance is defined by the Chezy equation. The following result was obtained:

$$\frac{\delta y/\delta x}{S_0} \propto S_0^{-5/3} \times \frac{1}{q} \frac{\delta q}{\delta t}$$
(4.9)

Where: q = flow rate per unit width

When considering channels with arbitrary geometry and utilizing more complex resistance laws it is more difficult to obtain results in a form similar to equation 4.9. However, the results shown by equation 4.9 can also be deduced from the more general equation 4.8. As the bed slope is increased. equation 4.9 shows that the relative importance of  $\delta y/\delta x$  will decrease. Similarly, as q increases or the rate of change of q with respect to time gets smaller, the importance of  $\delta y/\delta x$  is reduced.

In general, c and  $T_w$  increase when flow rate increases; also, c increases when the slope is increased. Thus as slope or flow rate increase, the absolute value of the  $\delta y/\delta x$  term will decrease. Decreasing the rate of change of flow rate with respect to time will also reduce the absolute value of the  $\delta y/\delta x$  term.

Comparing the spatial acceleration  $\frac{Q}{gA^2} \frac{\delta Q}{\delta x}$  with  $\delta y / \delta x$  vields

$$\frac{\frac{Q}{gA^2}\frac{\delta Q}{\delta x}}{\frac{\delta y}{\delta x}} = \frac{Q}{gA^2}\frac{\delta Q}{\delta y} = O(F_R^2)$$
(4.10)

Where:

Also:

$$F_R = Froude number = \frac{Q}{\sqrt{g \frac{A^3}{T_w}}}$$

$$\frac{1}{gA}\frac{\delta Q}{\delta t} = O(\frac{C}{gA}\frac{\delta Q}{\delta x}) = O(\frac{Q}{gA^2}\frac{\delta Q}{\delta x})$$
(4.11)

Thus the two acceleration terms  $\frac{Q}{gA^2} \frac{\delta Q}{\delta x}$  and  $\frac{1}{gA} \frac{\delta Q}{\delta t}$  are

of the same order of magnitude and are of no higher order than  $\delta y/\delta x$ 

unless  $F_R >> 1$  which would occur only on extremely steep slopes such as mountain torrents.

In a system that has a channel with a very gentle bed slope,  $F_R <<1$ . Thus, while  $\frac{\delta y}{\delta x} / S_o$  may be appreciable, the acceleration terms could be negligible.

This led Henderson to classify open channel, unsteady flow systems into three categories depending on the importance of the four slope terms.

The first of these categories was the steep sloped system. In this case, all the terms except the bed slope are negligible when describing friction slope. The second category consists of an intermediate sloped system in which all four of the terms defining friction slope are necessary to provide an accurate description of the flow rate. The third classification consists of a gentle sloped system where the acceleration terms are negligible and the friction slope is defined only by the bed slope So and  $\frac{\delta y}{\delta x}$ .

While these classifications are termed according to the bed slope, this is not the only factor which determines the relative importance of the various terms. These other factors include channel roughness, geometry of the channel, and the inflow hydrograph. It may be possible for a particular physical system to be classified as being steep if the inflow hydrograph rises and falls very slowly or it could be classified as intermediate if the hydrograph is very rapid and the flows are such that  $F_R^2$  is large enough so that acceleration terms are important. If the flows are in a range where  $F_R^2$  is low and the acceleration terms are negligible, but the changes in flow rate are rapid, the same system may be classified as gentle sloped. A later portion of this chapter describes an attempt to use these criteria to determine the importance of the various terms in the dynamic equation and to demonstrate the method of classifying open channel unsteady flow systems.

#### 4.2 THE SYSTEMS STUDIED

The two examples used for the comparison of dynamic and kinematic flood routing techniques were very similar in some aspects, but quite diverse in other areas.

The first system has been discussed earlier in the sections dealing with sensitivity tests of the two methods of analysis, and comprised a wide rectangular channel having a downstream boundary condition of uniform flow and subject to a triangular inflow hydrograph as shown typically in figure 2.5 This system was used with several bed slopes. In each case a combination of bed slope and roughness was utilized which provided a consistent stage-discharge curve for all the results, i.e. $\sqrt{S_0}/n$  = constant. All comparisons were made at a section

.97

some distance upstream from the downstream boundary in an attempt to eliminate errors arising from the assumption of an unnatural downstream control.

The second system studied was taken from the problem first proposed by Thomas in 1934. Amein (1967) reports that the channel was assumed to be very wide with a sinusoidal inflow hydrograph starting from a base flow of 50 cfs/ft., and peaking at 200 cfs/ft. width. A time base of 96 hours was used with this hydrograph.

The main differences between the two systems were a difference of scale.

- 1. The physical length of the first problem is of the order of 1/40 that of Thomas' problem.
- 2. Although the flow rate per unit width of the inflow hydrographs were of the same order of magnitude, the first system employs a hydrograph with a time base of 2.77 hours compared to 96 hours for the second example. The effect of this is to greatly change the magnitude of the acceleration terms in the momentum equation.

A concise description of the variables used for each problem is provided in table 4.1 A comparison of the inflow hydrographs is shown in figure 4.1

# TABLE 4.1

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VARIABLE	SYSTEM 1	SYSTEM 2
distance to outflow position width depth slope roughness	40,000 ft 100 ft 20 ft 0.0002 - 0.0100 0.0149 - 0,1050	300 miles very wide 30. 1 ft 1/5280 0.02985
INFLOW HYDROGRAPH		
shape	triangular	sinusoidal
time base	2.77 hrs	96 hrs
peak	166.7 cfs/ft	200 cfs/ft
baseflow	33.3 cfs/ft	50 cfs/ft

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# COMPARISON OF SYSTEMS

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The result of an order of magnitude analysis for the two systems is shown in table 4.2 From this analysis, the systems may be classified in the following manner. When a bed slope of 0,0002 is used in System 1, all the terms of the momentum are of the same order of magnitude or greater than the bed slope, thus this system would be termed as an intermediate sloped system. With a bed slope of 0.001, the importance of the acceleration term is diminished and the system may be classified as either intermediate or gentle sloped, depending on what criterion is used to define terms as being significant. The simulation with the bed slope of 0.002 can be classified as gentle sloped because the acceleration terms are only approximately 5% of the bed slope and the  $\frac{\delta y}{\delta x}$ term is still significant. This demonstrates the statement made earlier that bed slope is not the only criterion for determining how a system is to be classified. Because the roughness coefficient was increased as the slope increased, the gentle sloped system was encountered at a slope greater than the intermediate sloped system.

To demonstrate the results obtained for a steep sloped system, where the bed slope is much larger than the  $\frac{\delta y}{\delta x}$  or acceleration terms, an execution of the program was made with So = 0.0100.

The order of magnitude analysis shows that System 2 would also be classified as a system with an intermediate slope, and from these

# TABLE 4.2

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# ORDER OF MAGNITUDE ANALYSIS

So	$\left  \frac{\delta y}{\delta x} \right $	$\left \frac{1}{gA}\frac{\delta Q}{\delta t}\right $	F <sup>2</sup>	δy  δx   So	$\frac{\left \frac{1}{gA}\frac{\delta Q}{\delta t}\right }{\left \frac{\delta y}{\delta x}\right }$			
	SYSTEM I							
0.0002				0.875 - 2.37				
0.0010	1.75×10 <sup>-4</sup>	0.41 x 10 <sup>-4</sup>	0.099	0.175 - 0.474	0.087			
0.0020	4.74 x 10 <sup>-4</sup>	1.16×10 <sup>-4</sup>	0.107	0.0875 - 0.237	0.67			
0.0100				1.75×10 <sup>-2</sup> - 4.74×10 <sup>-2</sup>				
SYSTEM 2								
0.00019	2.4 × 10 <sup>-3</sup>	0.28x10 <sup>-4</sup>	0.035-0.045	1.28 × 10 <sup>-1</sup>	0.11x10 <sup>-1</sup>			

results, one would expect less attenuation of the flood wave in System 1 as the bed slope is increased. While a comparison between the two systems is difficult due to the differences of scale, one would expect a less marked attenuation of the flood wave in System 2 when compared with the attenuation obtained with System 1 using a bed slope of 0.0002, due to the relative sizes of  $\frac{\delta y}{\delta x}$  and bed slope.

#### 4.3 COMPARISON OF RESULTS FOR SYSTEM 1

The general kinematic flood routing technique with the nucleus located at the centre of the molecule  $(a=0.5, \beta=0.5)$  was used to obtain a kinematic solution for System 1. Because the ratio of bed slope and roughness was adjusted to maintain a constant ratio of  $\sqrt{S_0/n}$ , only one kinematic solution was required for the various bed slopes.

Figure 4.2 contains a comparison of the outflow hydrographs obtained from the four dynamic simulations of System 1, and from the kinematic solution. This clearly demonstrates the pronounced effect of the time variant terms on the resultant outflow hydrograph. The simulation of System 1 with a bed slope of 0.0002 shows the effect of terms which are of the same order of magnitude as the bed slope. In this case, the wave reached the outflow quicker than the other waves and subsided less rapidly. As slope increased, the attenuation of the



TIME (HOURS)

wave decreased and the wave form approached the shape predicted by kinematic wave theory.

Another demonstration of the effects of the terms in the dynamic equation can be seen by examining the stage-discharge curves in figure 4.3. With increased bed slope and roughness, the slope of the energy line was dominated by the bed slope term. Thus the hysteresis loop in the stage-discharge decreases as shown in the figure.

#### 4.4 COMPARISON OF RESULTS FOR SYSTEM 2

An approach, similar to the one utilized for comparing the results of System 1, was used to compare the results of this system. However, hydrographs at two sections were recorded and peak values at a point half way between the two hydrographs were also printed to provide further comparisons. This allowed the results of this study to be viewed in conjunction with the results of two other studies, Amein (1967), and Garrison (1968).

A comparison of the hydrograph 100 miles downstream from the upper end of the channel as obtained by kinematic and dynamic solutions is shown in figure 4.4. The attenuation of the flood wave is apparent; but differences in the wave shape are relatively negligible.

However, by the time the flood wave has reached a point 300 miles downstream, there are considerable differences in the wave





shapes resulting from the two methods of analysis as shown in figure 4.5. The wave front on the kinematic solution is much steeper and no attenuation has taken place. Conversely, the dynamic solution shows a wave which has steepened slightly, and has attenuated. The stagedischarge curve for this system 300 miles downstream is shown in figure 4.6.

From the order of magnitude analysis and a comparison with simulations of the first system, more attenuation may have been expected because of the high value associated with the  $\frac{\delta y}{\delta x}$  term. However, this may be explained by considering the manner in which the order of magnitude of the  $\frac{\delta y}{\delta x}$  term was obtained. In System 1, the trapezoidal inflow hydrograph provided a rate of change of flow rate  $\frac{\delta Q}{\delta t}$  that was constant when the hydrograph was rising or falling. With the sinusoidal inflow hydrograph,  $\frac{\delta Q}{\delta t}$  varied with time and the maximum rate of change of flow rate was used to determine the order of magnitude of the  $\frac{\delta y}{\delta x}$  term. During the major portion of the simulation,  $\frac{\delta y}{\delta x}$  was actually much less than the value specified in table 4. 2. This may account for the relatively small amount of attenuation demonstrated in the results of this simulation.

#### 4.5 FURTHER COMPARISONS

Kinematic solutions of unsteady open channel flow problems were first utilized by engineers and researchers before the advent of high





speed digital computers due to the enormous amount of calculations required to solve the dynamic equations. While these shorthand methods reduced the computational load, it was still necessary to have large increments in time and space to reduce the number of calculations required for the numerical solution by a kinematic technique. The use of step sizes which were too large may have provided solutions which were not realistic or which gave results that were extremely sensitive to the step sizes in time or distance.

Recording the peak values of depth and flowrate and the times at which these maximums occurred at a section 200 miles downstream in System 2, allowed a comparison to be made with results of the studies by Amein (1968) and Graves (1967). Table 4.3 contains a table which shows a comparison of peak values for flowrate and depths obtained using the various techniques.

When a kinematic solution was used with the nucleus in the centre of the molecule, no attenuation resulted. Moving the nucleus to the downstream boundary so that the channel was simulated by a series of reservoirs, caused some attenuation of the flood wave. The size of  $\Delta X$  dictated the amount of attenuation that was manifested in the kinematic solution with the nucleus located off of the centre of the molecule. There appears to be considerable difference in the timing of the peak values as the sizes of  $\Delta X$  and  $\Delta T$  were varied. This,

TABLE 4.3

# COMPARISONS TWO HUNDRED MILES DOWNSTREAM SYSTEM 2

METHOD	ΔΧ	ΔT	FLOW	RATE	DEPTH	
	(MILES)	(HOURS)	PEAK (CFS/FT)	TIME (HOURS)	PEAK (FEET)	TIME (HOURS)
!	10	2	199.9	75.1	30.06	75.1
KINEMATIC	10	3	200.0	74.6	30.07	74.6
<b>α</b> = 0.50 <b>β</b> =0.50	20	2	200.2	74.8	30.07	74.8
	20	3	<b>200</b> .0	74 8	30.06	74.8
· ·		2	194.2	75.1	29.54	75.1
KINEMATIC		3	194.4	75.2	29.55	75.2
<b>β</b> = 0.50	20	2	188.0	75.8 <sup>°</sup>	28.97	75.8
		3	188.1	75.9	28.98	75.9
EXPLICIT DYNAMIC SOLUTION	10	0.15	191.0	75.1	29.18	76.9
STORAGE ROUTING AMEIN (1968) METHOD PROPOSED BY THOMAS(1934)	25	12	_		29.2	78.0
		24			27.8	90.0
	12	12	_	_	27.7	78.0
		24	-		26.0	84.0
GRAVES(1968)	50	1	193.9	76.0	29.5	77.6

however, is due to the relatively flat top of the wave which resulted in the peak being ill defined. The depth did not vary by more than a few hundredths of a foot several time steps before or after the peak. The average value of the time of peak is 75.25 hours and the values obtained for individual runs do not vary from this value by more than one time step. The results obtained with the kinematic solution also compare favourably with the answers provided by the finite difference solution of the dynamic equations. While the attenuation was not accurately modelled, the times to the peak values compare reasonably well.

Amein (1968) reports the results of Thomas' storage-routing technique in a paper describing an implicit finite difference solution of the dynamic equation. He pointed out that these results seem to indicate that the storage routing technique is unacceptable due to the variation of the values recorded for peak stage and the time of peak. These problems may be symptoms of using large steps resulting in the solution becoming unstable. In all of the cases cited, the time of peak did not vary from the peak predicted using the dynamic solution by more than one time step. Thus the kinematic and the dynamic solutions are similar in that there are maximum sizes of  $\Delta X$  and  $\Delta T$  that can be used to achieve satisfactory results.

Graves (1967) used a flood routing technique based on the continuity equation and a stage discharge curve that was a function of the steady

flow conditions and the rate of change of flow with respect to time. This type of solution is neither kinematic or dynamic in a strict theoretical sense. However, it is noteworthy in that it is an example of methods which attempt to bridge the gap between kinematic solutions and the complete solution of the dynamic equation. The results of his study show a peak flow rate of 193.9 cubic feet per second per foot channel width occurring at a point 200 miles downstream 76 hours after the start of the simulation. The maximum stage was 29.46 feet 78 hours after the start of the simulation. These results compare quite favourably with the dynamic solution.

In making these comparisons, two problems have been manifested which are worthy of further discussion. First, when comparisons of results are obtained using numerical flood routing models, it is sometimes advantageous to compare stages while other situations warrant a comparison of flow rates. In reporting results of models which do not assume a single valued stage-discharge relation, it may be expedient to report both stage and flow rates. The other problem encountered dealt with the size of the time step used in the computation. Difficulties associated with using large step sizes in a finite difference scheme are well documented and are usually guarded against. However, the type of flood wave that was encountered in System 2 had a relatively flat top which made locating the true time of peak a difficult task. Using smaller time steps would reduce this uncertainty; but would increase the computation costs. Another approach to overcoming this problem would be to fit a curve to the data points in the region of the peak using either graphical methods or a numerical technique. The latter angle of attack was used with data where the peak did not appear to fall close to one of the time increments. A second order curve was utilized in this attempt. The equation has the form:

$$Q = a + bt + ct^2$$
 (4.12)

Where:

a, b, c, = coefficients determined from the data Q = flow ratet = time

The results shown in table 4.3 have been refined by this interpolation technique.

#### 4.6 CONCLUSIONS

In this chapter, a review of the terms in the momentum equation has been made and a method for determining the order of magnitude of these terms has been investigated. From this study, several broad qualitative classifications of unsteady flow phenomena have been demonstrated. These form a preliminary basis for determining the suitability of applying a kinematic solution to a particular unsteady flow situation.

The classifications, first proposed by Henderson, are:

(1) STEEP SLOPE - Only the bed slope is significant in this type of system. The absolute value of  $\frac{\delta y}{\delta x}$  is estimated using the equation:

$$\left|\frac{\delta y}{\delta x}\right| = \left|\frac{1}{C^2}\frac{\delta Q}{T_{0}\delta t}\right|$$
(4.8)

If the quantity

$$\left|\frac{\delta y}{\delta x}\right| / S_{o} <<1$$
(4.13)

the system may be classified as steep.

(2) GENTLE SLOPE - With this classification, only the bed slope and  $\frac{\delta y}{\delta x}$  terms of the dynamic equation are significant. It has been shown that if the system does not fit into the steep slope classification and

$$F_R^2 << 1$$
 (4.14)

the system may be classified as gentle sloped.

(3) INTERMEDIATE SLOPE - This system is encountered when all the terms in the momentum equation are important. A system which cannot be classified with the two previous categories will fit under this title.

Results of several simulations were presented which showed the comparisons of dynamic and kinematic solutions for various classifications of systems. As the bedslope became more significant in relation to the other terms, the solutions tended toward a kinematic solution. However, attenuation, which resulted from the dynamic effects in the system, was detected even when the bedslope was two orders of magnitude greater than the  $\frac{\delta y}{\delta x}$  term. It is felt that this observation is not a general observation that can be applied to other cases of unsteady flow phenomena. Other systems, which have a longer wave length hydrograph, may have much less attenuation.

It should also be noted that the classifications steep, intermediate or gentle slope are not the result of a particular bedslope in a channel. These classifications also reflect the flow resistance, channel geometry, ranges of flows encountered in the hydrograph and the rate of change of the flow rates. This was demonstrated by studying a system which, for a particular inflow hydrograph, changed from intermediate slope to gentle slope and finally to a steep slope system, as the bedslope and roughness were increased.

Comparisons of results from this study were made with results reported by Amein (1968) and Graves (1967) to demonstrate problems

associated with using excessively large time steps. The difficulties were manifested in the form of instabilities and in data interpretation. This demonstrates that both the kinematic methods and the dynamic methods have limitations to the sizes of  $\Delta X$  and  $\Delta T$ that can be successfully employed to provide a solution. An interpolation technique was employed to overcome problems associated with data interpretation.

The method proposed by Graves is interesting in that it is an example of methods which are employed to model the attenuation and movement of a flood wave with a quasi-kinematic approach. Another example of a technique which employs an approximation in conjunction with kinematic flood routing to model movement and attenuation is the kinematic solution with the nucleus moved from the centre of the molecule. These types of models have not been discussed at this point in the thesis. However, the next chapters present a discussion of several methods that may be utilized to extend the usefulness of kinematic models.

# CHAPTER 5

#### ATTENUATION AND KINEMATIC METHODS

This chapter provides a discussion of a method of modelling attenuation in conjunction with the general kinematic flood routing algorithm. The previous chapter, which contains a comparison of kinematic and dynamic solutions of two physical systems, demonstrates the ability of the two methods to model attenuation. The dynamic solution follows from a rational numerical representation of the actual physical phenomena. Kinematic methods, per se, do not predict attenuation but depend on a manipulation of the computational scheme to approximate the reduction of the flood wave, which may not, in fact, truly represent the process that is occurring in the prototype.

Thus, while dynamic techniques provide a more realistic method of predicting attenuation, it may be possible and computationally advantageous to utilize a modified kinematic routine, if the solutions can be shown, by a calibration process, to be valid.

In presenting this discussion three components of modelling attenuation will be explored. These are:

119

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- "Molecule effect:" the attenuation introduced by moving the nucleus from the centre of the molecule.
- "Cascade effect:" the way in which attenuation is affected by varying the number of storage units in series.
- "Storage effect:" the dependance of the attenuation on the relative magnitude of live storage and volume in the flood hydrograph.

To show the usefulness of this approach in the simulation of unsteady open channel flow, a number of numerical experiments are presented and the results are compared with the corresponding dynamic solutions. A description of the way in which an engineer might calibrate a model of this type is included along with the concluding remarks.

# 5.1 THE "MOLECULE EFFECT" - MODELLING ATTENUATION BY MOVING THE NUCLEUS

In designing the numerical experiments used with this study, two distinct physical systems were employed. These systems, described in detail earlier in this thesis were:

(i) a rectangular channel 20' deep and 100 ' widesubjected to a triangular shaped inflow hydrograph

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(ii) a very wide channel subjected to a sinusoidal inflow hydrograph.

The primary difference between the two systems was the scale. The time base of the inflow hydrograph for the first system was 2.77 hours while the second hydrograph had a time base of 96 hours. The length of the channel for the first system was 40,000 feet whereas a channel 300 miles long was used in the second system.

To examine the sensitivity of the flood routing solution to the position of the nucleus within the molecule, a program was employed which for fixed values of  $\Delta X$  and  $\Delta T$ , carried out a routing analysis for a range of values of  $\alpha$  and  $\beta$ . Both parameters were varied between 0.0 and 1.0 with increments of 0.25, so that a total of 25 possible solutions were obtained. It was recognized at the outset, that a significant number of these nucleus positions would not result in an acceptable solution, because of numerical instability at certain values of  $\alpha$  and  $\beta$ .

The results for the first system are shown in Tables 3.3 and 3.4 and 3.5 of Chapter 3. The elements of the matrix define the positions of the nucleus and the values contained in each element are the ratios of the peak outflow divided by the full bank flow rate. A similar set of results was compiled for the second system in tables 5.1, 5.2, 5.3 and 5.4.

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# PEAK VALUES OF THE OUTFLOW HYDROGRAPH

# SYSTEM 2

## 300 MILES DOWNSTREAM

# $\Delta T = 2$ HOURS

# $\Delta X = 10$ MILES

βα	1.00	0.75	.0.50	0.25	0,00
1.00	0.976	0.954	0.932	0.910	0.888
0.75	unstable	0.986	0.964	0.942	0.919
0.50	unstable	unstable	0.996	0.975	0.952
0.25	unstable	unstable	unstable	unstable	0.986
0.00	unstable	unstable	unstable	unstable	unstable

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# PEAK VALUES OF THE OUTFLOW HYDROGRAPH

### SYSTEM 2

### 300 MILES DOWNSTREAM

 $\Delta T = 2$  HOURS

# $\Delta X = 20$ MILES

βα	1.00	Q. 75	0.50	0.25	0.00
1.00	unstable	0.975	0.931	0.887	0.849
0.75	unstable	unstable	0.963	0.918	0.875
0.50	unstable	unstable	0.997	0.951	0.904
0.25	unstable	unstable	unstable	0.985	0.938
0.00	unstable	unstable	unstable	1.000	0.972

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# PEAK VALUES OF THE OUTFLOW HYDROGRAPH

## SYSTEM 2

### 300 MILES DOWNSTREAM

 $\Delta T = 3 HOURS$ 

 $\Delta X = 10$  MILES

βα	1.00	0.75	0.50	0.25	0.00
1.00	0.946	0.924	0.902	0.880	0.860
0.75	unstable	0.970	0.949	0.927	0.903
0,50	unstable	unstable	1.000	0.975	0.954
0,25	unstable	unstable	unstable	unstable	1,000
0.00	unstable	unstable	unstable	unstable	unstable

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### PEAK VALUES OF THE OUTFLOW HYDROGRAPH

# SYSTEM 2

### 300 MILES DOWNSTREAM

 $\Delta T = 3 HOURS$ 

 $\Delta X = 20$  MILES

βα	· 1.00	0.75	0.50	0.25	0.00
1.00	unstable	0.944	0.901	0.860	0.825
0,75	unstable	0.990	0.948	0.903	0.860
0.50	unstable	unstable	0.995	0.952	0.904
0.25	unstable	unstable	unstable	0.999	0.956
0.00	unstable	unstable	unstable	unstable	1.00

Before proceeding with the theoretical aspects of the "molecule effect" it is vital that the reader obtain a clear picture of the phenomena that is manifested by these results. If a three dimensional plot were made, showing the peak outflow versus α and , a smooth surface would be depicted. This surface may be B visualized by imagining contour lines across the tables which contain these values. For example, the results shown in table 3.3 for System 1, would form a surface with the line of steepest descent going from the central position toward the upper right hand corner of the table ( $\mathbf{\alpha} = 0.0$ ,  $\boldsymbol{\beta} = 1.0$ ) with the gradient of this line becoming shallower in the proximity of the corner. Similarly, table 5.1, which shows some results from System 2, presents data which would form a surface with the steepest descent upwards and to the right. If the strike of the surface were plotted to determine the direction of the dip, the strike being a horizontal line perpendicular to the dip or line of steepest descent, the direction of the dip would be found to be three increments upwards and two increments to the right. Examination of the other tables (i. e. 5.2 to 5.4) reveals a similar result, although it can be seen that as the size of  $\Delta X$ and  $\Delta T$  varies, the direction of the dip changes. When  $\Delta X'$  is made larger, the dip swings towards the right hand side; but if

 $\Delta T$  is increased, the dip points toward the upper portion of the matrix.

If the necessary key factors could be identified enabling an engineer or researcher to describe the characteristics of these surfaces without having to generate numerous computer simulations, it would then be possible to preselect the necessary values of  $\boldsymbol{\alpha}$ ,

 $\boldsymbol{\beta}$  etc. in order to produce attenuation of a required amount.

The following theoretical investigation is directed towards identification of the parameters which characterize the P( $\alpha$ ,  $\beta$ ) surfaces described above and towards rationalizing the introduction of numerical error, or "molecule effect" that is brought about by moving the nucleus away from the centre of the molecule.

Examination of the finite difference approximation of the continuity equation in differential form has revealed that the representation is given by equation 5. 1.

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} + \frac{1}{2}(2\alpha - 1)\Delta X \frac{\delta^2 Q}{\delta x^2} + \frac{1}{2}(1 - 2\beta)\Delta T \frac{\delta^2 A}{\delta t^2} + O(\Delta X^2, \Delta T^2)$$
  
=  $\bar{q}$  (5.1)

The mathematical proof of this statement is presented in Appendix C.

Ignoring error terms higher than first order allows the equation to be written as:

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} + \epsilon = \bar{q}$$
 (5.2)
where:

$$\epsilon = \frac{1}{2} (2\alpha - I) \Delta X \frac{\delta^2 Q}{\delta x^2} + \frac{1}{2} (I - 2\beta) \Delta T \frac{\delta^2 A}{\delta t^2}$$
(5.3)

To calculate the sensitivity of the size of the error terms to variations of  $\alpha$ ,  $\beta$ ,  $\Delta X$  and  $\Delta T$ , it is necessary to analyze the nature of the mathematical model. This information can be obtained by determining the relative size of the terms  $\frac{\delta^2 Q}{\delta x^2}$  and  $\frac{\delta^2 A}{\delta x^2}$ 

Describing the kinematic wave by the partial differential equation:

$$\frac{\delta y}{\delta t} + C \frac{\delta y}{\delta x} = 0$$
 (5.4)

implicitly defines the wave celerity as

$$C = \frac{-\delta y / \delta t}{\delta y / \delta x} = -\frac{\delta x}{\delta t}$$
(5.5)

Rearranging the differential form of the continuity equation yields:

$$\frac{\delta Q}{\delta x} = \bar{q} - \frac{\delta A}{\delta t}$$
(5.6)

or:

$$\frac{\delta A}{\delta t} = \bar{q} - \frac{\delta Q}{\delta x}$$
(5)

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Further differentiation of equations 5.6 and 5.7 leads to the second order terms of 5.8 and 5.9.

$$\frac{\delta^2 Q}{\delta x^2} = -\frac{\delta^2 A}{\delta x \, \delta t} \tag{5.8}$$

$$\frac{\delta^2 \dot{A}}{\delta t^2} = -\frac{\delta^2 Q}{\delta t \, \delta x}$$
(5.9)

Multiplying equation 5.8 by the kinematic wave velocity gives:

$$C\frac{\delta^2 Q}{\delta x^2} = -C\frac{\delta^2 A}{\delta x \delta t}$$
(5.10)

from which:

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$$C\frac{\delta^2 Q}{\delta x^2} = -\left(\frac{\delta x}{\delta t}\right)\frac{\delta A}{\delta x \,\delta t}$$
(5.11)

$$C\frac{\delta^2 Q}{\delta x^2} = \frac{\delta^2 A}{\delta t^2}$$
(5.12)

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This leads to several important points. First, by substituting equation 5.12 into the error expression 5.3 the following equation is obtained:

$$\epsilon = (2\alpha - 1)\frac{\Delta X}{2}\frac{\delta^2 Q}{\delta x^2} + (1 - 2\beta)\frac{C\Delta T}{2}\frac{\delta^2 Q}{\delta x^2}$$
(5.13)

$$=\frac{\delta^2 Q}{\delta x^2} \left( (\alpha - 0.5) \Delta X + (0.5 - \beta) C \Delta T \right)$$
(5.14)

This demonstrates the qualitative dependance of the error term on linear terms in  $\alpha$ ,  $\beta$ .

The second item that can be developed from the previous algebraic manipulations is a finite difference representation of the error terms.

By referring to figure 3.2 for a definition of nodal values in the finite difference molecule, it can be seen that:

$$\frac{\delta^2 A}{\delta t^2} = -\frac{\delta^2 Q}{\delta x \delta t} = -\left(\frac{(Q_4 - Q_3)}{\Delta T} - \frac{(Q_2 - Q_1)}{\Delta T}\right)$$
(5.15)

$$= -\left(\frac{Q_4 - Q_3 + Q_1 - Q_2}{\Delta X \Delta T}\right)$$
(5.16)

Similarly:  

$$\frac{\delta^2 Q}{\delta x^2} = -\frac{\delta^2 A}{\delta x \, \delta t} = -\left(\frac{(A_4 - A_3)}{\Delta T} - \frac{(A_2 - A_1)}{\Delta T}\right) \qquad (5.17)$$

$$= -\left(\frac{A_4 - A_3 + A_1 - A_2}{\Delta X \, \Delta T}\right) \qquad (5.18)$$

Using the following finite difference approximations:

$$\frac{\delta Q}{\delta x} = \frac{Q_4 + Q_3 - Q_2 - Q_1}{2\Delta X}$$
(5.19)

and:

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$$\frac{\delta A}{\delta t} = \frac{A_4 + A_2 - A_3 - A_1}{2\Delta T}$$
(5.20)

equation 5.1 may be rewritten in finite difference form. The equation is:

$$\frac{Q_4 + Q_3 - Q_2 - Q_1}{2\Delta X} + \frac{A_4 + A_2 - A_3 - A_1}{2\Delta T}$$
$$+ \frac{1}{2} (2\alpha - 1) \Delta X \left( \frac{-(A_4 - A_3 + A_1 - A_2)}{\Delta X \Delta T} \right)$$
$$+ \frac{1}{2} (1 - 2\beta) \Delta T \left( \frac{-(Q_4 - Q_3 + Q_1 - Q_2)}{\Delta X \Delta T} \right) = \bar{q}$$

131

(5.21)

Collecting terms yields:

$$\frac{Q_4 + Q_3 - Q_2 - Q_1}{2\Delta X} + \frac{1 - 2\beta}{2} \frac{Q_3 + Q_2 - Q_4 - Q_1}{\Delta X}$$

$$\frac{A_4 + A_2 - A_3 - A_1}{2\Delta T} + \frac{2\alpha - 1}{2} \frac{A_3 + A_2 - A_4 - A_1}{\Delta T} = \bar{q} \qquad (5.22)$$
Further simplification gives:

$$\frac{\beta(Q_4 - Q_2) + (1 - \beta)(Q_3 - Q_1)}{\Delta X}$$

$$\frac{(1-\alpha)(A_4-A_3) + \alpha(A_2-A_1)}{\Delta T} = \bar{q}$$
 (5.23)

It can be seen that equation 5.23 is precisely the same as equation 3.10. This leads to the conclusion that the error term related to the selection of  $\boldsymbol{\alpha}$  and  $\boldsymbol{\Delta X}$  has physical significance in that it indicates the amount of reservoir type storage, storage which is a function of outflow alone, that is assigned to the elementary reach.

Furthermore, because the error terms are linear additions to the continuity equation, the introduction of error related to the choice of  $\beta$  and  $\Delta T$  causes the model to behave in a manner similar to the introduction of reservoir storage in the channel. Thus, by the use of equation 5.14, the first order errors associated with  $\Delta T$  may be converted to equivalent reservoir storage.

This property has already been demonstrated in pictorial fashion in a previous chapter, where the shape of the outflow hydrograph for the first system was shown to be related to the peak of the hydrograph. (Figure 3.6) If the peak attenuated to a particular value, the flood wave appeared to have the same shape regardless of whether the attenuation resulted from the introduction of error in the order of  $\Delta X$  or error in the order of  $\Delta T$ .

A qualitative description of the equivalent reservoir storage in the channel may be obtained in the following manner. Dividing equation 5.12 by  $\frac{\Delta X}{2} \frac{\delta^2 Q}{\delta x^2}$  yields:

$$\frac{2\varepsilon}{\Delta X \frac{\delta^2 Q}{\delta x^2}} = S = (2\alpha - 1) + (1 - 2\beta) \frac{C\Delta T}{\Delta X}$$
(5.24)

where: S = stability number defined in Appendix A.

Using the dimensionless number, S, it is possible to describe the amount of equivalent reservoir storage in an elementary reach. If the nucleus is in the centre of the molecule, S = 0.0 and there is no reservoir storage in the channel. However, when the nucleus is located at the downstream boundary ( $\alpha = 0.0$ ,  $\beta = 0.5$ ) the value of S is -1.0 and the elementary unit is modelled as a reservoir.

The amount of apparent reservoir storage in the elementary reach can be increased even further by selecting a value of  $\beta$  greater than 0.5. The exact amount by which the apparent reservoir storage will increase depends upon the values of **C**,  $\Delta X$  and  $\Delta T$ . If  $C \frac{\Delta T}{\Delta X}$  is precisely equal to unity, the value of S will vary in direct proportion to both (2  $\alpha$  -1) and (1-2 $\beta$ ).

For a linear channel the kinematic wave velocity is constant with respect to y and there is no difficulty in determining the value of c to use in calculating S. In non-linear channel elements in which c is a function of the stage y, it is necessary to choose some value of stage or discharge for which c may be computed. Ideally this stage should be able to be selected from a knowledge of the inflow hydrograph. However, until some suitable guidelines are available to aid in this selection it is preferable to determine the appropriate value of c by simulation. Because the attenuation is a function of S it is possible to determine  $C \frac{\Delta T}{\Delta X}$  from the data presented in the tables:

For System 1,  $C \frac{\Delta T}{\Delta X}$  may be calculated in the following manner: (a) The time steps chosen for the simulation caused the peak of the inflow hydrograph to be truncated to 0.984 times the maximum full bank flow. Thus, the values chosen to determine c should be divided by 0.984.

(b) Using data from table 3.3, the sensitivity
to changes of *α* at the centre of the mole-

cule are:

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$$\frac{\Delta P}{\Delta(2\alpha-1)} = \frac{1.00 - \frac{0.924}{0.984}}{-0.50} = -0.122$$

The sensitivity to changes in  $\beta$  are:

$$\frac{\Delta P}{\Delta(1-2\beta)} = \frac{1.00 - \frac{0.925}{0.984}}{-0.50} = -0.120$$

Thus, S is slightly more sensitive to the value of  $\boldsymbol{\alpha}$  and

$$C\frac{\Delta T}{\Delta X} = \frac{-0.120}{-0.122}$$

or

$$C = \frac{-0.120}{-0.122} \times \frac{2500}{200} = 12.3 \text{ fps}$$

 (c) This type of calculation can be repeated for all of the data available and an average value determined. However, when applying this method it should be noted that attenuation may not, and indeed, usually will not, be a linear function of S. Thus, the sensitivity tests described above should be made using values of the attenuated peak outflow that are approximately equal.

If we replace the parameters  $\boldsymbol{\alpha}$  and  $\boldsymbol{\beta}$  by equivalent variables with the origin at the centre of the molecule we obtain

 $\theta = 2\alpha - 1; \phi = 1 - 2\beta$ 

Then the desired value of c may be expressed as

$$C = \frac{\Delta X}{\Delta T} \frac{\frac{\delta P}{\delta \phi}}{\frac{\delta P}{\delta \theta}}$$
(5.25)

If P is assumed to be 1.0 at the point  $\boldsymbol{\alpha} = 0.5$ ,  $\boldsymbol{\beta} = 0.5$ ( $\boldsymbol{\theta} = 0$ ,  $\boldsymbol{\phi} = 0$ ) then two simulation runs will provide the forward difference values for  $\frac{\delta P}{\delta \theta}$  and  $\frac{\delta P}{\delta \phi}$  from which c may be evaluated.

#### 5.2 "CASCADE EFFECT"

The "cascade effect" is defined as the extent to which overall attenuation is increased as the number of subreaches of constant  $\Delta X$ 

are increased. Dooge (1959) demonstrated that the peak outflow issuing from a series of linear reservoirs may be related to the Poisson (N-1, N-1). Evaluation of the Poisson function reveals that the term P(N-1, N-1) is proportional to  $\sqrt{N}$  where N is the number of elements chained together. A similar situation may be shown to exist for chains of non-linear reservoirs.

Several numerical experiments were performed to determine the relationship between the attenuation, (1-P), and the number of elements n when the constraint of linearity is removed. For these experiments, the cross section area from System 1 was utilized with reach lengths of 1, 250 and 5, 000 feet. The inflow hydrograph was defined as the triangular hydrograph used for System 1. The hydrograph was routed through various cascades with 2, 4, 8, 16, 32 and 64 units in each cascade.

In order to introduce a significant degree of attenuation and vary the molecule effect the computations were made with the following positions of the nucleus.

POSITION	ł	2	ŗ
a	0.0	0.0	0.5
β	ე.5	1.0	1.0

The results of these simulations are presented in figure 5.1. which shows a plot of attenuation (1-P), versus the number of cascaded reaches, N, in the system. Each of these numerical experiments resulted in a set of points that could be approximated by a straight line on a log-log plot. After fitting the straight lines to the data it could be seen that a family of nearly parallel curves resulted.

The slope of the various lines was determined to provide an estimate of the exponent, e, in the general equation

$$(1-P) = k N^{e}$$
 (5.26)

where:

P = peak outflow ratio
k = a coefficient
N = number of reaches in the simulation
e = an exponent

The results indicated a range of results from 0.542 to 0.574 for e as shown on figure 5.1.

These results are presented to demonstrate an interesting feature of the system and simulation. It is also interesting to note that if a linear system was utilized, the exponent would be of the order of 0.5. More tests are required to delineate the relation-ship of the exponent e to factors such as the input, and system



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characteristics. Further investigations into these items and the variance of the exponent, e, could provide information which may quantify the effect of system non-linearity on the response characteristics.

#### 5.3 "STORAGE EFFECT"

. "Storage effect" is defined as the attenuation that results from a finite quantity of live storage expressed as a proportion of the volume in the flood hydrograph. It is intuitive that more attenuation will occur as reservoir volume increases.

Since a constant inflow hydrograph was employed in all tests, this effect was studied by running further numerical experiments with System 1 in which the attenuation, 1-P, was dependent on the reach length  $\Delta X$ .

Following empirical evaluation of the molecule and cascade effect it was then intended to relate the attenuation as a function of the live storage. Although it may be anticipated that this relationship is highly non-linear over a wide range of attenuation values, it is reasonable to assume that within the relatively small range of

 $\Delta X$  values examined the function would not depart significantly from a straight line plot.

#### 5.4 ANALYSIS OF NUMERICAL EXPERIMENTS

For the purpose of evaluating the several functional relations described above, attention was directed to the result of numerical experiments on System 1 as recorded in tables 3.3 - 3.5. The various effects were isolated in a slightly different order to that in which they have been considered up to this point, i.e. cascade effect and molecule effect and storage effect.

#### 5.4.1 CASCADE EFFECT

In order to study the effect of increasing the number of serial elements in a system, the program was modified slightly so that  $\Delta X$ ,  $\alpha$  and  $\beta$  could be held constant, while n was increased geometrically. The results for various values of  $\Delta X$ ,  $\alpha$  and  $\beta$ are tabulated in table 5.5.

Following on the argument of section 5.2, a log-log plot of the attenuation (1-P) versus n (number of elements) revealed that the value of e in equation 5.26 was essentially constant with a mean value of 0.55. The slope of the various plots is indicated if figure 5.1 and ranges from 0.542 to 0.576.

The effect of this correlation is to allow the attenuation of a single reach to be determined as a function of (I - P)N<sup>0.55</sup>

TABLE 5.5

# TYPICAL RESULTS OF CASCADE EFFECT EXPERIMENTS

### SYSTEM 1

# $\Delta T = 200 \text{ SECONDS}$

#### **∆**X = 1250 FEET

	a	β	N	(1-P)
-			2	0.0164
	0.00	0.50	4	0.0247
			8	0.0381
			16	0.0525
	~		32	0.0755
			64	0.1089

where P is the resultant relative peak flow issuing from a cascade of N elements. Armed with this information it was now possible to extend the analysis to the molecule and storage effects.

#### 5.4.2 MOLECULE EFFECT

As discussed in section 5.1, the attenuation (1-P) may be related to the stability parameter S which is shown to be a function of  $\alpha$ ,  $\beta$ ,  $\Delta X$ ,  $\Delta T$  and c by equation 5.24.

Figure 5.2 shows a log-log plot of  $(1-P)/N^{.55}$  as a function of S. Each line represents values obtained for constant values of  $\Delta X$ . Once again the relations appear to be simple exponentials with the range of experiments covered. However, the exponent of the different experiments varies between 0.5 and 0.6. The average of these values is 0.55. The general result is that for constant values of N and  $\Delta X$ , the attenuation (1-P) is proportional to approximately-S<sup>0.55</sup>. Thus for constant N and  $\Delta X$ 

 $\frac{I-P}{-S^{0.55}} = CONSTANT$  (5.27)

#### 5.4.3 STORAGE EFFECT

Using the correlations determined in the foregoing two sections it is now possible to isolate the dependence of attenuation on

# FIGURE 5.2 ATTENUATION PER ELEMENTARY REACH Vs STABILITY NUMBER

 $\Delta T = 200$  SECONDS



the live storage. This was achieved by reducing the results of the numerical tests to obtain values of

$$\frac{(1-P)}{-S^{0.55}N^{0.55}} = f(\Delta X)$$
(5.28)

To express the relative live storage as a volumetric ratio, it is convenient to use as abscissa the non-dimensional term

$$\frac{\Delta X \cdot \Delta A}{Vol}$$

where: 
$$\Delta A = A(Qmax) - A(Qbase)$$

Vol = volume of inflow hydrograph above the base flow This choice of parameter is somewhat arbitrary and valid here only because of the constancy of the inflow hydrograph employed and the prismatic nature of the channel. The results of this analysis are shown in figure 5.3 in which it may be seen that the "reduced" attenuation is approximately a linear function of the relative live storage parameter.

$$\frac{(I-P)}{-S^{0.55}N^{0.55}} = k \frac{\Delta X \cdot \Delta A}{Vol}$$
(5.29)

where: k = 0.485 for the results examined



#### 5.4.4 EXTENSION TO OTHER SYSTEMS

The results of the foregoing analysis are of interest but are rather dependent on the assumptions and constraints of the particular system from which the numerical data were obtained. It must be recognized that a more generally applicable relation may require inclusion of the effect of the hydrograph shape (defined perhaps by higher statistical moments) and the degree to which the expression  $\Delta X \Delta A$ truly represents the live storage in the system.

It must be recognized that the functional relations developed above are valid only for the system from which numerical data were obtained and that application of these correlations to other systems is unwarranted until such time as further numerical tests can be carried out. Until such time, however, it is desirable to develop some simple guidelines for the development, use and calibration of kinematic routing models. The following section attempts to develop such a methodology on the basis of the analysis presented earlier and the application in various circumstances is illustrated.

# 5.5 GUIDELINES FOR CALIBRATION OF KINEMATIC ROUTING MODELS

Since it is recognized that no attempt has been made to quantify the dependance of attenuation on hydrograph shape and storage

effect, it follows that the engineer engaged in constructing a kinematic routing model must make use of trial numerical experiments as an aid to calibration. This being so, it is desirable to identify a compound parameter which combines the several separate effects described above and which will enable computational effort to be minimized in the process of model development and adjustment.

In the process of calculating the relation of attenuation to S,  $\Delta X$  and N, it was shown that a correlation existed which promised to become a useful tool to efficiently calibrate a numerical model.

Rearranging equation 5.14 indicates that

$$2\varepsilon = \frac{\delta^2 Q}{\delta x^2} \cdot S \cdot \Delta X$$
 (5.30)

The term  $\frac{\delta^2 Q}{\delta x^2}$  is a function of the system and the inflow hydrograph and remains fixed for any independent values of Q,

 $\beta$ ,  $\Delta X$  and  $\Delta T$ . Assuming that the attenuation is some function of the error term, a correlation between P and-S $\Delta X$  should exist. Figure 5.4 contains a plot of P versus-S $\Delta X$  for System 1; a similar plot for System 2 is shown in figure 5.5. The results plotted on these curves are shown in tables 5.6 and 5.7.







# TABLE 5.6

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# TYPICAL VALUES USED TO DEFINE

#### THE P VS SAX CURVE

# SYSTEM 1

# 40,000 FT. DOWNSTREAM

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ΔT (SEC.)	ΔX (FT.)	α	β	S	S Δ X (F T.)	PEAK
	2, 500	0.50	0,5	0.0	0	0 988
		0,25		-0,5	- 1250	0,924
		0.00		-1:0	-2500	0,890
	5, 000	0,50		0.0	0	0.965
200		0,25		-0.5	-2500	0,890
		0.00		- 1. 0	- 5000	0.840
		0.50		0.0	0	0,941
	10,000	0,25		-0.5	- 5000	0.840
		0,00		- 1. 0	- 10000	0,771

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# TABLE 5.7

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# TYPICAL VALUES USED TO DEFINE

# THE P VS SAX CURVE

#### SYSTEM 2

ΔT (HRS.)	ΔX (MI.)	α	β	S	S <b>∆</b> X (Ml.)	PEAK
2	10	0.00	0.25 0.50 0.75 1.00	-0.285 -1.000 -1.715 -2.430	- 2.85 -10.00 -17.15 -24.30	0.986 0.952 0.919 0.888
		0.25	0.50 0.75 1.00	-0.500 -1.215 -1.930	- 5.00 -12.15 -19.30	0.975 0.942 0.910
		0.50	0.50 0.75 1.00	0.000 -0.715 -1.430	0,00 - 7.15 -14.30	0.996 0.964 0.932
		0.75	0.75 1.00 1.00	-0.215 -0.930 -0.430	- 2.15 - 9.30 - 4.30	0.986 0.954 0.976
3	20	0.00	0.25 0.50 0.75 1.00	-0.463 -1.000 -1.537 -2.075	- 5.26 -20.00 -30.74 -41.50	0.956 0.904 0.860 0.825
		0.25	0.25 0.50 0.75 1.00	+0.037 -0.500 -1.037 -1.537	+ 0.74 -10.00 -20.74 -30.74	0.999 0.952 0.903 0.860
		0.50	0.50 0.75 1.00	0.000 -0.537 -1.075	00.00 -10.74 -21.50	0.995 0.948 0.901
		0.75	0.75	-0.037 -0.575	- 0.74 -11 5	0.990 0.944

# 300 MILES DOWNSTREAM

It seems reasonable to assume that similar curves may be obtained for any system. A procedure may therefore be developed whereby these curves may be obtained from two or more simulation runs with carefully selected parameters. This curve may then be used as a guideline for the selection of the required value of  $S \Delta X$  corresponding to the observed attenuation. A suggested procedure for the selection of parametric values is given below in the form of a check test.

- 1) Select an appropriate value of  $\Delta T$  to model the inflow hydrograph. If the channel is arbitrarily shaped it may be necessary to specify the  $\Delta X$  increments due to the constraints of geometric data.
- Obtain an estimate for the kinematic wave velocity, c, based on the system and the inflow hydrograph.
- 3) Assuming that the  $\Delta X$  steps can be varied, select a value of  $\Delta X \doteqdot C\Delta T$  (say) with the constraint that  $\frac{L}{\Delta X}$  is an integer. In the first simulation selecting  $\alpha = 0.0$ ,  $\beta = 1.0$  will provide the maximum amount of attenuation that is possible with the specified  $\Delta X$  and  $\Delta T$ In addition, the numerical calculations will be unconditionally stable.

- 4) The above steps would define S∆X and performing the necessary mathematical calculations would provide a result with an attenuated outflow hydrograph.
- 5) If the predicted attenuation is greater than the observed peak then adopt a smaller value of  $S\Delta X$ , say one half of the previous value. There are a number of ways in which  $S\Delta X$  may be varied by adjusting one of four variables  $\alpha$ ,  $\beta$ ,  $\Delta X$  or  $\Delta T$ . Unless the value of c is well established, it is suggested that  $\alpha$  and  $\beta$  be adjusted and  $\Delta X$  and  $\Delta T$  held fixed. This may provide more information which may be helpful in determining c.
- 6) Repeat the calculations with the reduced value of c to obtain a new estimate of the peak outflow and hence define another point on the relationship

#### $P = f(S\Delta X)$ (5.31)

from which it may be possible to approximate the  $S\Delta X$ required to provide the needed amount of attenuation.

7) If the first value of the attenuated hydrograph is less than the desired amount, it is necessary to increase  $S\Delta X$ . This can be done by increasing either  $\Delta X$  or  $\Delta T$ . If the channel has been defined with arbitrary cross sections it would be necessary to adjust  $\Delta T$ . Otherwise an adjustment of  $\Delta X$  may provide the most efficient solution.

8) Successive adjustments of  $S\Delta X$  would produce the relationship of P to  $S\Delta X$  and allow the user to calibrate the model.

There may be other variations on the previously described calibration process that may be employed to achieve the desired results. For example, if the value of c is unknown, the preliminary calculations may be performed with  $\beta = 0.5$  and thus eliminate the need to have a value for the kinematic wave velocity. Another numerical experiment with  $\alpha = 0.5$  and  $\beta > 0.5$  would yield the data required to determine c. In selecting values for the parameters  $\alpha$  and  $\beta$  and step sizes  $\Delta X$  and  $\Delta T$  the user should remember to follow stability criteria guidelines (see Chapter 3) and must also utilize sound engineering judgment. This will ensure that unreasonably large numerical errors and approximations are not introduced into the calculations.

# 5.6 DISCUSSION OF HYDROGRAPHS RESULTING FROM ATTENUATED KINEMATIC SOLUTIONS

The results obtained from the kinematic simulations are shown in figures 5.6 - 5 8 for System 1 and in figure 5.9 for System 2.







(HOURS) TIME



When the bed slope was very shallow, dynamic effects were significant and the kinematic solution was lagged relative to the dynamic solution as shown in figure 5.6. However, as the bed slope was increased and dynamic effects became less significant the dynamic solution and the attenuated kinematic solution were very close as demonstrated on figures 5.7 and 5.8. Figure 5.9 illustrates that with System 2 there was a very close correspondence of the results from the two methods of analysis. Further discussion of criteria relating to the applicability of kinematic routing is found in Chapter 4.

#### 5.7 CONCLUSIONS

Numerical experiments have been performed to study the way in which attenuation is modelled by the:

- 1) Molecule effect
- 2) Cascade effect
- 3) Storage effect

Results of using the general kinematic flood routing method to simulate the above effects reveal that in certain cases it is possible to accurately model unsteady open channel flow. In the past kinematic techniques have been subject to criticism due to the sensitivity of the solution to the size of  $\Delta X$  and  $\Delta T$ . However, investigating this sensitivity and determining how it affects the solution, provides valuable insight into a potentially useful numerical tool. Varying the size of  $\Delta X$  changes the molecule effect, depending on the selection of  $\alpha$  and  $\beta$ , it changes the storage effect and for a fixed total reach length it affects the cascade effect.

Similarly the selection of  $\Delta T$  is a significant parameter in determining the molecule effect if the nucleus is not located on the centre of the molecule.

In the course of the numerical experiments, it was noted that the peak outflow, if plotted over the finite difference molecule, forms a relatively smooth surface. A limited relationship has been developed which can be used to predict the value of the peak at any point over the molecule after calibration. With this relationship, which combines molecule, cascade and storage effects, it is possible to establish a kinematic model based on results of a dynamic model and/or actual recorded input and output of the prototype. Using S and  $\Delta X$  it is possible to select values of  $\alpha$ ,  $\beta$ ,  $\Delta X$  and  $\Delta T$ to produce a particular outflow hydrograph, provided the dynamic effects are not too dominant.

Further research is required to determine the influence of other factors such as the inflow hydrograph characteristics and

physical properties of the system on the outflow hydrograph. This may enable a user to reduce the number of numerical experiments required to calibrate a model and may make it possible to predict the outflow hydrographs after the system has been altered without subsequent recalibration.

#### CHAPTER 6

#### MODELLING ATTENUATION WITH AN IMAGINARY RESERVOIR

The previous chapter was devoted to a discussion of apparent attenuation which resulted from the inclusion of storage in an elementary reach. This storage was introduced by modifying the computational approach so that combination channel-reservoir units were modelled. This chapter is devoted to studying a system where the channel and reservoirs were modelled as two distinct phenomena.

The prime objectives of this chapter are; (i) to provide a picture of the basic idea of a channel and reservoir in series, (ii) to study the sensitivity of the solution to the location of the reservoir and (iii) to investigate differences resulting from simulations where the reservoirs have different degrees of nonlinearity.

The results of the study are summarized in the concluding remarks. As well, several suggestions are included which point out areas where future research of this topic may be beneficial and where the concept of a channel and imaginary reservoir in series may be useful.
### 6.1 THE PHYSICAL SYSTEM:

Using a channel and an imaginary reservoir in series to simulate a natural channel system is an often used hydrologic tool. This technique, commonly known as the lag and route method, utilizes two components in performing the numerical simulation. These are (i) a linear channel and (ii) a linear reservoir. A linear channel is defined as a channel in which the flow rate is a linear function of the cross-sectional area. Thus, the kinematic wave velocity  $\frac{dQ}{dA}$  is a constant for any specific cross-section independent of flow rate (or stage). Translation of the wave is modelled by this portion of the model. The linear reservoir, which produces the attenuation of the flood wave, is defined as a reservoir in which the storage is a linear function of the outflow.

There are several characteristics of linear systems that allows the simulation process to be simplified to a degree. Firstly, when linear components are utilized in the study the response of the system and therefore the solution are not sensitive to the input and it is possible to utilize the principle of superposition. Thus, it is not necessary to consider the order in which the components occur.

The present study enlarges upon the basic lag and route approach by allowing for the inclusion of non-linear channel and non-linear reservoir elements in the system. In particular, the following variables were considered:

1. The location of the reservoir.

2. The degree of non-linearity of the reservoir.

3. The magnitude of live storage associated with the reservoir.

In studying the sensitivity of the solution to the location of the reservoir, the limiting cases were investigated. The traditional lag and route method, with the reservoir located at the downstream end of the channel was used along with what is hereafter defined as the route and lag method, a system with the imaginary reservoir at the upstream limit of the channel.

System 1 was utilized for this portion of the study. The channel was modelled as eight elementary reaches, each 5000 feet long. The single imaginary reservoir was described using the following equation:

	$ST = KQ^W$	(6.1)
where:	ST = storage	
	Q = flow rate	
	K = a parameter	
	w = a parameter	

A schematic representation of the system is shown in figure 6.1.

When studying System 1 using the finite difference solutions of the momentum and continuity equations, attenuated peaks of the order of 0.7 and 0.8 were obtained. By trial and error, an approximate res-

# FIGURE 6. I SCHEMATIC OF CHANNEL - RESERVOIR SYSTEM



ervoir size was determined which would result in attenuation of approximately the same magnitude. This was used to study the sensitivity of the solution to the location of the reservoir.

### 6.2 THE POSITION OF THE RESERVOIR

Table 6.1 contains a set of results obtained from a series of simulations using the route and lag analysis (reservoir upstream) while table 6.2 contains a set of results for similar simulations using the lag and route (reservoir downstream) method of analysis.

For each system studied, nine reservoir types were examined using value of K and w as follows:

> K = 12000, 14000, 16000w = 0.8, 0.9, 1.0

For each system and reservoir type, the attenuated peak, P, and two measures of lag were recorded. The lag parameters used were defined as follows:

Tp = time to peak (hr)

Tc = time to centroid (hrs) of hydrograph above baseflow The values of P, Tp, Tc are given for each system and reservoir type in tables 6.1 and 6.2. As might be expected, the attenuation P varies inversely with K and w since each of these terms tend to increase the sensitivity of live storage to the discharge Q (see equation 6.1). The lag terms similarly increase with increase of storage

# TABLE 6.1

# DATA FROM RESERVOIR -- CHANNEL SIMULATION

# SYSTEM 1

# 40,000 FT DOWNSTREAM

# \*Truncated

ĸ	0.8	0.9	1.0
12000	P = 0.833	P = 0.632	P = 0.437
	Tp = 1.24	Tp = 1.69	Tp = 2.13
	Tc = 1.47	Tc = 1.91	Tc = 1.89*
	P = 0.808	P = 0.599	P = 0.412
14000	Tp = 1.30	Tp = 1.74	Tp = 2.19
	Tc = 1.54	Tc = 1.93	Tc = 1.87*
	P = 0.780	P = 0.570	P = 0.391
16000	Tp = 1.36	Tp = 1.80	Tp = 2.19
	Tc = 1.60	Tc = 1.96	Tc = 1.85*

P = Peak Flow/Full Bank Flow

Tp = Time to Peak (hrs)

Tc = Time to Centroid (hrs)

See figure 6.3 for definition of above terms.

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# TABLE 6.2

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# DATA FROM CHANNEL--RESERVOIR SIMULATION

### SYSTEM 1

# 40,000 FT DOWNSTREAM

# \*Truncated

ĸ	0.8	0.9	1.0
12000	P = 0.818 Tp = 1.13	P = 0.616 Tp = 1.58	P = 0.429 Tp = 2.02
	Tc = 1.42	Tc = 1.80	Tc = 1.80*
14000	P = 0.791	P = 0.584	P = 0.405
	Tp = 1.19	Tp = 1.63	Tp = 2.08
	Tc = 1.49	Tc = 1.78*	Tc = 0,80*
16000	P = 0.766	P = 0.556	P = 0.386
	Tp = 1.25	Tp = 1.69	Tp = 2.08
	Tc = 1.56	Tc = 1.74*	Tc = 0.10*

P = Peak Flow/Full Bank Flow

Tp = Time to Peak (hrs)

Tc = Time to Centroid (hrs)

See figure 6.3 for definition of above terms.

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(K and w increasing) although with large amounts of lag and attenuation, the outflow hydrograph is significantly truncated resulting in serious error in evaluation of Tc. These cases are marked in tables 6.1 and 6.2.

The initial observation is that whereas the value of P is sensitive to values of K and w, the sensitivity to the position of the reservoir is low. In general, values of P in table 6.2 are 2% lower than in table 6.1. The effect of moving the reservoir downstream is small and even this small difference is due not to any fundamental change in the performance of the model, but rather to numerical "softening" of the peak in the channel routing stage which is more pronounced when the unattenuated hydrograph is routed down the channel. Thus, for this particular system, the attenuation is quite insensitive to the location of the reservoir. This can be attributed to the fact that the modulating effect of the reservoir is dependent upon wave shape, among other things, and there is no significant change in the wave shape as it passes down the channel.

While there was no significant differences in the amount of attenuation when the reservoir location was varied from the upstream end of the channel to the downstream limit, there was an appreciable change in the shape of the outflow hydrograph as shown in figure 6.2. With a route and lag system, the peak outflow occurred later and the



TIME (HOURS)

rising limb of the hydrograph was steeper than the simulation of the lag and route configuration. This phenomena results from the variation of kinematic wave velocity with stage.

For the system studied, higher stages have a higher kinematic wave velocity than lower stages or flow rates, thus, the peak of the unattenuated wave would travel down the channel faster and arrive at the downstream end passing through the reservoir sooner than the wave which was attenuated by a reservoir and then flowed down the channel. This variation in kinematic wave velocity also causes waves to steepen; but with the reservoir at the downstream end, the steepening caused in the channel was smoothed out. However, the waves from the reservoir located at the upstream end of the channel were steepened by the passage down the channel.

A comparison of the kinematic solutions with the dynamic solutions indicates that there is a considerably longer lag in the outflow from the kinematic simulations. This can be attributed to the introduction of the imaginary reservoir and the fact that the kinematic wave velocity of the channel is less than the dynamic wave velocity.

Although these preliminary tests confirmed the general qualitative dependence on K and w, further tests were used to extend this knowledge. In these tests the route and lag configuration was used

to minimize the errors that were introduced into the kinematic channel routing calculations.

# 6.3 MEASUREMENT AND DEFINITION OF THE RESPONSE

# CHARACTERISTICS

In order to provide a quantitative base for the comparison of the results of the numerous simulations that were made during the process of studying the effects of varying K and w, the following items were quantified:

- 1. Peak outflow.
- 2. Centroidal lag of the outflow hydrograph.
- 3. Skewness of the outflow hydrograph.

The first two items were computed in the simulation program and obtained from the computer printouts. A triangular inflow hydrograph, defined earlier in the description of System 1 was used in this study. As a result of the discretization of the hydrograph to conform with the time steps used in the numerical calculations, the peak of the hydrograph was truncated. Therefore, the inflow hydrograph was trapezoidal in shape with a peak value of 0.984 times the peak of the triangular hydrograph. However, this truncation was ignored and no corrections were made to the data produced by the computer. The centroidal lag value calculated by the computer included all of the base flow. Corrections were made to the data generated by the computer to obtain the centroid of the outflow hydrograph alone. A graphical representation of these and other values used to describe the hydrograph is contained in figure 6.3.

The third parameter used to describe the hydrograph was an attempt to quantify the general shape of the hydrograph in relation to the centroid and the peak discharge. This empirical skew factor was defined as:

SF = 
$$\frac{Tc - T_p}{T_b(P - 0.2)}$$
 (6.2)

where:

SF = skew factor

P = peak flow ratio
Tc = time of centroid
Tp = time of peak discharge
Tb = time base of inflow hydrograph.
0.2= base flow rate ratio

The skew factor describes the slope of the line from the base flow rate and the centroid of the hydrograph to the peak outflow and time of peak outflow as shown in figure 6.3.

## 6.4 SYSTEM PARAMETERS.

A measure of the system parameters (independent variables) was required to complement the response variables (dependent



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variables) so that the effects of non-linearity and the size of the reservoir could be studied.

The channel properties and the configuration of the system were held constant. The system was characterized by the ST versus Q curve and the two parameters that were utilized to describe this curve were the slope of the chord and the "bow" or departure from linearity in the curve over the range of flow rate that was experienced by the system. These parameters are respectively:

CHORD SLOPE = 
$$\frac{\Delta ST}{\Delta Q}$$

where: ST = storage

Q = flow rate

and

$$NL = \frac{\frac{dST}{dQ}}{\frac{\Delta ST}{\Delta Q}}$$
(6.4)

where:

$$\frac{dST}{dQ} = wKQ^{W-1}$$

A definition sketch of these parameters is included in figure 6.4.

(6.3)

# FIGURE 6.4 DEFINITION OF SYSTEM PARAMETERS



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Because the range of flows that were encountered was constant for all of the tests NL reduced to a function of w, however, if the range of flows varied, NL would reflect the change in the curvature of the ST versus Q curve for different ranges of flows.

## 6.5 DISCUSSION OF THE RESULTS

To determine the relationship of the response characteristics to the system parameters, a number of plots were made to show the relationship of P, Tc and SF to  $\frac{\Delta ST}{\Delta Q}$  and NL. In each of the cases,  $\frac{\Delta ST}{\Delta Q}$  was found to be the dominant physical parameter. The solutions were sensitive to the non-linearity of the reservoir as will be shown in the discussion of the various graphs.

# 6.5.1 RELATION OF PEAK FLOW RATIO, P, TO SYSTEM

## PARAMETERS.

Figure 6.5 contains a plot of P versus  $\frac{\Delta ST}{\Delta Q}$ . Each of the co-ordinate pairs plotted was identified to show the NL number for the particular simulation from which the value was obtained. This figure indicates that as  $\frac{\Delta ST}{\Delta Q}$  increases, P decreases. In other words, as the surface area of the reservoir increases, there is a larger amount of attenuation.

There was a small amount of scatter in the points that were plotted. By identifying each of the points according to the non-





linearity of the reservoir, it was obvious that this scatter could be correlated to NL.

As the non-linearity number, NL, increased algebraically from negative values to positive values, a larger  $\frac{\Delta ST}{\Delta O}$ was required to yield a given amount of attenuation. It should be noted that  $\frac{\Delta ST}{\Delta Q}$  was calculated on the basis of base flow and the maximum flow rate of the input, which in this study was full bank flow. If  $\frac{\Delta ST}{\Delta O}$  had been calculated using base flow and the peak of the attenuated outflow hydrograph, the variation of the values of  $\frac{\Delta ST}{\Delta Q}$ with respect to NL may have been reduced.

# 6.5.2 RELATION OF TIME OF CENTROID, Tc, TO SYSTEM

 $\frac{PARAMETERS}{\Delta Q}$  Shown in figure 6.6 indicates that as  $\frac{\Delta ST}{12}$  increases there is a corresponding increase in the centroidal lag. This is a result which is intuitively correct for as the size of a reservoir increases, it is reasonable to expect an increase in the delay caused by storage. The plot showed a small amount of scatter that again could be attributed to the non-linearity of the reservoir. With an algebraic increase of NL, there was a decrease in the cetroidal lag associated with a particular value of  $\Delta ST$ This can be attributed to the properties of a storage unit which has a proportionally larger storage with high flow rates when compared





with lower discharges. During the early portions of the inflow hydrograph, when the flows are low, there is no delay in the flow through the storage element. This contributes to the earlier centroid of the outflow hydrograph.

# 6.5.3 RELATION OF SKEW FACTOR, SF, TO SYSTEM

 $\frac{\text{PARAMETERS}}{\text{The plot of SF versus}} \frac{\Delta ST}{\Delta 0} \text{ is shown in figure 6.7. The}$ two previous graphs have shown relationships that were very nearly single valued. However, there is a large amount of scatter that can be attributed to the non-linearity of the reservoir and to what appears to be numerical errors in the computation of the skewness factor. This numerical error is a result of truncation in the computer output. From the data available, it appears that as the nonlinearity increases, the skewness decreases. An increase of  $\frac{\Delta ST}{\Delta O}$ causes a larger skewness. It will be shown later that the time of the peak outflow is a unique function of P. Thus, for a given P value, the time of the peak will be fixed, but the centroid of the hydrograph will vary in relation to NL. As NL increases and Tc decreases, SF will decrease.

The curves that are presented in the previous discussion are only valid over a finite range. This is due to the fact that as





increased, there was an increase in the length of the falling limb of the outflow hydrograph, which in turn resulted in a truncation of a portion of the hydrograph due to the termination of the simulation. The curve showing P versus  $\frac{\Delta ST}{\Delta Q}$  is not subject to errors of this nature, but the plots of the centroidal lag and the skewness factor are influenced by this error when  $\frac{\Delta ST}{\Delta Q}$  is larger than approximately 4,000 seconds.

Reductions in the amount of scatter in the figures which show peak outflow, centroidal lag and skewness factor as a function of  $\frac{\Delta ST}{\Delta Q}$  may be obtained if the chord slope is computed using the base flow and the peak outflow rather than base flow and peak inflow. However, the use of the latter technique may remain the most useful due to the fact that peak outflow is an unknown until the simulation is completed.

# 6.6 FURTHER COMPARISONS

To provide a comparison of the shapes of the variance of the hydrograph shape as the parameters K and w were varied, two of the many hydrographs obtained from these simulations were plotted on figure 6.8. Each of these hydrographs had a peak ratio of approximately 0.83. With K = 10 and w = 1.5, the hydrograph rose slightly quicker and dropped slightly sooner than the simulated hydrograph that was obtained using K = 12,000 and



w = 0.8. In the region of the peak there was no significant difference between the hydrographs and the two plots coincide.

The reason for the very close correspondence of the two floodwaves in the region of the peak can be attributed to the characteristics of the system being studied. For a given inflow hydrograph, kinematic channel system and imaginary reservoir the time at which a particular peak outflow occurs is constant regardless of the value K and w chosen to simulate the reservoir. For example, with the reservoir located at the upstream end of the channel, the hydrograph will be attenuated and the peak will occur at the intersection of the outflow hydrograph and the falling limb of the inflow hydrograph, thus uniquely determining the time of peak outflow from the reservoir independent of the reservoir characteristics, though the reservoir properties determine the peak. Similarly, the travel time for a specific flow rate is constant for a particular kinematic channel. A plot of the peak outflow ratio was plotted along with the times that these peaks occur. This data is contained in figure 6.9. As the peak outflow increased the speed at which the wave propagated increased and thus the time lag decreased.

Scatter resulted from the discrete representation of the hydrograph which caused the occurences to be represented at intervals one time step apart.

# FIGURE 6.9 PEAK OUTFLOW VS Tp

RESERVOIR - CHANNEL, SYSTEM I



Figure 6. 10 shows three storage versus flow rate curves for reservoirs which attenuated the flood wave to a peak ratio of 0.7 orll, 672 cfs. It can be seen that the maximum amount of storage utilized varied significantly as the non-linearity of the reservoir characteristics varied. With K = 15 and w = 1.544, the maximum storage required was twenty seven million cubic feet. Forty two million cubic feet of storage was required with K = 16,000 and w = 0.833. This fact may not be particularly important where the reservoir is imaginary, however, if the reservoir was being constructed to provide maximum attenuation, it would be prudent to design and operate the control structures so that the storage-outflow relationship is defined by an equation with a high index, w. In essence, this means that the reservoir should not be filled during the lower flows of the rising limb of the hydrograph, and the available storage should be reserved for larger flows immediately before and after the peak inflow.

#### 6.7 CONCLUSIONS

This chapter has provided a brief review of lag-route and route-lag methods of simulated unsteady flow in open channels. It may be of interest to note that the use of a reservoir in series with a kinematic channel is analogous to the continuously stirred tank





FLOW RATE (THOUSAND CFS)

reactor and the plug flow tank reactor utilized in chemical engineering simulations. With reservoir and channel systems the objective is to predict flow rates while the concentration of a tracer is the important variable with the tank reactor systems.

The basic properties of the reservoir channel system that were studied included;

- 1. The location of the reservoir.
- 2. The degree of non-linearity of the reservoir.
- 3. The magnitude of live storage associated with the reservoir.

Reservoir characteristics were described using the following equation:

	$ST = KQ^{m}$	(6.1)
where:	ST = storage	
	Q = flow rate	
	K = a parameter	
	w = a parameter	

The general kinematic flood routing method was utilized to perform the channel routing computations.

The results of the study may be summarized as follows:

1. The general kinematic flood routing technique seems to perform more satisfactorily when the flood wave is less peaked.

- 2. For the particular system studied, the position of the imaginary reservoir affects the shape and timing of the hydrograph. Moving the reservoir toward the downstream boundary results in an outflow hydrograph that occurs sooner and which rises more slowly. The amount of attenuation predicted by the simulations indicates that the reduction of the flood peak is relatively insensitive to the location of the reservoir.
- 3. The non-linearity of the reservoir affects the peak of the outflow hydrograph, the centroidal lag of the hydrograph and the skewness factor. However, the chord slope of the relationship between storage and discharge,  $\frac{\Delta ST}{\Delta Q}$ , appears to be the dominant factor in determining values of these response characteristics. A comparison of several hydrographs, which had approximately the same peak outflow ratio indicated a relatively small change in the outflow hydrograph as a result of varying the parameters K and w. This lends substance to the statement of Dooge and Harley (1967) that the assumption of a linear reservoir for simulation purposes does not appear to be particularly restrictive.

The simulation of System 1 by means of the dynamic 4. or "complete" solution with So = 0.0002 and n = 0.0149yields an outflow hydrograph which is attenuated to a peak value approximately equal to 0.75 times full bank flow. Comparisons with the results of the route and lag method indicate that some thirty to forty million cubic feet of storage would be required to produce an attenuation equal to that provided by dynamic dispersion. If the construc-' tion of a reservoir on a channel similar to System 1 was being contemplated, the attenuation of the flood wave as it passes down the river must be considered. Ignoring this attenuation would result in the dam being larger than required to produce a given reduction of the flood wave. However, if the channel attenuation is accounted for, it is possible that the size of the dam could be reduced knowing that the flood wave would be attenuated to the required amount before flowing out of the system.

The comparison of the route-lag and lag-route results with the data obtained from the dynamic simulation indicates that the lag predicted by the complete solution is significantly shorter than the lag predicted by either of the

other two methods. This can be attributed to the lag that was introduced by the inclusion of the imaginary reservoir and the fact that the kinematic wave velocity is slower than the dynamic wave velocity. The routelag model has the longest lag of the systems studied.

More study is required to provide guidelines for selecting the appropriate parameters to use in describing the imaginary reservoir for a particular physical system. However, it appears that utilizing a linear reservoir and calibrating the simulation tool with recorded or simulated data is a viable and extremely useful tool for hydrologic simulation.

# CHAPTER 7

# A KINEMATIC FLOOD ROUTING MODEL

Comparisons of kinematic flood routing simulations with rigorous solutions of the momentum and continuity equations describing unsteady flow phenomena have revealed that unsteady flow problems may be classed in three different categories. These are:

- Situations where kinematic theory provides as accurate an answer as could be obtained by the use of the more expensive rigorous solution.
- 2. Cases where a good approximation is obtained by using the kinematic theory in conjunction with another approximation such as an imaginary reservoir in series with the channel to account for attenuation.
- 3. Problems where the changes in flow are so rapid or the slope so shallow that the kinematic wave theory provides answers which are markedly different from those observed or predicted by a dynamic solution.

At present, criteria for differentiating between the various classes of problems are more qualitative than quantitative and these

criteria require further study. Development of an efficient and versatile computer program to perform kinematic flood routing calculations can be justified on the grounds that it can be used as a research tool to help establish quantitative guidelines for determining into which of the above mentioned categories a given physical problem falls. Further comparisons between kinematic wave solutions and results obtained by solving the momentum and continuity equations could provide the key to development of the rules for successfully applying kinematic wave methods. The use of the program as an engineering tool must not be overlooked and, to this end the program should be capable of handling natural channels defined by arbitrary geometry.

This chapter deals with the objectives and development of an interactive computer program capable of performing kinematic flood routing calculations for systems of arbitrary geometry. After outlining the objectives of the program and elementary operations used to route a flood, several applications of the program are shown. Concluding remarks are provided, which outline specific aspects of the method and the computer program that could be developed further.

#### 7.1 OBJECTIVES

The primary objective of this aspect of the study was to develop a computer program capable of being used as an engineering

or research tool. To accomplish this goal, the following specific program capabilities were identified.

- Kinematic routing methods must be able to be applied to natural channels of arbitrary geometry. The general kinematic method outlined in Chapter 3 was used as the theoretical basis and the numerical calculations were performed using a direct finite difference technique.
- 2. It has been shown in Chapters 3 and 5 that the values selected for the two parameters *α* and *β* have a very noticeable effect on the results of the computer simulation. To enable the use of this program as a research aid, it is vital that there be flexibility in the selection of *α* and

 $\boldsymbol{\beta}$ . Also, changing the two parameters may be beneficial in engineering applications.

- 3. To predict the attenuation of the flood wave, it may be necessary to insert a reservoir in series with the channel. This storage unit may be included either by variation of the parameters in the kinematic routing algorithm, as outlined above, or by placing an imaginary reservoir in series with the channel.
- 4. A command oriented procedure is desirable to enable the user to have maximum flexibility in the manner in which the particular system is analyzed.

5. The program should be designed to take advantage of the facilities of time sharing since this mode of operation allows the maximum amount of flexibility and interaction with the numerical representation of the system as the flood is being routed along the channel.

#### 7.2 DEVELOPMENT OF PROGRAM PROCEDURE

The development of a program that fulfils the above objectives was carried out within the context of the Civil Engineering Program Library (CEPL), Smith (1970), Smith (1974), Walden (1973). This allowed the existing resources of the program library to be utilized and, hopefully, will ensure that the finished program is accessible for use by others.

Currently, the CEPL is comprised of FORTRAN routines capable of analyzing a number of elementary engineering programs, and designed to perform a number of basic operations. These programs may be classified into three general categories.

- Basic routines provided to perform, one elementary function or calculation. A subroutine, which determines the cross-section properties, falls into this category.
- Slightly more comprehensive routines which require calls of the lower level subroutines to perform an elementary operation fall into the second group.

3. Managerial routines which utilize the other routines in the process of solving a particular problem.

A "building block" approach may be taken to develop a new program utilizing small modules. In this way, documentation is simplified and users can easily apply the routines.

One of the features of the CEPL, which is particularly useful in achieving the objectives outlined earlier, is the method of defining the arbitrary geometry of a natural channel. These properties are specified by section number, chainage from the upstream end and a roughness measure. The cross-section data is entered by specifying a series of coordinates which define the outline of the channel. Figure 7.1 contains a pictoral representation of the manner in which the cross-section is defined.

The computation of steady state profiles is a vital step in the process of routing a flood using the kinematic method outlined in Chapter 3.

A number of routines are available in the CEPL to calculate these flow profiles. For example, subroutine EZRA performs the backwater calculation for each reach between two cross-sections. However, subroutines CRITIC and CONTRO are available to calculate the flow depth that may occur at a channel transition such as a bridge or weir.

# FIGURE 7.1

# DEFINITION OF THE CHANNEL CROSS-SECTION




New subroutines, known as KINRUT and RESVOR, designed to be consistent with CEPL, were written to perform the kinematic flood routing and the reservoir routing. Documentation of these subroutines is provided in Appendix D. KINRUT routes the flood through one elementary reach using the functional relationships which are defined prior to calling the subroutine. This allows flexibility in the selection of values for the parameters

 $\alpha$  and  $\beta$ . The other method of including a pseudo-reservoir into the simulation is facilitated by RESVOR. Storage in the imaginary reservoir is defined using the following equation:

ST = KQ<sup>W</sup>

where: ST = Storage

Q = Flow rate

K = A parameter

w = A parameter

#### 7.3 THE PROGRAM

The program developed to analyze the passage of a flood wave in a channel using kinematic wave theory is predominently a subroutine of the "managerial" type. A small driver program serves

(7.1)

merely to dimension the arrays required for the calling statement. The chief function of the managerial subroutine is to provide a forum for its interactive operation of the various elementary computational subroutines through the use of operational data entered from a time-shared console.

The flow diagram in figure 7.2 shows the general operation of this routine. After the execution of the program begins, the required storage space is dimensioned in the main program. Subsequently, subroutine RIVER3 is called and the geometric data is read in from a previously defined tape. Further operation of the program requires explicit direction by commands entered in the operational data.

The operational data includes commands which direct the program to allow entry of data describing the downstream water levels, flow rates along the channel and the resistance law to use for computation of the steady flow state profiles, among other things. Data describing the system geometry may be printed and this information may be adjusted, if desired. The extent of the options available is best described by the list of commands available in RIVER3 shown in figure 7.3.

#### 7.4 USE OF THE PROGRAM

Before presenting an example of the use of the program, a brief discussion of the steps required to perform the kinematic

### FIGURE 7.2

### PROGRAM FLOW CHART



# FIGURE 7.3

# LIST OF COMMANDS

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AFTER INVITATION TO TYPE ":-" GIVE ONE OF
THE FOLLOWING COMMANDS
DISCHARGE TO SPECIFY FLOW
D/S WLTO DEFINE DOWNSTREAM CONTROL LEVEL
INFLOWTO DEFINE INFLOW HYDROGRAPH
RESISTANCE TO SET FLOW RESISTANCE LAW
OLD SECTION TO PRINT COORDS OF A SECTION
NEW SECTION TO REDEFINE COORDS OF A SECTION
OLD COEFFTO PRINT ROUGHNESS MEASURE
NEW COEFFTO REDEFINE ROUGHNESS MEASURE
CRITIC TO COMPUTE CRITICAL DEPTH AT A SECTION
CHANGES TO PRINT CHANGES OF COORDS OR ROUGHNESS
TABLE TO PRINT TABLE OF ALL SECTIONS DATA
PROFILES TO PRINT OUT SURFACE PROFILES
COMPUTE TO COMPUTE PROFILES
ROUTE TO ROUTE THE FLOOD
RESERVOIR TO ROUTE THROUGH A RESERVOIR
RESTART TO BEGIN AGAIN
HELP FOR COMMAND OPTIONS
STOP TO TERMINATE

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routing computations will be discussed. To calculate the functional relationships required for the flood routing, it is necessary to perform a steady state analysis of the watercourse and to calibrate the model.

The data required, and the order of operations is generally as follows:

- Define the downstream control by specifying the stagedischarge curve using a series of coordinates.
- Define the lowest and highest flow rates to be used for computing the flow profiles and the number of profiles to be computed. These flow rates may be varied along the channel to account for lateral inflow or outflow.
- 3. Define the resistance law to be used in the computations.
- 4. After entry of the above data, the profiles may be computed.
- 5. If a printout of the surface profiles is desired, it may be obtained using the PROFILES command.

At this point, the user is provided with data describing the steady state performance of the system and field observations, if available, may be compared with the profiles. If the two sets of information show significant discrepancies, physical data, such as the roughness coefficients, may be varied and the profiles recalculated. This process could then be repeated if necessary until the computed results agree with the observed profiles. Alternately, the user may wish to use only the steady state capabilities of the program and could study the effects of channel modifications by varying the cross-section data or could discontinue the execution of the program.

If the steady-state calculations are found to be satisfactory after completing an examination and/or calibration process, the unsteady analysis may be started. This involves the separate operations of (i) defining the inflow hydrograph and (ii) performing the required flood routing.

Currently the program is capable of handling only one inflow hydrograph; but this hydrograph may be defined at any section specified by the user in the operational data.

### 7.4.1 Definition of Inflow Hydrograph

The inflow hydrograph is described by time and flow rate coordinates. Straight lines are assumed between these points. The flow rate for time less than the first coordinate is assumed to be the flow rate of the first coordinate and the flow rate for time greater than the time of the last coordinate is assumed to be the flow rate of the last coordinate.

### 7.4.2 Performing the Routing

Kinematic or reservoir routing may be performed in any order. The only limitation, at present, is that the routing must start at the location where the inflow hydrograph is defined or where the last routing ended. This allows a user to route a flood for several reaches, insert a reservoir, if desired, and/or continue with the kinematic routing. Alternately, the process may be restarted at the location where the inflow hydrograph was defined.

If the kinematic routing through the channel is used, an option allows the user to redefine the parameters  $\mathbf{a}$  and  $\mathbf{\beta}$  for a number of elementary reaches. The default value for these two parameters is 0.5. Other information defining the time step, sections where the routing is to begin and end, and the start and finish of the time period under consideration must also be entered. The reservoir routing option requires the same type of data regarding time step, as well as start and finish time, but the size of the reservoir, as defined by the parameters K and w, and the location of the reservoir must be specified.

Using these features, a calibration of the response hydrograph may be performed or sensitivity tests may be performed. In general, the other command options may be used at any time to provide information or to modify data. Checks are built into the options to warn the user if other data must be entered before the specified operation may be carried out or if changes have been made, say to cross-seclion data, and vital information such as surface profiles, has not been

recalculated. The exception to this rule is the STOP command. This command may be used only at the end of the computations. No other commands may be entered after STOP.

#### 7.5 EXAMPLE APPLICATIONS

This section of the chapter is devoted to demonstrating the use of RIVER3 in providing an effective means of studying kinematic flood waves in natural channels for research or engineering purposes. Two systems and the results from each are presented. Each description outlines the physical system and results from the various numerical experiments.

#### 7.5.1 Application One

The first system studied consisted of approximately 10 miles of channel defined by 58 cross-sections. The profile of the invert level along the watercourse is shown in figure 7.4 along with some of the typical cross-sections. Appendix G contains a listing of the geometric data used to describe the channel.

An inflow hydrograph with a trapezoidal shape having a base flow of 350 cfs and a peak of 4,000 cfs was utilized. Figure 7.5 shows this hydrograph and the outflow hydrograph obtained using an implicit dynamic analysis. Walden (1973).



### FIGURE 7.5 HYDROGRAPHS FOR APPLICATION ONE



This particular physical system showed very little attenuation, which indicates that a kinematic solution may prove to be a very good approximation to the more complete dynamic solution. The first tests performed were used to study the effects of changing the size of the time step. The values chosen for the two tests were 1, 200 seconds and 300 seconds. The outflow hydrographs predicted from these simulations are shown on figure 7.5 along with the solution predicted using the dynamic analysis. Both of the kinematic solutions predicted the general shape of the outflow hydrograph. However, there were signs of slight instability on the falling limb of the hydrograph with  $\Delta T=1$ , 200 seconds. Neither kinematic solution correctly predicted the peak outflow nor did they simulate the dynamic effects demonstrated by the earlier rise of the outflow hydrograph.

To simulate the attenuation of the floodwave, several tests were performed with the reaches between section numbers 41 and 45 simulated using:

$$a = 0.0$$
  
 $b = 0.5$ 

The other reaches were simulated using the standard default values of 0.5 for these parameters. The results of these tests, which did not successfully predict the attenuation of the flood wave as it passed

through the channel are shown in figure 7.6. This may have been due, in part, to slight instability of the computation in the region of the peak. Nevertheless, the general shape of the outflow hydrograph was modelled.

A number of other tests were performed using various values for the parameters and time steps, etc. In general, these tests did not successfully model large amounts of attenuation. This was not a severe restriction in application one due to the fact that the physical prototype did not appear to manifest large amounts of attenuation. The introduction of a significant amount of attenuation, if that was deemed necessary, could be made by the inclusion of an imaginary reservoir in series with the channel. The simulation shown in application one was relatively successful in simulating the shape of the outflow hydrograph. This was due largely to the fact that the system was relatively steep and a kinematic analysis was a good approximation. The useful ness of the kinematic analysis as an engineering tool relies upon the verification of the results. In the present study, this verification was provided by a more rigorous dynamic analysis. Other means of substantiating the results could be (i) verification by comparison of the computed results with hydrologic data recorded in the system or (ii) development of guide-

### FIGURE 7.6 HYDROGRAPHS FOR APPLICATION ONE



lines that would enable a user to determine the applicability of kinematic analysis to the particular system without the use of the dynamic analysis or extensive hydrologic data. A preliminary basis for these guidelines is presented in Chapter 4.

#### 7.5.2 Application Two

The second prototype that was modelled using RIVER3 was a natural channel that had two major constrictions in the lower reaches. Figure 7.7 contains a profile of the invert elevations of application two and some typical cross sections of this waterway, the geometric data is listed in Appendix G. These constrictions were large embankments with relatively small culverts extending through the lower portions. During a major flood, substantial ponding of water would occur behind these embankments, particularly the upstream embankment. This ponding would result in the attenuation of a flood wave if the embankment was stable under the severe load imposed by the innundation of the upstream side. The system described above exists, in the Lower Ancaster Creek near Hamilton, Ontario. The upstream embankment carries a railroad track and a highway traverses the downstream embankment,

In this study, several modifications were made to the data to make the problem more tractable. Firstly, during a major flood, the culverts would be flowing full, the assumption of open channel flow





is not valid. Thus, instead of culverts through embankments, the structures were modelled as embankments with deep, narrow slots, through which the water flowed. Secondly, Sulphur Creek, a major tributary, joins the main stream above the first large constriction.

The program was designed to handle a single water course as opposed to a network. The complication arising from the existence of Sulphur Creek tributary was avoided by ignoring the tributary inflow and modifying the storm hydrograph accordingly.

It should be noted that these difficulties were a result of restrictions in the CEPL routines designed to compute the steady state profiles rather than limitations of the flood routing algorithm. Further development of the computer program could lead to a solution of these problems without the above mentioned simplifications.

The primary purpose of application two was to demonstrate the use of RIVER3 in a situation where a significant amount of attenuation would be manifested.

The input hydrograph, shown in figure 7.8 was routed from section number 37 to the first embankment at section 51 and thence to section 64, the downstream limit of the stream.

The first routing, using the nucleus in the centre of the finite difference molecule proved to be unstable. To provide a solution which was stable, the nucleus of the molecule was moved downstream

## FIGURE 7.8 HYDROGRAPHS FOR APPLICATION TWO

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(i.e. maintaining  $\beta = 0.5$ ). The results of these simulations are shown in figure 7.8.

With  $\mathbf{a}$  set to 0.0, the outflow hydrograph rose slowly and peaked at a value of approximately 1,850 cfs. There were no signs of numerical instability. Using a value of  $\mathbf{a} = 0.2$ , a similar soluation was obtained, except there was a sharp rise to a peak of about 2,200 cfs., and then a drop to approximately the same hydrograph predicted by the solution with  $\mathbf{a} = 0.0$ . The last solution shown on figure 7.8 was obtained using  $\mathbf{a} = 0.4$ . In the early portions of the outflow hydrograph, there is no significant differences between it and the other two solutions. However, in the region of the peak there are indications of numerical instability. The tests results shownon figure 7.8 were compiled using a time step of one half an hour. To determine the sensitivity of the numerical analysis to changes in the size of the time step, two tests were performed with

 $\mathbf{\alpha} = 0.5$  and  $\mathbf{\beta} = 1.0$  These results are shown in figure 7.9 It can be seen that with both of the simulations, a relatively long, well rounded hydrograph resulted and the peak outflow was attenuated to approximately one-half of the peak inflow. Using  $\mathbf{\Delta}T = 1$  hour resulted in the flood wave being attenuated to a peak of 1, 420 cfs. while the peak obtained using  $\mathbf{\Delta}T = \frac{1}{2}$  hr. was 1,530 cfs.

## FIGURE 7.9 HYDROGRAPHS FOR APPLICATION TWO



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Figure 7.10 contains a comparison of several outflow hydrographs at the downstream limit of the stream. When the nucleus was located in the centre of the molecule for all of the reaches in the simulation, some numerical instability appeared after the peak. The same simulation was performed using  $\Delta T = \frac{1}{2}$  hour, and the solution showed even greater numerical errors, therefore, the results were not plotted. In general, numerical stability problems were reduced when the simulation was performed with the nucleus shifted from the centre of the molecule for at least a few of the reaches upstream of the first large embankment.

The results of these simulations indicate that it is possible to model a system where the flood wave is attenuated as it passes through the system. However, it is necessary that the results be verified by a more rigorous analysis.

An attempt to provide the data to verify the kinematic model using an implicit dynamic technique was unsuccessful. Amein (1969), Walden (1973). The computer program utilized a four point implicit method to approximate the differential equations governing unsteady flow in open channels. The system of non-linear equations that resulted from the finite difference equations was solved using a Newton-Raphson iteration technique. However, for this application, the interative procedure failed to converge to a solution and execution

## FIGURE 7.10 HYDROGRAPHS FOR APPLICATION TWO



was abnormally halted. It appears that alternate methods of solving the system of equations must be employed in order to obtain a solution.

Application two may be handled by an alternate approach which involves assuming a reservoir in the reaches above the first embankment. The properties of the reservoir would be determined by the volume of water stored in the valley as a function of outflow. This type of approach could possibly yield a more valid solution than the previous method which considered the large volume in the valley as a series of channel reservoirs. If the valley section was simulated as one unit, more attenuation may result and the time of travel through the valley would probably be reduced. This remains an area that requires further study.

### 7.6 CONCLUSIONS AND DISCUSSION OF RESULTS

The computer program has been developed to allow a user to route a flood through an open channel system defined by arbitrary geometry. The routine is designed to be utilized in time sharing mode with operational data entered on a teletype console. The results are printed on the same unit. Execution of the program is directed by commands entered as operational data. The general capabilities of the program include the ability to perform steady state analysis

of a system and store the data for use during the analysis of unsteady flow conditions. Additional commands are included to provide efficient data handling capabilities and to allow modification of the geometric data during the execution of the program.

Several applications of the program have been presented and have yielded results which indicate that the kinematic analysis is useful in analyzing physical systems provided the solutions can be verified by either (i) a more rigorous analysis or by (ii) recorded hydrologic data. Attenuation may be simulated by varying the values of  $\boldsymbol{\alpha}$  and  $\boldsymbol{\beta}$  or by the introduction of an imaginary reservoir in series with the channel. Commands incorporated in RIVER3 facilitate both of these devices.

Stability was one of the prime factors which governed the success of the various numerical experiments that were performed to demonstrate the two applications. Numerical instability may be reduced by positioning the nucleus away from the centre of the molecule in a more stable region. However, another approach may be utilized to improve stability. Because the program analyzes a complete time history of one reach before computing the conditions in the next portion of the watercourse, it is possible to use a time step that varies as a function of the reach length and the kinematic wave velocity through that section of the stream. The output hydrograph would

be reduced to a common time step via an interpolation routine. This requires further study and would involve some rearrangement of the routines currently being employed in RIVER3. Some experimentation with the selection of the value of  $\Delta T$  may result in a solution which is stable and yet avoids the need to have a variable time step. Alternately, the flood may be routed through one series of reaches using one time step and then routed through another series using a different increment of time. One unique feature of kinematic routing is the fact that, in some cases, a large time step may produce more stable results.

Compatibility with the CEPL has been one of the major guidelines utilized in the development of RIVER3, and it is hoped that an improved version will become an addition to this library of routines. There are several specific areas that should be investigated and possibly improved. These are outlined in the following paragraphs.

The improvements that are envisaged for RIVER3 are primarily rearrangements of the input and output modes to allow greater flexibility and easier operation. For example, the channel analysis routines are capable of handling cross sections defined by arbitrary geometry, however, the reservoir characteristics must be defined by the two parameters K and w of the exponential relation (7.1). By providing the necessary input options, it would be possible to define arbitrary reservoir characteristics by using co-ordinates to define storage as a function of flow rate. RESRUT, a subroutine currently available in the CEPL, is capable of employing this type of data, or a simple rearrangement of the data, to perform a reservoir routing.

The basic command options that are available in RIVER3 have, by and large, been adopted from other similar programs currently found in the CEPL. However, several of the commands have been enlarged and other options have been added. As a result, the amount of operational data has increased. To alleviate the amount of input required from the teletype terminal, some of the data may need to be incorporated in files similar to the data file used for geometric data. Particular examples of this include the values of  $\boldsymbol{\alpha}$  and  $\boldsymbol{\beta}$  which may be included directly in the crosssection data. Reservoir data and inflow hydrograph information may be stored on files and read in at the user's command.

There are several suggestions that may also be investigated as major improvements to extend the power of RIVER3 and the CEPL. The item which could receive first priority is the provision of graph plotting capabilities for the program. This could be done either as a graph on the remote terminal or via storage of the necessary data to provide plots on the off line plotter associated with a computer facility. Other improvements could allow watershed (overland flow) simulation by the inclusion of routines to handle the data that describes the watershed as a large number of surfaces specified by length, width, slope and roughness. In addition, the routine would then require additional input capabilities to allow definition of the rainfall hyetograph(s). The third improvement would enable a user to verify the results obtained using a kinematic analysis by utilizing. the dynamic flood routing capabilities of the CEPL. This might be done by programming a command option into RIVER3 which would allow the user to generate the input cards necessary to perform the more rigorous analysis in a batch mode, for example.

There are many other items that could be included in a routine such as RIVER3, however, the selection of the items and the approach used to carry out these functions would depend upon individual tastes, and the degree of sophistication desired in the program. The application and further development of RIVER3 remains an almost limitless area for future work.

### CHAPTER 8

#### CONCLUSIONS

The preceeding chapters include detailed conclusions relating to the specific topics. This final chapter contains only a summary of the major points.

The objectives of this thesis were as follows:

- Provide a precise data base which can be utilized to compare various flood routing algorithms.
- (2) Investigate kinematic flood routing methods to develop a general framework for comparison of the variety of techniques found in the literature.
- (3) Determine the numerical characteristics of the general kinematic flood routing method. This will help determine the limitations of the finite difference schemes.
- (4) Compare the results of the kinematic simulations with the data base in an attempt to determine practical limitations of the kinematic algorithms.
- (5) Investigate methods of modelling attenuation with kinematic flood routing methods.

(6) Provide a versatile computer program that will enable a user to apply kinematic flood routing techniques and the methods of modelling attenuation in channel systems of arbitrary geometry.

These objectives have been achieved in the following manner.

Chapter 2 outlines the steps that were taken to provide the data base. An explicit staggered mesh finite difference scheme was used to provide a numerical solution to the differential equations describing the conservation of linear momentum and mass in prismatic channels of simple geometry. Tests were performed on the stability and convergence of the calculations to ensure as far as possible that the results of the numerical model were indeed reliable. Comparisons with other finite difference schemes which had been employed successfully show that the explicit method compared very well in representing unsteady flow in a system with simple geometry.

The general framework to compare the various kinematic algorithms was developed in Chapter 3. This method employs a rectangular finite difference molecule with the system variables defined at the corners of the molecule. (See figure 3.2) The difference between this method and the various other approaches to kinematic flood routing lies in the selection of the point about which the continuity equation is

expanded; a further distinction lies in the relationship adopted between the flow rate and the cross section area. Applying the continuity equation about the centre of the finite difference molecule models the elementary channel unit as an ideal channel section where storage is related equally to inflow and outflow. A reservoir is simulated when the continuity equation is expanded about a point on the downstream boundary of the molecule. The stability of the numerical schemes is affected significantly by the location of the nucleus, the point about which the continuity equation is expanded. Three areas of stability were identified, a zone of unconditional instability, a zone of conditional stability, and a point of unconditional stability. (See table A-1 in Appendix A.)

When the continuity equation is applied to a point on the downstream edge of the molecule and on the highest time level of the molecule, the scheme is unconditionally stable. The definition of the stable and unstable zones of the molecule was shown to depend on the size of the space and time incremental values of Q and  $\beta$ .

Moving the nucleus from the centre of the molecule results in an increase in the amount of error that is introduced in the numerical calculations. An analytic investigation of the finite difference equation of continuity indicates that first order errors are not introduced into the calculations when the nucleus is located at the centre of the

molecule. However, as the nucleus is moved in a downstream direction or to an increasing level of time, there is an increase in the amount of first order error introduced into the computations. These errors are always negative in sign and introduce a pseudoattenuation into the kinematic representation of a flood wave passing through a channel system.

Chapter 4 contains the results of a number of kinematic simulations compared with results produced by the solution of the momentum and continuity equations. An investigation was made of the order of magnitude of the various terms of the momentum equation to help evaluate the differences between the various simulations. The results show that as the order of magnitude of the terms reduce in comparison to the bedslope, the kinematic results agree closely with the more precise results produced by the numerical solution of the momentum and continuity equation. No guidelines were developed to indicate the difference in the relative size of the various terms in order to have close agreement between the results of the two methods of simulations. Preliminary results indicate that as the time base of the inflow hydrograph increases, the difference in the relative size of the terms can decrease in comparison with the relative size necessary to produce good results with a peaky, fast rising hydrograph.

Without intuitive or pre-calibrated use of the pseudo-attenuation phenomena referred to above, kinematic flood routing methods do not correctly predict attenuation of the flood wave as it passes through the channel. Chapters 5 and 6 outline modifications that can be made to the numerical model to introduce attenuation into the simulated flood wave. The method outlined in Chapter 5 is the introduction of a calibrated numerical error into the computational scheme. Alternately, the procedure used, moving the nucleus away from the centre of the finite difference molecule, may be viewed as modelling the elementary reaches as combination channel-reservoir units. The results of the numerical experiments indicate that in certain cases it is possible to simulate unsteady open channel systems accurately. This is true when dynamic effects demonstrated by an early rise of the outflow hydrograph, are not the dominant process in the physical system.

values of attenuated peak outflow obtained as function of the coefficients  $\boldsymbol{\alpha}$  and  $\boldsymbol{\beta}$ , defining displacement of the nucleus from the centre of the molecule, constituted a set of points comprising a well conditioned and slightly concave surface denoted by P( $\boldsymbol{\alpha}, \boldsymbol{\beta}$ ). The properties of this surface were examined and it was demonstrated to be a useful tool in the calibration of kinematic flood routing models based on

The series of numerical experiments demonstrated that the

results of either a dynamic (or 'complete') solution and/or actual prototype observations. The calibrated model may then in turn be used for further simulations. A detailed example of the use of this tool is included in Chapter 5.

Experiments, of a preliminary nature, were performed with System 1 to determine the effect of arranging identical elementary reaches in a cascade. The results, for the system modelled, indicate that as the number of elementary reaches increased the total attenuation increased. The amount of attenuation was found to vary with the number of reaches to the power of 0.55. (See equation 5.26)

Chapter 6 was devoted to studying the inclusion of a reservoir in series with a kinematic channel as a tool to model attenuation. The device of combining a channel and an imaginary reservoir in series is commonly used to model a hydrologic system. The method is known as the lag and route technique. The study reported in Chapter 6 utilized non-linear components in an attempt to identify the sensitivity of the solution to the reservoir location and different degrees of non-linearity in the reservoir. The numerical experiments indicated that the position of the imaginary reservoir affects the timing and shape of the hydrograph. Moving the reservoir toward the downstream boundary results in an outflow hydrograph that occurs sooner and rises more slowly. The simulations indicate that the amount of

attenuation is relatively insensitive to the location of the reservoir. Furthermore, the non-linearity of the reservoir affects the peak of the outflow hydrograph, the centroidal lag of the hydrograph and the 'skewness factor. However, the chord slope of the relationship between storage and discharge (as illustrated in figure 6.4) appears to be the dominant factor in determining the response characteristics.

Finally, Chapter 7 describes the development of an interactive computer model that is capable of performing kinematic flood routing in natural channels using the general method outlined in Chapter 3 as well as allowing for the inclusion of a reservoir at any section along the channel. Execution of the program is directed by commands entered as operational data. The general capabilities of the computer program include the ability to perform steady state analysis as well as allowing for the modification of the geometric data, which describes the channel, during the execution of the program.

Two examples of the program are included in the thesis. Both of these experiments demonstrate the effect of moving the nucleus from the centre of the molecule. Not only was there an increase of attenuation as the nucleus was moved from the centre; there was also an increase in the stability of the numerical analysis.

This computer program, known as RIVER3, has been designed to be compatible with the Civil Engineering Program Library. This

has been done to facilitate easy access by potential users of this program. There are an almost limitless number of modifications that can be made to RIVER3 to improve and expand its capabilities. Hopefully, this thesis has provided some insight into these various possibilities and enhances the understanding of kinematic flood routing techniques as very useful hydraulic and hydrologic simulation tools. 233

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### APPENDIX A

### STABILITY ANALYSIS OF THE GENERAL KINEMATIC METHOD

This appendix is devoted to the stability analysis of the general kinematic method. Stability is a primary consideration used to evaluate the performance of a numerical solution. An unstable scheme will cause small errors to amplify and dominate, thus, masking the solution.

This analysis is based on a linearized version of the continuity equation. For nonlinear equations, the method is not exact, but if the increment being considered is small and the coefficients of the derivatives are smooth functions, the approximation of constant coefficients is reasonable. Hopefully, this analysis will identify schemes which are obviously unstable.

The finite difference grid used to solve the continuity equation is:



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The continuity equation is:

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} = \bar{q}$$

Where:

Q = flow rate

 $A = c_{ross}$  section area

 $\mathbf{x} = \mathbf{distance}$ 

t = time

 $\overline{\mathbf{q}}$  = rate of lateral inflow

Because:

$$c = \frac{dQ}{dA}$$
 (A.2)

Where: c = kinematic wave velocity

The continuity equation may be rewritten as:

 $c\frac{\delta A}{\delta x} + \frac{\delta A}{\delta t} = \bar{q} \qquad (A.3)$ 

At any point J, K the numerical solution  $A_J^K$  is equal to the true solution A(  $J\Delta X$ ,  $K\Delta T$ ) plus an error term  $\overline{A}_J^K$  or:

$$A_{J}^{K} = A(J\Delta X, K\Delta T) + \bar{A}_{J}^{K}$$
 (A.4)

(A.I)

Because the system is linear, the error term can be written as one term of a Fourier Series:

$$\bar{A}_{J}^{K} = Ao EXP(i(m\sigma\Delta X + n\gamma\Delta T))$$
 (A.5)

Where:

Ao = constant

i = **√**-1

 $\sigma, \gamma$  = wave numbers in space and time

By writing the linearized finite difference equation of continuity in terms of the numerical solution and subtracting the true solution, the finite difference equation in terms of the errors is obtained. For the general kinematic method this equation is:

$$c\left(\frac{\beta(\bar{A}_{J}^{K}) + \bar{A}_{J}^{K+1} + (1-\beta)(\bar{A}_{J+1} - \bar{A}_{J}^{K})}{\Delta X}\right)$$

$$+ \frac{(I-\alpha)(\bar{A}_{J+i}^{K+1} - \bar{A}_{j+1}^{K}) + \alpha(\bar{A}_{J}^{K+1} - \bar{A}_{J}^{K})}{\Delta T} = 0 \quad (A.6)$$

For the point J, K, m and n may be assumed equal to zero,

with no loss of generality. Thus:

$$\bar{A}_{J}^{K} = Ao \qquad (A.7)$$

$$\bar{A}_{J}^{K+1} = A_{0} E X P(i Y \Delta T)$$
(A.8)

$$\bar{A}_{j+1}^{K} = Ao EXP(i\sigma\Delta X)$$
 (A.9)

$$\bar{A}_{J+l}^{(+)} = Ao EXP(i(\sigma \Delta X + Y \Delta T))$$
(A.10)

Solving equation A.6 for the unknown  $\vec{A}_{J+l}^{K+l}$  and multiplying by  $\Delta T$  yields:

$$C \frac{\Delta T}{\Delta X} \beta \bar{A}_{J+1}^{K+1} + (1-\alpha) \bar{A}_{J+1}^{K+1} = \alpha \bar{A}_{J}^{K} + (1-\beta) C \frac{\Delta T}{\Delta X} \bar{A}_{J}^{K}$$

$$- C \frac{\Delta T}{\Delta X} (1 - \beta) \bar{A}_{J+1}^{K} + (1 - \alpha) \bar{A}_{J+1}^{K}$$

$$+ C \frac{\Delta T}{\Delta X} \beta \bar{A}_{J}^{K+1} - \alpha \bar{A}_{J}^{K+1}$$
(A.11)

Collecting terms:

$$\left( (1-\alpha) + C\beta \frac{\Delta T}{\Delta X} \right) \bar{A}_{J+1}^{K+1} = \left( \alpha + (1-\beta)C \frac{\Delta T}{\Delta X} \right) \bar{A}_{J}^{K}$$

$$+ \left( (1-\alpha) - (1-\beta)C \frac{\Delta T}{\Delta X} \right) \bar{A}_{J+1}^{K}$$

$$+ \left( -\alpha + \beta C \frac{\Delta T}{\Delta X} \right) \bar{A}_{J}^{K+1}$$

$$(A.12)$$

To facilitate easy algebraic manipulation, the following substitutions are employed, the D terms corresponding to the numbered nodes of Figure 3.2.

$$DI = \alpha + (I - \beta)C \frac{\Delta I}{\Delta X}$$
(A.13)

$$D2 = -\alpha + \beta C \frac{\Delta T}{\Delta X}$$
 (A.14)

$$D3 = (I - \alpha) - (I - \beta)C\frac{\Delta T}{\Delta X}$$
(A.15)

(A.I6)

$$D4 = (1 - \alpha) + \beta C \frac{\Delta T}{\Delta X}$$

Thus, equation A.12 becomes:

$$D4\bar{A}_{J+1}^{K+1} = DI\bar{A}_{J}^{K} + D2\bar{A}_{J}^{K+1} + D3\bar{A}_{J+1}^{K}$$
 (A.17)

To ensure stability, errors must not amplify, thus:  $\begin{vmatrix} \bar{A} & K+1 \\ J+1 \\ \hline{A} & J+1 \\ \hline{A} & J+1 \end{vmatrix} = 1$ (A.18)

The error term must lie within the unity circle. That is:

$$|EXP(i(\sigma\Delta X + \gamma\Delta T))| = \left|\frac{D3\bar{A}_{J+1}^{K}}{D4}\right| \left|\frac{D1\bar{A}_{J}^{K} + D2\bar{A}_{J}^{K+1} + D3\bar{A}_{J+1}^{K}}{D3\bar{A}_{J+1}^{K}}\right| \leq I \quad (A.19)$$

Inserting the error terms written as components of a Fourier

Series gives:

$$|EXP(i(\sigma \Delta X + \gamma \Delta T))| = \left| \frac{DI + D2EXP(i\gamma \Delta T) + D3EXP(i\sigma \Delta X)}{D4} \right| \leq I (A.20)$$

 $t_{\rm c}$ 

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equation A. 20:

 $EXP(i\sigma\Delta X) = COS \sigma\Delta X + iSIN \sigma\Delta X$  (A.21)

 $EXP(iY \Delta T) = COSY \Delta T + iSINY \Delta T$  (A.22)

Substituting equations A, 21 and A, 22 into equation A, 20 yields:

$$|EXP(i(\sigma \Delta X + \gamma \Delta T))| = \frac{1}{|D4|} |DI + D2(\cos \gamma \Delta T + i \sin \gamma \Delta T)$$

+ D3(COS  $\sigma \Delta X$  + iSIN  $\sigma \Delta X$ ) = [

The worst case will occur when:

 $\cos \sigma \Delta X + i \sin \sigma \Delta X = \pm 1$ 

and

 $\cos \sqrt{\Delta T} + i \sin \sqrt{\Delta T} = \pm I$ 

(A.25)

(A.24)

(A.23)

Thus, to ensure stability. the following equation should be satisfied:

Rearranging equation A. 26 gives:

Expanding equation A, 27 leads to:

$$\begin{vmatrix} \alpha + (1-\beta) c \frac{\Delta T}{\Delta X} \end{vmatrix} + \begin{vmatrix} -\alpha + \beta c \frac{\Delta T}{\Delta X} \end{vmatrix}$$
$$+ \begin{vmatrix} (1-\alpha) - (1-\beta) c \frac{\Delta T}{\Delta X} \end{vmatrix} - \begin{vmatrix} (1-\alpha) + \beta c \frac{\Delta T}{\Delta X} \end{vmatrix} \neq 1$$
(A.28)

(A.26)

(A.27)

 $D_1$  and  $D_4$  will always be positive thus:

$$\begin{vmatrix} -\alpha & +\beta c \frac{\Delta T}{\Delta X} \end{vmatrix} + \begin{vmatrix} (1-\alpha) - (1-\beta) c \frac{\Delta T}{\Delta X} \end{vmatrix}$$
$$+ \alpha + (1-\beta) c \frac{\Delta T}{\Delta X} - 1 + \alpha - \beta c \frac{\Delta T}{\Delta X} \leq 0$$
(A.29)

Or:

$$\left| -\alpha + \beta C \frac{\Delta T}{\Delta X} \right| + \left| (1 - \alpha) - (1 - \beta) C \frac{\Delta T}{\Delta X} \right|$$
$$+ (2\alpha - 1) + (1 - 2\beta) C \frac{\Delta T}{\Delta X} \leq 0$$
(A.30)

If the first two terms of equation A. 30 are of the same sign, then:

$$(1-2\alpha) + (2\beta-1)C\frac{\Delta T}{\Delta X}$$

$$+(2\alpha - 1) + (1 - 2\beta)C\frac{\Delta T}{\Delta X} = 0$$
 (A.31)

Defining:

$$S = (2\alpha - 1) + (1 - 2\beta)C\frac{\Delta T}{\Delta X}$$
 (A.32)

Where:

S = Stability number

$$|-S| + S \neq 0$$
 (A.33)

Thus stability could be achieved if:

and if the following two terms have the same sign.

$$CI = -\alpha + \beta C \frac{\Delta T}{\Delta X}$$
(A.35)  
$$C2 = (I - \alpha) - (I - \beta) C \frac{\Delta T}{\Delta X}$$
(A.36)

If the two conditions do not have the same sign then the following equation must be satisfied in order to ensure stability.

$$|C| + |C2| + S = 0$$
 (A.37)

Several examples are presented to demonstrate the application of these rules.

Consider the location of the nucleus at a point specified by

$$\boldsymbol{\alpha} = 0.0$$
$$\boldsymbol{\beta} = 0.0$$
$$C_1 = 0.0$$
$$C_2 = \mathbf{I} - \mathbf{C} \frac{\mathbf{\Delta T}}{\mathbf{\Delta X}}$$

At that point:

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Thus the two terms have the same sign and the only condition for stability is equation A. 34:

$$S = -I + C \frac{\Delta T}{\Delta X} \leq 0$$

Or:

If the nucleus was located at a point specified by:

$$\alpha = 0.0$$
  
 $\beta = 0.5$ 

then:

$$C_{1} = 0.5 C \frac{\Delta T}{\Delta X}$$
$$C_{2} = 1 - 0.5 C \frac{\Delta T}{\Delta X}$$

The two conditions will have the same sign if

$$C\frac{\Delta T}{\Delta X} \leq 2$$

and it will be necessary to satisfy only equation A.34.

For  $C \frac{\Delta T}{\Delta X} \ge 2$ , equation A. 37 must be satisfied:

$$\left| 0.5 \, \mathrm{C} \frac{\Delta \mathrm{T}}{\Delta \mathrm{X}} \right| + \left| 1 - 0.5 \, \mathrm{C} \frac{\Delta \mathrm{T}}{\Delta \mathrm{X}} \right| - 1 = 0$$

$$C\frac{\Delta T}{\Delta X} \neq 1$$

$$C\frac{\Delta T}{\Delta X} - 2 \neq 0$$

This condition is satisfied only when  $C\frac{\Delta T}{\Delta X} \leq 2.0$ . Thus it was concluded that the solution is stable only under the following condition.

$$C\frac{\Delta T}{\Delta X} \leq 2$$

This type of analysis was continued for the various points specified in table A.1.

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Or:

## RESULTS OF STABILITY ANALYSIS

$$KN = C \frac{\Delta T}{\Delta X}$$

GENERAL METHOD:

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β	1.00	0.75	0.50	0.25	0.00
1.00	∞≥KN≥I	∞≥KN≥ <u>3</u> 4	∞≥KN≥ <u>1</u> 2	∞≥KN≥ <u>1</u>	∞≐KN≏O
0.75	UNSTABLE	KN = 1	2 = KN= <u>2</u>	3≥K⊵: <u>1</u>	4≥KN≥0
0.50	UNSTABLE	UNSTABLE	KN = I	<u>3</u> ≥KN≥ <u>1</u> 2	5 <b>≠ K N ≥ O</b>
0.25	UNSTABLE	UNSTABLE	UNSTABLE	KN = 1	<u>4</u> ≥KN≥O
0.00	UNSTABLE	UNSTABLE	UNSTABLE	UNSTABLE	I ≥ KN ≥ O

LAX - WENDROFF: KN = I

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### APPENDIX B

### DEGREE OF APPROXIMATION FOR THE DYNAMIC ANALYSIS

By expanding the differential formulations of the momentum and continuity equations using a Taylor Series and observing the terms that are truncated when the finite difference approximations are made, it is possible to determine the degree of approximation.

The Taylor Series for a function of two variables is:

$$f(x+g,t+h) = f(x,t) + g \frac{\delta f}{\delta x} + h \frac{\delta i}{\delta t} + \frac{g^2 \delta^2 f}{2 \delta x^2} + \frac{h^2 \delta^2 f}{2 \delta t^2} + gh \frac{\delta^2 f}{\delta x \delta t} + \frac{g^3 \delta^3 f}{6 \delta x^3} + \frac{g^2 h}{2 \delta x^2 \delta t} + \frac{gh^2}{2 \delta x \delta t^2} + \frac{h}{6 \delta x^3} \delta^3 f + \cdots$$
(B.1)

For the continuity equation:

$$g = \Delta X/2$$
$$h = \Delta T/2$$

$$\frac{\delta Q}{\delta x} = \frac{\Omega(x + \Delta X/2, t) - \Omega(x - \Delta X/2, t)}{\Delta X}$$
$$= \frac{1}{\Delta X} \left( \Delta X \frac{\delta Q}{\delta x} + \frac{\Delta X^3}{24} \frac{\delta^3 Q}{\delta x^3} \cdots \right)$$
(3.2)

$$\frac{\delta A}{\delta t} \stackrel{\text{d}}{=} \frac{A(x, t + \Delta T/2) - A(x, t - \Delta T/2)}{\Delta T}$$

$$\stackrel{\text{d}}{=} \frac{1}{\Delta T} \left( \Delta T \frac{\delta A}{\delta t} + \frac{\Delta T^3}{24} \frac{\delta^3 A}{\delta t^3} \cdots \right) \qquad (B.3)$$

Substituting into the continuity equations yields:

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} + O(\Delta X^2, \Delta T^2) = \bar{q}$$
 (B.4)

For the momentum equation, the following approximations will be found:

$$\frac{\delta Q}{\delta t} \stackrel{\sharp}{=} \frac{1}{\Delta T} \left( \Delta T \frac{\delta Q}{\delta t} + \frac{\Delta T^3}{24} \frac{\delta^3 Q}{\delta t^3} \cdots \right)$$
(B.5)  
$$\frac{\delta h}{\delta x} \stackrel{\sharp}{=} \frac{1}{\Delta X} \left( \Delta X \frac{\delta h}{\delta x} + \frac{\Delta X^3 \delta^3 h}{24 \delta x^3} \cdots \right)$$
(B.5)  
and

$$\frac{\delta h}{\delta t} = \frac{1}{\Delta T} \left( \Delta T \frac{\delta h}{\delta t} + \frac{\Delta T^3}{24} \frac{\delta^3 h}{\delta t^3} + \frac{\Delta T \Delta X^2}{8} \frac{\delta^3 h}{\delta x^2 \delta t} \cdots \right) \quad (B.7)$$

Therefore the approximation is in the order of  $\Delta T^2$  and  $\Delta X^2$ . However, another approximation should be analyzed. That is the manner in which  $Q^2$  terms are represented.

$$Q(x,t)^2 \cong Q(x,t + \Delta T/2) Q(x,t - \Delta T/2)$$
 (B.8)

$$Q(x,t)^2 \equiv Q(x,t) - \frac{\Delta T^2}{4} \left( \frac{\delta Q}{\delta t} \right) + \cdots$$
 (B.9)

Thus, the approximation of the momentum equation is of the order of  $\Delta X^2$  and  $\Delta T^2$ .

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### APPENDIX C

### DEGREE OF APPROXIMATION FOR THE GENERAL KINEMATIC ROUTING METHOD

Analyzing the degree of approximation involves expanding the differential equation into a Taylor Series and determining what terms are truncated when the finite difference approximations are made.

For a function, f(x, t), the value of the function at f(x+g, t+h)is given by the Taylor Series:

$$f(x+g,t+h) = f(x,t) + g\frac{\delta f}{\delta x} + h\frac{\delta f}{\delta t} + \frac{g^2}{2}\frac{\delta^2 f}{\delta x^2}$$

+ 
$$gh \frac{\delta^2 f}{\delta x \, \delta t}$$
 +  $\frac{h^2 \, \delta^2 f}{2 \, \delta t^2}$  +  $\frac{g^3 \, \delta^3 f}{6 \, \delta x^3}$   
+  $\frac{g^2 h}{2 \, \delta x^2 \, \delta t}$  +  $\frac{gh^2}{2 \, \delta x \, \delta t^2}$  +  $\frac{h^3 \, \delta^3 f}{6 \, \delta t^3}$  + ... (C.1)

To obtain the general kinematic flood routing method, the continuity equation is applied using finite difference methods, to a space time diagram as shown below.



The finite difference approximations are:

ΔT

δt

$$\frac{\delta Q}{\delta x} = \frac{\beta(Q_4 - Q_2) + (1 - \beta)(Q_3 - Q_1)}{\Delta X}$$
(C.2)  
$$\frac{\delta A}{h} = \frac{(1 - \alpha)(A_4 - A_3) + \alpha(A_2 - A_1)}{(C.3)}$$

Expanding about point A:

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$$Q_{4} - Q_{2} = (x_{1} + x_{2}) \frac{\delta Q}{\delta x} + \frac{1}{2} (x_{1}^{2} - x_{2}^{2}) \frac{\delta^{2} Q}{\delta x^{2}}$$

$$+ (x_{1}t_{2} + x_{2}t_{2}) \frac{\delta^{2} Q}{\delta x \, \delta t} + \frac{1}{6} (x_{1}^{3} + x_{2}^{3}) \frac{\delta^{3} Q}{\delta x^{3}}$$

$$+ \frac{1}{2} (x_{1}^{2}t_{2} - x_{2}^{2}t_{2}) \frac{\delta^{3} Q}{\delta x^{2} \, \delta t} + \frac{1}{2} (x_{1} t_{2}^{2} + x_{2} t_{2}^{2}) \frac{\delta^{3} Q}{\delta x \, \delta t^{2}} + \cdots (C.4)$$

$$Q_{3} - Q_{1} = (x_{1} + x_{2}) \frac{\delta Q}{\delta x} + \frac{1}{2} (x_{1}^{2} - x_{2}^{2}) \frac{\delta^{2} Q}{\delta x^{2}}$$

$$- (x_{1}t_{1} + x_{2}t_{1}) \frac{\delta^{2} Q}{\delta x \, \delta t} + \frac{1}{6} (x_{1}^{3} + x_{2}^{3}) \frac{\delta^{3} Q}{\delta x^{3}}$$

$$- \frac{1}{2} (x_{1}^{2}t_{1} - x_{2}^{2}t_{1}) \frac{\delta^{2} Q}{\delta x^{2} \, \delta t} + \frac{1}{2} (x_{1} t_{1}^{2} + x_{2} t_{1}^{2}) \frac{\delta^{3} Q}{\delta x \, \delta t^{2}} + \cdots (C.5)$$

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Substituting expressions for  $Q_4 - Q_2$  and  $Q_3 - Q_1$  into equation C-2, the following expression is obtained:

$$\frac{\delta Q}{\delta x} = \frac{1}{\Delta X} \left( \Delta X \frac{\delta Q}{\delta x} + \frac{1}{2} (x_1^2 - x_2^2) \frac{\delta^2 Q}{\delta x^2} + \frac{1}{6} (x_1^3 + x_2^3) \frac{\delta^3 Q}{\delta x^3} + \beta (x_1 t_2 + x_2 t_2) \frac{\delta^2 Q}{\delta x \ \delta t} - (x_1 t_1 + x_2 t_1) \frac{\delta^2 Q}{\delta x \ \delta t} + \beta (x_1 t_1 + x_2 t_1) \frac{\delta^2 Q}{\delta x \ \delta t} + \frac{\beta}{2} (x_1^2 t_2 - x_2^2 t_2) \frac{\delta^3 Q}{\delta x^2 \ \delta t} - \frac{1}{2} (x_1^2 t_1 - x_2^2 t_1) \frac{\delta^3 Q}{\delta x^2 \ \delta t} + \frac{\beta}{2} (x_1^2 t_1 - x_2^2 t_1) \frac{\delta^3 Q}{\delta x^2 \ \delta t} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x^2 \ \delta t} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{1}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{1}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\beta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta^3 Q}{\delta x \ \delta t^2} + \frac{\delta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta}{\delta x \ \delta t^2} + \frac{\delta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta}{\delta x \ \delta t^2} + \frac{\delta}{2} (x_1 t_1^2 + x_2 t_1^2) \frac{\delta}{\delta x \ \delta t^2} + \frac{\delta}{\delta x \ \delta t^2} +$$

expressed in the following manner:

$$x_1 = \alpha \Delta X$$
  $x_1^2 = \alpha^2 \Delta X^2$ 

$$x_2 = (1 - \alpha) \Delta X \qquad \qquad x_2^2 = (1 - 2\alpha + \alpha^2) \Delta X^2$$

$$t_1 = \beta \Delta T$$
  $t_1^2 = \beta^2 \Delta T^2$ 

 $t_1^2 = (1 - \beta) \Delta T \qquad \qquad t_2^2 = (1 - 2\beta + \beta^2) \Delta T^2$ 

This yields:

$$\frac{\delta Q}{\delta x} = \frac{1}{\Delta X} \left( \Delta X \frac{\delta Q}{\delta x} + \frac{1}{2} (2\alpha - 1) \Delta X^2 \frac{\delta^2 Q}{\delta x^2} + \frac{1}{6} (1 - 3\alpha + 3\alpha^2) \Delta X^3 \frac{\delta^3 Q}{\delta x^3} + \frac{1}{2} (\beta - \beta^2) \Delta X \Delta T^2 \frac{\delta^3 Q}{\delta x \delta t^2} + \cdots \right) \quad (C.7)$$
  
Thus:

$$\frac{\delta Q}{\delta x} = \frac{\delta Q}{\delta x} + O(\Delta X^2, \Delta T^2)$$
 (C.8)

when

 $\alpha = 0.5$ 

If  $\alpha \neq 0.5$ 

$$\frac{\delta Q}{\delta x} = \frac{\delta Q}{\delta x} + O(\Delta X, \Delta T^2)$$
 (C.9)

The magnitude of the  $O(\Delta X)$  error is related to (2 $\alpha$ -1).

The general description of truncation error is:

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$$\frac{\delta Q}{\delta x} = \frac{\delta Q}{\delta x} + (2\alpha - 1)O(\Delta X) + O(\Delta X^2, \Delta T^2)$$
 (C.10)

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Similarly expanding the  $\frac{\delta A}{\delta t}$  term about point A:

$$\frac{\delta A}{\delta t} = \frac{1}{\Delta T} \left( \Delta T \frac{\delta A}{\delta t} + \frac{1}{2} (1 - 2\beta) \Delta T^2 \frac{\delta^2 A}{\delta t^2} \right)$$

$$\div \frac{1}{6} (1 - 3\beta + 3\beta^2) \Delta T^3 \frac{\delta^3 A}{\delta t^3} + \frac{1}{2} (\alpha - \alpha^2) \Delta X^2 \Delta T \frac{\delta^3 A}{\delta x^2 \delta t} \cdots ) (C.11)$$

When  $\beta = 0.5$  the approximation will be

$$\frac{\delta\Lambda}{\delta t} = \frac{\delta\Lambda}{\delta t} + O(\Delta X_{t}^{2} \Delta T^{2})$$
(C.12)  
If  $\beta \neq 0.5$ 

$$\frac{\delta \Lambda}{\delta t} = \frac{\delta \Lambda}{\delta t} + O(\Delta X^2, \Delta T)$$
 (C.13)

Generally, the approximation may be written as

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$$\frac{\delta \Lambda}{\delta t} = \frac{\delta A}{\delta t} + (1 - 2\beta) O(\Delta T) + O(\Delta X^2, \Delta T^2)$$
 (C.14)

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The expression for the general kinematic flood routing method can be written as follows:

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} + (2\alpha - 1)O(\Delta X) + (1 - 2\beta)O(\Delta T)$$

$$+ O(\Delta X_{1}^{2} \Delta T^{2}) = \hat{q}$$
 (C.15)

As the location about which the equation of continuity is moved further away from the centre of the molecule, the amount of error introduced by approximations on the order of  $\Delta X$  and  $\Delta T$ increases linearly.

As  $\boldsymbol{\alpha}$  is varied, terms on the order of  $\boldsymbol{\Delta X}$  are modified while varying  $\boldsymbol{\beta}$  changes the way the errors on the order of

 $\Delta T$  are introduced.

A plot of  $(2 \ \mathbf{a} \ -1)$  over the molecule will reveal a plane sloping toward the downstream side of the molecule (denoted by points 3 and 4). This is shown in figure C-1. Similarly  $(1-2 \ \mathbf{\beta})$  will be a plane tipped in the direction of increasing time (denoted by points 2 and 4).

Because the errors  $O(\Delta X)$  and  $O(\Delta T)$  introduced by varying  $\alpha$  or  $\beta$  respectively, are independent of each other, it would be expected that a plot of errors would be a plane which has a dip

and strike determined by the relative size of the errors as well

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as the sign of the errors.

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## PLOT OF $(2\alpha - 1)$





# PLOT OF (1-2B)

# FIGURE C.2

<u>|</u>...

### APPENDIX D

### DOCUMENTATION OF

### COMPUTER ROUTINES

\*\*\*FINDIF

\*\*\*KINDIF

HPLOT

KINFUN

KINRUT

RESVOR

RIVER3

#### UNSTEADY FLOW ANALYSIS

- PURPOSE: This program provides a finite difference solution of the partial differential equations describing one-dimensional unsteady flow in a rectangular channel.
- (2) METHOD: The input data is supplied to the program using data cards. Problem variables are initialized and the input is printed to provide a check before the solution of the equations is begun. An explicit finite difference scheme based on a staggered mesh space-time diagram is utilized to provide the solution. Output is in the form of printer plots which show hydrographs at any two sections defined by the user. Depth versus time curves are also plotted for any two sections which are specified at the time of execution. Similarly. a stage discharge curve is plotted for a user defined section. Options are available to have the hydrographs. etc., punched on cards for utilization with other programs and for storage purposes.

- (3) PROGRAM:
- (a) DECK NAME: FINDIF
- (b) CALLING SEQUENCE: None. This is a driving program.
- (c) INPUT:

Data Card Type	Data	Format			
1	Length of channel, width,	F10.1, F10.2			
	Mannings n, slope. depth	2F10, 8. F10.2			
2	Number of reaches, Courant Number for choice of time				
	increment, time at which				
	the analysis is to be terminated				
	(seconds).	13, 2F10.2			
3	Test number (for identification) 13				
4	Punch code to produce hydrographs				
	and depth versus time plots.				
	Punch code for stage-discharge				
,	curve. (Enter 0 for no cards, 1				
	for punched output.)	213			
5	Section for which stage-discharge				
	curve should be plotted.	13			
б	Section numbers for which	hyd <b>ro-</b>			
	graphs are to be plotted.	213			

7 Section numbers for which stage versus time curves are to be plotted. 213 8 Number of points defining inflow hydrograph. 13 9 The ratio of flow to full bank flow and time (seconds) for one point on the inflow hydrograph. F10.2, F10.0 Note: Repeat card 9 for each of the points which define the inflow hydrograph. The output consists of a listing of various OUTPUT: variables for the problem being tackled as well as two hydrographs, two stage versus time curves, and a stage discharge curve. Options allow the stage discharge curve, hydrographs, and depth versus time curves to be punched on cards for future use. (Also see discussion.) RESTRICTIONS: This program is limited to the solution of

(d)

(e)

problems involving uniform rectangular channels where flow resistance is defined by Mannings equation. The time step is determined from full bank flow conditions and the Courant Number. Stability problems

may be encountered if the Courant Number is too large. (The program is unable to simulate supercritical flow.) A uniform flow depth is assumed as the downstream control. (See discussion.)

(f) OTHER DECKS: INTER1, MANNGQ, PLOTPT, OUTPLT

(4) EXAMPLE:

(a) THE PROBLEM: A rectangular channel 50,000' long, 100' wide and 20' deep is subject to a triangular inflow hydrograph. Twenty-five reaches are utilized in the analysis and bedslope = 0.0002 is specified with a Mannings n = 0.0149. The Courant Number vas set equal to 0.5 and the analysis was carried out over a period of 24,000 seconds. A copy of the data cards and sample output is shown in Appendix G of the source.
(5) DI SCUSSION: New users should refer to documentation of all related program and routines. Further

details of the method of analysis is available from the source.

A variation of this program was developed to perform the same type of calculations with a wide channel. This program was further

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modified to provide output which could be used to generate an animated movie of the flood wave moving down the channel.

 (6) SOURCE: "A Comparison of Kinematic Flood Routing Methods". by Fred Biesenthal (A Master of Engineering Thesis) McMaster University Hamilton, Ontario

### UNSTEADY FLOW ANALYSIS

- (1) PURPOSE: This program models the movement of a flood wave down a uniform rectangular channel using a kinematic flood routing technique. For each problem tackled, twenty-five solutions are provided. This enables a user to study the effects of varying the position of the "nucleus" within the finite difference "molecule".
- (2)METHOD: After obtaining the input data, the program variables are initialized, and the flood routing portion is carried out. The actual flood routing calculations are performed by subroutine "KINRUT" and a further description of the algorithm may be obtained from the documentation of that subroutine as well as from the source. The flood routing computations are incorporated within nested "DO" loops so that the parameters ALPHA and BETA are systematically varied to provide solutions with twentyfive positions of the nucleus. With each solution a title page and two graphs are provided to document the results of the simulations.

(3) PROGRAM:

(a) DECK NAME: KINDIF

- (b) CALLING SEQUENCE: None, this is a driving program.
- (c) INPUT:

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Data Card Type	Data	Format			
1	Length of channel, width, Fl	.0.1, F10.2,			
	Manning n, slope, depth 21	F10.8. F10.2			
2	Number of reaches, time				
	step (seconds). Time at				
	which analysis is to be				
	terminated (seconds). I3	, 2F10.2			
3	Test number (for identif-				
	ication)	13			
4	Punch code to produce hydro-				
	graphs on cards.	13			
5	Section numbers for which				
	hydrographs are to be plotted	213			
6	Section numbers for which				
	stage versus time curves are				
,	to be plotted.	213			
7	Number of points defining				
	inflow hydrograph	13			
8	The ratio of flow to fullbank				
	flow and time (seconds) for one	2			
	point on the inflow hydrograph.	F10.2. F10.0			
Note: Repeat card 8 for each of the points which define the inflow hydrograph.

- (d) OUTPUT: The output consists of a listing of various parameters for the problem being tackled in the form of a title block as well as two hydrographs and two stage - time curves for each of the solutions attempted. An option is incorporated in the program to allow a user to record the hydrographs on punched cards.
- (e) RESTRICTIONS: This program is limited to the solution of problems involving uniform rectangular channels where flow resistance is defined by Mannings equation. (See discussion.)
- (f) OTHER DECKS: HPLOT, INTERI, MANNGQ, KINRUT, KINFUN
- (4) EXAMPLE:
- (a) THE PROBLEM: A rectangular channel 50,000' long 100' wide and 20' deep is subject to a triangular inflow hydrograph. Twenty reaches are used to analyze the channel which has a bedslope of 0.0002 and a Mannings "n" of 0.0149. The time step was set equal to 200 seconds and the analysis was made for a period of 24,000 seconds. A copy of the data input

cards and a sample output is shown in Appendix G of the source.

(5) DISCUSSION: New users should refer to documentation of all related programs and routines. Further details of the method of analysis is available from the source.

> A variation of this program was developed to perform the same calculations with a very wide channel.

Another variation was developed which allowed the user to install an imaginary reservoir in series with the channel. This reservoir could be located at a section in the channel and the size of the reservoir was defined in the input The "nucleus" was located in the centre of data. the finite difference molecule when the channel was analyzed. Nested "DO" loops were incorporated to allow several simultation to be performed with different sizes of reservoirs. "A Comparison of Kinematic Flood Routing Methods", by Fred Biesenthal (A Master of Engineering Thesis). McMaster University Hamilton, Ontario

(6) SCURCE:

### HYDROGRAPH PLOTS

- PURPOSE: This subroutine plots two hydrographs on a set of axis.
- A grid. which is 50 printer lines high and METHOD: (2) 118 spaces wide, is utilized. The hydrograph is specified by a series of zeros or plus signs positioned in the grid. Automatic scaling is provided within the subroutine and the calibration marks are placed on the scales. Time units are plotted (and labeled) as hours) across the bottom scale. An interpolation routine allows hydrographs with two different time steps to be plotted on the same scale. (The largest time step is used for the time scale.) For further details of the method refer to the subroutine listing.
- (3) PROGRAM:
- (a) DECK NAME: HPLOT
- (b) CALLING SEQUENCE: HPLOT (QOUT2D. DTIMAR, NQOU'T, PEAKAR, TSTART) QOUT2D = A two dimensional floating

point array containing the flowrates of the two hydrographs to be plotted. The flow rates are specified at equal time intervals though the time step for the two hydrographs need not be the same value.

- DTIMAR = A floating point array containing two values. The first value is the time step for the first hydrograph and the second value contains the time step for the second hydrograph.
- NQOUT = A floating point array defining the number of points in each hydrograph.
- PEAKAR = A floating point array which defines the peak flow rate of each hydrograph.
- TSTART = A floating point value which defines the start time of the hydrograph,

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(c) OUTPUT: A call of this subroutine results in a plot
 of two hydrographs on one page of com puter paper. A line printer which is at
 least 130 spaces wide is required.

- (d) PESTRICTIONS: Negative values will not be plotted correctly.
- (e) OTHER DECKS REQUIRED: None
- (4) EXAMPLE: The sample printouts available from the source provide examples of this subroutine.
- (5) DISCUSSION: The use of the subroutine is restricted to batch mode operations. It is possible, with further modification to utilize this approach with interactive terminals.
- (6) SOURCE: Modified from Williams, Jimmy R. and
   Roy W. Hann, "HYMO, A Problem Oriented
   Language for Building Models", Water
   Resources Research, Vol. 8, No. 1, 79 86. February, 1972.

FLOOD ROUTING USING THE CONTINUITY EQUATIONS

- PURPOSE: This routine calculates the functional relationships of an elementary river reach prior to the routing of a floodwave using subroutine "KINRUT".
- (2) METHOD: The functional relationships are defined as follows:

FUS = BETA # QIN - ALPHA \* STOR/DTIM

for the upstream section and

GDS = BETA \* QOUT + (I.- ALPHA) \* STOR / DTIM

for the downstream section.

It should be noted that STOR is the storage in the total length of the elementary reach with a steady state condition. In the case of FUS(QIN), the flow rate along the elementary reach is assumed to be that of the inflow while a steady flow rate of QOUT along the elementary reach is assumed when calculating GDS(QOUT).

- (3) PROGRAM:
- (a) DECK NAME: KINFUN

(b) CALLING

SEQUENCE: CALL KINFUN (ALPHA, BETA, DTIM, STORUS, QUSAR, NFUS, STORDS, QDSAR, NFDS, FUSAR, GDSAR)

WHERE ALPHA = Floating point variable containing ALPHA.

- BETA = Floating point variable containing BETA.
- DTIM = Floating point variable containing the value of the time increment.
- STORUS = Floating point array of size NFUS which contains values for steady state storage to define coordinates of storage - inflow relationship.
- QUSAR = Floating point array of size NFUS which contains values of inflow to define coordinates of storage inflow relationship.
- NFUS = Integer describing the number of points which define the relationship between steady state storage and inflow.
- STORDS = Floating point array of size NFDS
  which contains values for steady
  state storage to define coordinates

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of storage - outflow relationship.

- QDSAR = Floating point array of size NFDS which contains values of outflow to define coordinates of storage outflow relationship.
- NFDS = Integer describing the number points which define the relationship between steady state storage and outflow.
- FUSAR = Floating point array of size NFUS containing values of the function relationship associated with inflow.
- GDSAR = Floating point array of size NFDS containing values of the function relationship associated with outflow.
- (c) OUTPUT FORMAT: The function relations are assigned to array FUSAR for the upstream section and GDSAR for the downstream section.
- (d) OTHER DECKS REQUIRED: None
- (e) RESTRICTIONS: The relationship between storage and flow rate must be a single valued function.
  - (4) EXAMPLE: Refer to the source for an example.
  - (5) DISCUSSION: None

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(6) SOURCE: "A Comparison of Kinematic Flood Routing

Methods", by Fred Biesenthal

(A Master of Engineering Thesis),

McMaster University

Hamilton, Ontario

FLOOD ROUTING USING THE CONTINUITY EQUATIONS

- PURPOSE: This routine is used to determine the outflow from a system where the unsteady flow regime can be reasonably modelled using kinematic wave theory.
- (2) METHOD: The routine uses a finite difference technique to solve the continuity equation. Two parameters ALPHA and BETA are used to determine the way in which the finite differences are calculated. Figure KINRUT1 shows the finite difference approximations and a portion of the space time grid. The calculations are made using functional relationships of inflow and outflow. The process of routing a flood through an elementary reach is very similar to the process of routing a flood through a reservoir using the storage indication method. A user should refer to the source for a more complete description of the numerical system used and the properties of the finite difference scheme. Refer to the documentation of KINFUN for the definition of the functional relationships.



GRAPHICAL REPRESENTATION

NUMERICAL APPROXIMATIONS



$$\frac{\delta A}{\delta T} = \frac{(I-\alpha)(A_4-A_3) + \alpha(A_2-A_1)}{\Delta T}$$

$$\frac{\delta Q}{\delta X} = \frac{\beta(Q_4 - Q_2) + (I - \beta)(Q_3 - Q_1)}{\Delta X}$$

- MOLECULE THE SPACE-TIME ELEMENT BOUNDED BY POINTS 1243
- NUCLEUS THE POINT "P" ABOUT WHICH THE FINITE DIFFERENCE EQUATION IS APPLIED

- (3) PROGRAM:
- (a) DECK NAME: KINRUT
- (b) CALLING SEQUENCE: CALL KINRUT (GDSAR, QDSAR, NFDS, FUSAR, QUSAR, NFUS, QLINAR, TIMARL, NPTSL, DX, DTIM, NPTS, QINAR, QOUTAR)
  - WHERE GDSAR = Floating point array of size NFDS containing values for the functional relation of outflow.
    - QDSAR = Float point array of size NFDS containing values of outflow related to the array GDSAR.
    - NFDS = Integer defining number of points which describe the function relation of outflow.
    - FUSAR = Floating point array of size NFUS containing values for the functional relation of inflow.
    - QUSAR = Floating point array of size NFUS containing values for inflow related to the array FUSAR.
    - NFUS = Integer defining the number of points which describe the function relation of inflow.

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- QLINAR = Floating point array of size NPTSL containing the values of flow per unit length which describe the lateral inflow hydrograph.
- TIMARL = Floating point array of size NPTSL containing the coordinates of time which describe the lateral inflow hydrograph.
- NPTSL = Integer defining the number of points which describe the lateral inflow hydrograph.
- DX = Floating point variable describing the length of the channel or reservoir which is subject to lateral inflow.
- DTIM = Floating point variable containing the time increment.
- NPTS = Integer defining the number of points which describe the inflow and outflow hydrographs.
- QINAR = Floating point array of size NPTS which contains the points describing the inflow hydrograph.

QOUTAR = Floating point array of size NPTS

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which contains the computed outflow hydrograph.

(c) OUTPUT FORMAT: The outflow hydrograph is assigned to array QOUTAR. Points of the outflow hydrograph are specified at equal time increments of value DTIM.

(d) OTHER DECKS REQUIRED: INTER 1

(e)

RESTRICTIONS: The user must specify the first value of the outflow hydrograph before calling KINRUT, and the points in the inflow hydrograph must be defined at equal time increments. Coordinates of the outflow hydrograph will correspond to the points on the inflow hydrograph.

> If lateral inflow is not a factor in the particular problem, it will be set to zero by defining the first value of TIMARL as a real negative value.

The functional relationships between flow and storage must be defined before calling KINRUT.

(4) EXAMPLE; Refer to the Source.

(5) DISCUSSION:

None

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 (6) SOURCE: "A Comparison of Kinematic Flood Routing Methods", by Fred Biesenthal (A Master of Engineering Thesis,) McMaster University Hamilton, Ontario

#### RESERVOIR ROUTING

PURPOSE: This subroutine generates the functional rating curves used by "KINRUT" to route a flood through a reservoir. A call of "KINRUT" performs the actual routing.
 METHOD: For a detailed description of the method used to route the flood refer to document-

ation of "KINRUT" and "KINFUN" as well

as the source. The rating of the reservoir

is given by the equation

STOR = FPVK\*QOUT\*\*FPVW

WHERE: STOR = Storage

FPVK = A constant

QOUT = Flow rate

FPVW = A constant

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- (3) PROGRAM:
- (a) DECK NAME; RESVOR
- (b) CALLING SEQUENCE: CALL RESVOR, (FPVK, FPVW, FLOWAR,

NFLOW, DTIM, NQIN, QINAR, QOUTAR)

WHERE: FPVK = A floating point variable

containing the value of the

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constant K.

- FPVW = A floating point variable containing the value of the constant W.
- FLOWAR = A floating point array containing the flow rates to be used when the functional relationships are to be calculated.
- NFLOW = An integer value which specifies the number of
- DTIM = A floating point variable

which specifies the time

step used in the calculations.

values in the vector FLOWAR.

(Same units as time unit of

flow rates.)

- NQIN = An integer value which specifies the number of points in the flow hydrographs.
- QINAR = A floating point array which contains the flow rates of the inflow hydrograph specified at equal time intervals defined by DTIM.

- QOUTAR = A floating point array which contains the flow rates of the outflow hydrograph specified at equal time intervals as defined by DTIM.
- (c) OUTPUT: The routed flow is returned to the calling sequence via QOUTAR. A plot of the inflow and outflow hydrographs is provided by a call of subroutine "HPLOT".
- (d) RESTRICTIONS: To provide a dimensionless representation

   of the hydrograph, a floating point value is
   transferred to the subroutine via a COMMON
   statement labelled "FLOW". The flow's
   were divided by a flowrate so that the result
   was a value of Q where:

 $| \ge Q \ge 0$ 

(e) OTHER DECKS REQUIRED: KINRUT, HPLOT
(4) EXAMPLE: An example of the use of this program is available from the source.
(5) DISCUSSION: The main purpose of this subroutine is to route a flood through an imaginary reservoir. If it is more advantageous to utilize reservoir rating curves which are storage as a function of elevation and outflow as a function of elevation see RESRUT.

(6) SOURCE: "A Comparison of Kinematic Flood Routing Methods", by Fred Biesenthal (A Master of Engineering Thesis)
McMaster University Hamilton, Ontario (1) PURPOSE: This program is designed specifically for use in a time sharing mode and is intended to set up, calibrate and subsequently modify a numerical model of a natural river channel under steady and unsteady flow conditions.

> A stretch of river is described by a number of cross-sections each of which is defined by a series of points. the initial co-ordinate values of which are referred to arbitrary datums for level and horizontal distance. Each cross-section is identified by a fixed chainage (and section number), and is also assigned an initial value of roughness coefficient. The resistance to flow may be defined by any one of a number of laws; the choice being made during the run. The program operates on the system thus defined and for specific values of discharge and downstream rating curve computes the water surface elevation(s) and energy level(s) at the cross sections specified during execution.

291

In the course of execution the user has the option to vary the discharge along the channel, the number of profiles calculated. the location and rating curve of the downstream control. the portion of the river over which the profile(s) is to be computed, the print out of the profile data, the roughness coefficient and the geometry of any selected section or any combinations of a the options mentioned above. In addition, there are several other features of the program devoted to steady state analysis which enables the user to calculate the critical flow depth at a selected section, list any changes that have been made to the data or to print the existing cross section data.

An initially defined system may therefore be adjusted to correctly reproduce an observed flow profile and then be used to examine the effect of a chosen design flow and to experiment with changes to the cross section geometry. The profiles that are calculated may be utilized in the unsteady analysis described below,

The analysis of water profiles incorporates a method of handling transition sections, such as weirs or bridges, which may occur along the channel.

The unsteady analysis incorporated into this subroutine allows a hydrograph of varying flow to be routed down portions of the channel using kinematic wave theory. An option is available to allow a user to route the flood through a reservoir which may be incorporated within the channel. Data which describes the hydrograph and other variables such as the time step for computation purposes are all entered during execution of the program.

METHOD: The calculation consists essentially of repeated applications of the subroutines
 EZRA and CONTRO, starting from the farthest downstream reach and proceeding upstream for each of the profiles.
 Flood routing is performed using kinematic wave theory as formulated in the subroutine KINRUT

- (3) PROGRAM:
- (a) DECK NAME: RIVER3
- (b) CALLING SEQUENCE: RIVER3 (HORZ2D, VERT2D, NPTSAR, CHAINR, RCAR, NXSEC, MAXPTS, G, NR1, NR2, NW, NF, AREAR, ELEVAR, FLOW2D, WLAR, QDAR, FUSAR, GDSAR, QUSAR, QDSAR, QRAR.)

WHERE: HORZ2D =Two-dimensional array containing the horizontal coordinates for the points describing the series of cross-sections, The first subscript of the array represents the section number and the second subscript represents one of the horizontal coordinates. VERT2D =Two-dimensional array containing the vertical coordinates for the points describ-

> ing the series of cross-sections. The first subscript of the array represents the section number

## RIVER3/5 295

and the second subscript represents one of the

vertical coordinates.

NPTSAR = One- dimensional array containing the number of coordinate points for each cross-

section in the series of sections.

- CHAINR = One-dimensional array containing the chainage values for the series of crosssections.
- RCAR = One-dimensional array containing the roughness co-

efficients for each cross -

section in a series of sections.

- NXSEC = The number of cross-sections in the river channel.
- MAXPTS = The maximum number of coordinate points required to describe any cross-section for the series of sections.
   G = Gravitational acceleration

constant.

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RIVER3/6 296

- NR1 = Periph. Device No. for Geometry file.
- NR2 = Periph. Device No. for keyboard input.
- NW = Periph, Device No. for terminal output,
- NF = The number of points in the working arrays and tables define cross-section prop-

erties versus flow rate.

AREAR = Two-dimensional array containing cross-section area for each section and flow used to describe the steady state performance. Dimensioned NF\*NXSEC. ELEVAR = Two-dimensional array

> containing water surface elevation for each section and flow used to describe the steady state performance. Dimensioned NF\*NXSEC.

RIVER3/7

297

- FLOW2D = Two-dimensional array containing the flowrates at each section used to calculate the steady state performance.
- WLAR = Vector array which contains the water level coordinates for the rating curve of the downstream section.
- QDAR = Vector array which contains the discharge coordinates for the rating curve of the downstream section.
  - FUSAR = Vector array which contains the function coordinates of the rating curve for the upstream section used to perform the kinematic flood routing.

GDSAR = Vector array which contains the function coordinates of the rating curve for the downstream section used to perform the kinematic flood routing.

RIVER3/8

- QUSAR=Vector array which containsthe flow coordinates of the<br/>rating curve for the upstream<br/>section used to perform the<br/>kinematic flood routing.QDSAR=Vector array which contains<br/>the flow coordinates of the<br/>rating curve for the downstream<br/>section used to perform the<br/>kinematic flood routing.
  - QRAR = Vector array which contains the flow coordinates of the rating curve used for reservoir routing.
- (c) OUTPUT: The output of the subroutine consists of invitations to enter data, such as questions requesting input, as well as the results of the various calculations. Generally, information provided consists of section number, chainage, water surface elevation and energy level for steady state analysis and section number and time and flow rates for a hydrograph calculated during unsteady flow analysis,

298

(d) RESTRICTIONS: The flow resistance law must be chosen from

aniong the following: Chezy, Manning. Strickler, Colebrook-White, Nikuradse's logarithmic Smooth Turbulent or Nikuradse's logarithmic Rough Turbulent. It is the user's responsibility to ensure that the selected law is appropriate both to the river system and to the roughness coefficient contained in the geometric data.

If the smooth turbulent law is used, the roughness coefficient is ignored but arbitrary data must still be provided.

The number of cross-sections cannot be varied from that defined in the data file. Neither can the initially defined maximum number of points per section be exceeded. The discharge is assumed to be uniform along each sub-reach. However. a number of sub-reaches may be defined within the river system. When performing the calculations a sub-critical flow regime is assumed.  (e) OTHER DECKS REQUIRED: READXS, SELSEC, PROPS, CRITIC, BOTTOM, SFROMQ, CHEZYQ, MANNGQ, STRICQ, COLEQ, SMOTHQ, ROUGHQ, EZRA, CONTRO, KINRUT

- (4) EXAMPLE:
- (a) INPUTS: A rectangular channel 50,000' long, 100' wide and 20' deep is subject to a triangular inflow hydrograph. Eleven sections with with four points each were utilized to describe the channel which had a bed slope = 0,0002. The Mannings roughness coefficient was n = 0.0149. The geometrical data was stored on a file and was read in as TAPE1. For each section the data comprises:
  (i) Section number (sections must be
  - numbered sequentially from No. 1 at the upstream end, but need not be read in that order.)
  - (ii) distance
  - (iii) roughness measure
  - (iv) Number of points used to describe

the section.

:

(v) A sequence of pairs of coordinates defining the cross-section geometry.
A copy of TAPEI is shown in Appendix G of the source. The format of this input file is dictated by the format used for the read statements in subroutine READXS.
Other information relating to discharge, control level, inflow hydrographs and subsequent system changes etc., is input from the console during the run.
Each input is detailed in the sample output which is largely self-explanatory.
A sample output is shown in Appendix G

- (b) OUTPUT: A sample output is shown in Appendix G of the source.
- (5) DISCUSSION: New users should refer to documentation of all related programs and routines. Further details of the method used for analysis of unsteady flow are available from the source.

(6) SCURCE: Modified from RIVER2 (Dr. A. A. Smith.

McMaster University, Hamilton, Ontario.)

"A Comparison of Kinematic Flood Routing

Methods". by Fred Biesenthal

(A Master of Engineering Thesis)

McMaster University

Hamilton, Ontario

## APPENDIX E

# DERIVATION OF THE LAX-WENDROFF METHOD

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The continuity equation is:

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$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} = \bar{q} \qquad (E.1)$$

This equation can be written in the conservation form:

$$\frac{\delta Q}{\delta x} + \frac{\delta A}{\delta t} - \bar{q} = 0 \qquad (E.2)$$

Expanding  $A(X, T + \Delta T)$  using a Taylor Series, the

following is obtained:

$$A(X, T+\Delta T) = A(X,T) + \Delta T \frac{\delta A}{\delta t} + \frac{\Delta T^2 \delta^2 A}{2 \delta t^2} + O(\Delta T^3) \qquad (E.3)$$

.

From equation E.2

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$$\frac{\delta A}{\delta t} = -\left(\frac{\delta Q}{\delta x} - \bar{q}\right) \tag{E.4}$$

and

$$\frac{\delta^2 A}{\delta t^2} = -\frac{\delta}{\delta x} \left( \frac{\delta Q}{\delta t} \right) + \frac{\delta \bar{q}}{\delta t}$$
(E.5)

Assuming a singled valued relationship between Q and A:

$$\frac{\delta Q}{\delta t} = f(A) \frac{\delta A}{\delta t}$$
(E.6)

Substituting equations E. 6 and E. 4 into equation E. 5 yields:

$$\frac{\delta^2 A}{\delta t^2} = \frac{\delta}{\delta x} \left( f(A) \frac{\delta Q}{\delta x} - f(A) \overline{q} \right) + \frac{\delta \overline{q}}{\delta t}$$
(E.7)

Therefore:

$$A(X,T+\Delta T) = A(X,T) - \Delta T \left(\frac{\delta Q}{\delta x} - \bar{q}\right) + \frac{\Delta T^{2}}{2} \left(\frac{\delta}{\delta x} \left(f(A)\frac{\delta Q}{\delta x} - f(A)\bar{q}\right) + \frac{\delta \bar{q}}{\delta t}\right) \quad (E.8)$$

:

Using the space-time grid shown in figure 3.2 equation E.8 can be expressed in terms of finite differences.

Figure 3.2 is reproduced below for reference purposes.



The equation in finite difference form is

$$A(I, J+I) = A(I, J) - \Delta T \left( \frac{Q(I+I, J) - Q(I-I, J)}{2\Delta X} - \frac{1}{2} (\bar{q}(I+I, J) + \bar{q}(I-I, J)) \right) + \frac{\Delta T^2}{2\Delta X} \left( f(A) \frac{Q(I+I, J) - Q(I, J)}{\Delta X} - \frac{f(A)}{2} (\bar{q}(I+I, J) + \bar{q}(I, J)) - f(A) \frac{Q(I, J) - Q(I-I, J)}{\Delta X} + \frac{f(A)}{2} (\bar{q}(I, J) + \bar{q}(I-I, J)) \right) + \frac{\Delta T}{2} \left( \bar{q}(I, J+I) - \bar{q}(I, J) \right)$$
(E.9)

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It should be noted that the value f(A) is a unique value for each position on the space time diagram. That is,

$$f(A(X,T)) = \frac{Q(X,T)}{A(X,T)}$$
 (E.10)

Substituting equation E. 10 into equation E. 9 gives:

$$A(I, J+I) = A(I, J) - \Delta T \left( \frac{Q(I+I, J) - Q(I-I, J)}{2\Delta X} \right)$$

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$$-\frac{1}{2}(\overline{q}(I+I,J)+\overline{q}(I-I,J))) + \frac{\Delta T^{2}}{2\Delta X}\left(\frac{Q(I+I,J)^{2}}{A(I+I,J)\Delta X}\right)$$

$$-2\frac{Q(I,J)^{2}}{A(I,J)\Delta X} + \frac{Q(I-I,J)^{2}}{A(I-I,J)\Delta X} - \frac{Q(I+I,J)\overline{q}(I+I,J)}{2A(I+I,J)}$$

$$+ \frac{Q(I-I,J)\overline{q}(I-I,J)}{2A(I-I,J)} + \frac{\Delta T}{2} \left(\overline{q}(I,J+I) - \overline{q}(I,J)\right) \quad (E.II)$$

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After deriving A(I, J+1), Q(I, J+1) is obtained from the relationship between flow rate and area. This function is single valued for kinematic waves.

Using the above formulation, the solution advances downstream on a particular time level. Succeeding time levels are considered in following passes down the channel.

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#### APPENDIX F

#### KINEMATIC FLOOD ROUTING--METHOD OF CHARACTERISTICS

Kinematic waves have one set of characteristics which travel downstream with a velocity,  $\bf C$ , where

$$C = \frac{dQ}{dA}$$
 (F.1)

Q = flow rate

$$A = area$$

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To route a flood wave through a channel, using a kinematic method, it is necessary to calculate only the wave velocity for various flow rates, determine how much time is required for the particular flow rate to move through the length of channel and thus plot points which determine the time history of the outflow hydrograph.

Determination of

Using Mannings Formula:

$Q = \frac{1.49}{n} A R^{\frac{2}{3}} \sqrt{S}$	(F.2)
$=\frac{1.49}{n}\sqrt{S}\frac{A^{\frac{5}{3}}}{P^{\frac{2}{3}}}$	(F.3)

$$K = \frac{1.49}{n} \sqrt{S}$$
 (F.4)

FOR A RECTANGULAR CHANNEL:

Y = DEPTH

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$$\langle = \frac{1110}{n} \sqrt{S}$$
 (F.4)

$$K = \frac{1110}{n} \sqrt{S}$$
 (F.4)

$$A^{5/3} P^{-2/3} = (T_w X Y)^{5/3} (T_w + 2Y)^{-2/3}$$
  
T = TOPWIDTH

$$\frac{dQ}{dA} = \frac{I}{T_w} \frac{dQ}{dY}$$
(F.6)

$$\frac{dQ}{dY} = \frac{d}{dY} \left( K (T_W Y)^{5/3} (T_W + 2Y)^{-2/3} \right)$$
 (F.7)

$$\frac{dQ}{dY} = K \left( \frac{5}{3} \frac{T_w^{\frac{5}{3}} Y^{\frac{2}{3}}}{(T_w + 2Y)^{2/3}} - \frac{4}{3} \frac{(T_w Y)^{5/3}}{(T_w + 2Y)^{2/3}} \right)$$
(F.8)

$$\frac{dQ}{T_w dY} = \frac{K(T_w Y)^{5/3}}{T_w(T_w + 2Y)^{2/3}} \left( \frac{5}{3Y} - \frac{4}{3(T_w + 2Y)} \right)$$
(F.9)

$$\frac{dQ}{dA} = \frac{Q}{T_w} \left( \frac{5}{3Y} - \frac{4}{3(T_w + 2Y)} \right)$$
(F.10)

(F.5)

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The following program was written to calculate flow rate,

ration of flow rate to full depth flow rate and kinematic wave

velocity all as a function depth.

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100=		PROGRAM KIN(INPUT, OUTPUT, TAPE5=INPUT, TAPE6=OUTPUT)
110=		S=0.0002
120=		RC=0.0149
13Ø=		G=32,2
140=		$W = 100 \cdot 0$
150=		DMAX=20.0
160=		DH=DMAX/40.
170=		P=DMAX*2.0+100.0
180=		A=W*DMAX
190=		CALL MANNGQ(A, P, S, RC, G, QMAX)
200=		H=0.0
210=		DO 10 1=1,40
220=		H=H+DH
23Ø=		A=H*W
240=		P=W+2.0*H
250=		CALL MANNGQ(A, P, S, RC, G, Q)
26Ø=		VK=Q/W*(5.Ø/(3.Ø*H)-4.0/(3.Ø*W+6.Ø*H))
270=		QR=Q/QMAX
28Ø=	10	WRITE(6,100) H, Q, QR, VK
290=	100	FORMAT(1X, 4F12.4)
300=		STOP
310=		END

The results are shown in table F.1

#### TABLE F-1

### FLOW RATE, FLOW RATIO AND KINEMATIC WAVE VELOCITY

#### AS A FUNCTION OF DEPTH SYSTEM I

DEPTH	FLOW RATE	Q/Qfull bank	С
(ft)	(cfs)		(fps)
.5000	44.3121	• 0027	1.4712
1.0000	139.7612	.0084	2.3111
1.5000	272.9276	.0164	2.9972
2.0000	438.0072	• Ø263	3.5939
2.5000	631.2877 /	•0379	4.1284
3.0000	850.0642	.0510	4.6157
3.5000	1092.2224	•0655	5.0650
4.0000	1356 • 0353	.0813	5.4827
4.5000	1640.0486	.0983	5.8736
5.0000	1943.0102	.1165	6.2412
5.5000	2263.8242	•1358	6.5881
6.0000	2601.5186	.1560	6.9167
6.5000	2955.2223	.1772	7.2288
7.0000	3324 • 1 48 1	•1993	7.5258
7.5000	3707.5789	•2223	7.8092
8.0000	4104.8584	•2462	8.0800
8.5000	4515.3820	•2708	8+3391
9.0000	4938.5908	•2962	8.5875
9.5000	5373.9658	• 3223	8.8259
10.0000	5821.0237	• 3491	9.0549
10.5000	6279.3127	• 3766	9.2752
11.0000	6748.4098	• 4047	9•4873
11.5000	7227.9175	• 4334	9.6917
12.0000	7717.4617	• 4628	9.8889
12.5000	8216.6895	• 4927	10.0791
13.0000	8725.2675	•5232	10.2629
13.5000	9242.8803	•5543	10.4406
14.0000	9769.2287	•5858	10.6124
14.5000	10304.0287	•6179	10.7787
15.0000	10847.0104	•6505	10.9397
15.5000	11397.9171	•6835	11.0957
16.0000	11956.5040	.7170	11.2470
16.5000	12522.5375	•7509	11.3936
17.0000	13095.7949	•7853	11.5360
17.5000	13676.0631	.8201	11.6741
18.0000	14263.1384	•8553	11.8083
18.5000	14856.8258	•8909	11.9386
19.0000	15456.9383	•9269	12.0653
19.5000	16063.2967	•9633	12.1885
20.0000	16675.7291	1.0000	12.3083

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This data was used to plot a curve of flow rate versus wave velocity. (Figure F.1) By dividing wave velocity for a particular flow rate into the length of the channel in feet, the lag between inflow and outflow of that flow rate can be determined. Points on an inflow hydrograph can be transposed downstream and the outflow hydrograph can be determined.

System 1 was analyzed using this method. The results of the analysis are given in table F.2 and the hydrograph is plotted in figure 3.8.

## TABLE F-2

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RESULTS OF KINEMATIC FLOOD ROUTING SYSTEMI

				the second se
Inflow Ratio	Velocity (ft/sec)	Travel Time (sec)	Inflow Time (sec)	Ou <b>tflow</b> Time (sec)
 0.2	7,52	5319	1500	6820
0.4	9.46	4228	2750	6980
 0.6	10.73	, 3728	4000	7730
0.8	11.60	3448	5250	8700
1.0	12.31	3249	6500	9750
 0.8	11.60	3448	7750	11200
0.6	10.73	3728	9000	12730
0.4	9.46	4228	10250	14480
0.2	7.52	5319	11500	16820
		1	8	1

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#### APPENDIX G

LISTINGS OF COMPUTER INPUT FILES, ROUTINES AND OUTPUT

INPUT FILES:	FINDIF
	KINDIF
	TAPE1
	ANCDAT
	WHTDAT
ROUTINES:	FINDIF
	KINDIF
	KINFUN
	KINRUT
	HPLOT
	RESVOR
•	RIVER3
OUTPUT:	FINDIF
	KINDIF
-	RIVER3

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TYPICAL INPUT FILE: FINDIF

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100=H986.T400. BIESENTHAL 110=ATTACH.RIVEPF.ID=HPAG.MP=1. 120=ATTRCH.CIVLIB.ID=HRAG.MR=1. 130=FTH(I=RIVERE) 140='DOET(LIB=CIVLIB) 150=660. 151=+EDE 170=50000. 100. 0.0149 0.0002 20. 24000. 180= 250.5 190=190 200 = 1 1210= 21 220= 1 21 230= 1 21 240= 3 250=0.2 1500. 260=1.0 6500. 11500. 270=0.2 271=+EDP 272=+EDF

## TYPICAL INPUT FILE: KINDIF

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100=HP95.T400.
BIEGENTHAL
110=ATTRCH.KINDIF.ID=HRAG.MP=1.
120=ATTACH.CIVLIE.ID=HPAG.MR=1.
130=FIN(I=KINDIF)
140=LDSET(LIP=CIVLIB)
150=160.
151=*ED2
                                 0.0002 20.
170=50000.
              100.
                      0,0149
                24000.
180= 20200.
190=190
200= 1 1
220= 1 17
230= 1 17
240= 3
              1500.
250=0.2
              6500.
260=1.0
              11500.
270=0.2
271=+EDR
272=+EDF
```

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## GEOMETRIC DATA FILE : TAPEI

1 0.0.0149 4 0.00130,00 0.00110.00 109.00110.00 100.00130.00 2 5000.0 .0149 4 0.00129.00 0.00109.00 100.00109.00 100.00129.00 З 10000.0 .0149 1 0.00128.00 0.00168.00 106.00108.00 100.00128.00 15000.0 .0149 4 4 0.99107.00 100.00107.00 100.00127.00 0.00127.00 5 20000.0 .0149 - 4 0.00106.00 100.00106.00 100.00126.00 0.00126.00 25000.0 .0149 6 <u>t</u>: 0.00105.00 100.00105.00 100.00125.00 0.00125.00 7 30000.0 .0149 4 0.00104.00 100.00104.00 100.00124.00 0.00124.00 8 35000.0 .0149 - 4 0.00123.00 0.00103.00 100.00103.00 100.00123.00 9 40000.0 .0149 4 0.00122.00 0.00102.00 100.00102.00 100.00122.00 10 45000.0 .0149 4 0.00121.00 0.00101.00 100.00101.00 100.00121.00 11 50600.0 .0149 4 9.90120.00 0.00100.08 100.00100.00 100.00120.00

## INPUT FILE: ANCDAT (GEOMETRIC DATA FOR APPLICATION TWO)

			* .					
$\int$	$\begin{array}{c} 0 & 0 & 0 & 0 & 0 & 0 & 10 \\ 0 & 0 & 720 & 0 & 32 & 0 & 7 \\ 109 & 0 & 707 & 0 & 166 & 0 & 7 \\ 0 & 0 & 707 & 0 & 166 & 0 & 7 \end{array}$	15.0 55.0 710.0 10.0 192.0 715.0	91.0 707.0 312.0 720.0	95.0 703.0	105.0 703.000	0100		
	0.0705.01007 54.0592.06000	20.0 695.0 83.0 700.0	36.0 692.0 120.0 705.0	40.0 688.0	50,0 688,000	0140		
≻	5.5 710 0 5.5 7 5.5 658 0 110 6	0000206980 89.5 10.8 691.0			<u>9.0 689</u> 500	0170		ł
	4 320,0,0150 13 5.5 710.0 5.5 6 5.5 687.0 11.0 6	29.0 2.0 697.0 88.5 10.8 690.0	1.1 695.0	9.0 697.0	00 0.0 688.500 5.6 699.000	0200		
	5 5 7 1 0 0 5 3 2 0 0 0 5 0 0 1 0 0 0 7 5 0 3 6 0 7	100.0 45.0 695.0	57.0 690.0	69.0 687.0	00 00 80.0 687.000	0230		
	80,0 890,0 80,0 6 382,0 0500 10 0,0 705,0 23,0 7	95.0 <u>80.0</u> 700.0 20.0 35.0 695.0	45.0 690.0	89.0 685.0	101.0 685.000	0260		
	105,0 690,0 115,0 6 3H2,0 0150 16 0,0 705,0 23,0 7	95.0 115.0 700.0 700.0 35.0 695.0	115.0 705.0 38.0 693.5	95.0 693.4	00 95.0 692.500	0290		
_ ,	8936924 8936 11206935 11506 400000500 10	85.0 100.8 665.0 95.0 115.0 700.0	100.8 692.4 115.0 705.0	95,1 692,5	95,1 693,400 00 00	0320		
	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	00.0 35.0 695.0 95.0 115.0 700.0	45.0 690.0 115.0 705.0	89,0 685,0	101,0 685,000 00 00	0350		
1	0.0 655.0 23.0 6 65.0 673.0 80.0 6 600.0 0150 16	80.0 35.0 675.0 75.0 110.0 680.0	36.0 673.0 140.0 685.0	45,0 668,0	54.0 668.000 00 00	0380	2 W V	
		80.0 35.0 675.0 68.0 54.0 668.0 75.0 110.0 680.0	36,0 674,0 54,0 673,0 140,0 685,0	49.5 673.9	49,5 673,100 49,5 673,900 00	0410	~	
	0.0685.023.06 65.0673.080.06	80.0 35.0 675.0 75.0 110.0 680.0	36.0 673.0 140.0 685.0	45,0 668,0	54.0 668.000	0440	······	
	$\begin{array}{c} 0.0 & 672.0 \\ 72.5 & 660.0 \\ 100.0 & 6 \\ 72.5 & 660.0 \\ 100.0 & 6 \\ 100.0 & 100.0 &$	70.0 40.0 665.0 65.0 142.0 670.0	42.5 660.0	45.0 655.0	55.0 655.000	0480		
•	0.0 672.0 20.0 6 53.2 657.0 57.2 6 70.4 654 5 70 9	70.0 40.0 665.0 57.0 57.2 659.4 64.4 70.9 664.0	50.0 664.9 60.4 659.5 78 9 664.0	50.0 659.5 60.4 664.9	53.2 659.400 70.4 665.000	0510		
1	60 5 665 0 100 6 1050 0 0200 8 0.0 670 0 25.0 6	65.1 142.0 670.0	142.0 672.0	04 0 657 0		0540		
	90 0 665 0 140 0 6					0570		-
1		20.0 70.0 616.0	80.0 616.0	85,0 620,0	97.0 625.000	0600		
1	0,0615,03,56	10.0 6.0 609.0	16.0 609.0	18,5 610,0	22,0 615,000	0620		
	40, 4.610, 043, 7.6 45, 1.614, 643, 7.6	15.0 45.0 614.6 09.0 47.1 610.0 15.0 90.0 615.5	42.4 613.6	40.3 612.6		10650 10650 10670	· · · · · · · · · · · · · · · · · · ·	-
2	1500,0,0500,0 0,0,615,0,35,0,6 1860,0,0500,6	42,5 607,0	47.5 607.0	55.0 610.0	90.0 615.000	0680		
2	0.0 610,0 501,0 6 2370,0 0500 8 0,0 605,0 240,0 6	05.0 503.0 602.0 00.0 314.0 597.0	510,0 602,0 315.0 595.0	512,0 605,0 325.0 595.0	931.0 610.000 00 326.0 597.000	0710		
2		05,0 314,0 595 0	314.0 593.5	317.0 598 4	117 0 593 000	0740		/
	323 5 593 0 323 5 5	93 4 327 0 593 5	327 0 595 0	410 0 500 C	500,0 605 000	0770		

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#### **INPUT - ANCDAT/2**

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### INPUT - ANCDAT/3

5 15800 D	0700 17 -				•		· ·
-110.0 285.	0 5.0 280.0	12.0 275.0	25.0 271.0	95.0 270.0	100,0 265,00	01450	
425 0 270	0 445 0 271 0	455.0 275.0	457 0 280 0	463.0 285.0		01470	•
440 0 260	0 18 0 275 0 0 450 0 265 0	25.0 270.0 505.0 266.0	360.0 266.0	410.0 265.0	420.0 260.0Ŭ	01490	
0.0 280.	0700 12	22.0 270.0	32.0 265.0	388.0 263.0	398.0 258.00	01510	
	1500-10 263.0	532.0 265.0	722,0 270,0	763.0 275.0	790,0 280,00	01530	
431 0 260	0 484 0 265 0	530.0 270.0	591.0 275.0	413.0 25/.0	428,0 257,00	01550	
323.0 260	0 20 0 270 0 0 345 0 265 0	55.0 265.0 448.0 270.0	298.0 260.0	303.0 256.0	318.0 256.00		
•730, 331,	30300 6 285,	0,0 256,	4,5 256	4,5 285,	250 331,0	01600	
-730, 331,	03000 8 285.	0,0 256,	4,5 256,	4,5 285,	250, 331,0	01620	
375.0 260.	0 20.0 270.0	47.0 265.0.	335.0 260.0	332.0 254.0	355.0 254.00	01640	·· -
19070 0	0700 12 270.0	26.0 265.0	344.0 260.0	389.0 257.0	397.0 253.00	01670	
420 0 253	0 422 0 257 0	424.0 260.0	426 0 265 0	428 0 270 0	430 0 275 00	01690 01700	
19300.0	0 0.0 280. 10310 6 280	0.0 253.	6.0 253.	6,0 280,	130, 300,0	01710	
19300.0	100010	55 0 250 0	6,0 253, 403 0 258 0	0,0 280, 408 0 252 0	130, 300,00	01750	
427 0 258	0 430 0 260 0 1500 12	440.0 265.0	455.0 270.0	400,0 232,0		01760	
28.0 251	30,0 255,0	31.0 259.0	290.0 259.0	10.0 255.0 361.0 265.0	375.0 273.00	01780 01790	
77.0 270	0700 11 0 19.0 265.0 1 81.0 258.0	29.0 260.0	437.0 258.0	481+8 358+1	69.0 250.00	01810	
22075.0	0650 15 265.0	31.0 260.0	53.0 255.0	115.0 253.0	156.0 252.50		
159 0 248	5 165 0 248 4 0 669 0 265 0	171 0 248 5 595 0 270 0	174 0 252 5	215 0 253 0	618.0 255.00		·
	0000 15	29.0 257.0	42.0 255.0	42.0 320.1	65.0 250.00	01870	
119 0 257	0 390 0 260 0 0600 13	420.0 265.0	0340 53040	114,0 230,1	01	01900	
0.0 265	0 12 5 260 0 2 73 5 245 3	25.0 255.0	37.5 250.0 96.5 250.0	109.0 255.0	121,5 260,00	01920	
237.50.0	0600 13 260-0	25.0.255.0	37.5 250 0		0	01940	
07.0 246	ž 73.5 ž46.3	76.5 249.0	96.5 250.0	109,0 255,0	121,5 260,00	01970	<b></b>
23700.0	0600 13 260.0	25.0 255.0	37.5 250.0	57.5 249.0	60.5 246.30	00050	
134 0-265	<u>c 73,5 246,3</u>	10.5 249.0	96,5 250,0	109.0 255.0	121,5 260,00	02020	
0.0 262	0 22.0 251.0	28.0 248.0	47.0 247.0	49.0 245.0	60,0 245,00	02040	

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# INPUT FILE: WHTDAT (GEOMETRIC DATA FOR APPLICATION ONE)

: ]

1 5 1.00 / 2220.0	000100
g.g113.y 6.g 96.5 56.0 86.5 65.0 90.0 90.0115.0	000200
-11.010.0-10.0 89.6 0.0 89.6 0.0 80.0 49.0 80.0 54.0 88.2 97.0 94.0 97.0110.	0000550
$\begin{bmatrix} 3 & 7 & 1 & 50 & 50 & 40 & 0 \\ - & -1 & -0 & 9 & 1 & 0 & 20 & -0 & 72 & -0 & -70 & -70 & -90 & -73 & -9 & -80 & -80 & -84 & -0 & -84 & -98 &$	000230
4 8 1.50 5940 h	000250
	000270
	000220
$\begin{bmatrix} 60.0 & 95.0 & 60.0 & /3.4 & -1.0 & 76.7 & -1.0 & 67.0 & 68.0 & 71.0 & 73.5 & 11.0 & 96.0 \\ 7.6 & .20 & 72.60.0 & \end{bmatrix}$	000310
0.0 9 <sup>2</sup> .0 0.0 57.0 46.0 68.0 71.0 14.0 86.0 81.0 86.0 81.0 95.0	00017 <u>0</u>
0.0 95.0 0.0 70.0 30.0 67.0 59.0 69.0100.0 76.0100.0 95.0	000340 000350
-10.0 94.0 20.0 55.0 60.0 64.0 80.0 67.0 80.0 94.0	
1 0 0 95 0 0 0 7 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	000340
11 7 .50 11800.0 0.0 9 0 0 0.0 77.0 28.0 64.0 28.0 58.0 -78.0 -58.0 88.0 71.0 88.0 97.0	
12 8 .50 11800.0 20.0 90.0 20.0 72.0 0.0 64.0 0.0 58.0 50.0 58.0 50.0 64.0 30.0 72.0 30.0 90.	000410
13 7 .50 11800.0 0.0 90.0 0.0 77.0 28.0 64.0 28.0 58.0 78.0 58.0 88.0 71.0 88.0 90.0	000430
147 +50 12500 0 0.0 90.0 0.0 68 0 19.0 57.6 29.0 56.5 69.0 58.0 84.0 65.0 94.0 90.0	ÖÖU450 000460
15 5 0.90 13950.0 15 5 0.90 13950.0	000470
	000490
	000500
18.6 .10 15340.0	000530 000530
0.0 83.0 0.0 70.0 50.0 56.5 65.0 51.0 85.0 56.5 85.0 85.0 19 6 .10 16500.0	000540 000550
0.0 85.0 0.0 70.0 10.0 53.5 75.0 53.5 80.0 70.0 80.0 85.0	
-8.0 74.0 12.0 50.0 12.0 49.0 54.0 49.0 70.0 52.0 80.0 57.0100.0 62.0100.0 74.	nôn0540
0.0 80.0 0.0 49.2 62.0 49.2 62.0 80.0	000600
-15.0 7.0 12.0 49.5 12.0 48.4 54.0 49.0 70.0 52.0 80.0 57.0100.0 62.0100.0 75.	000420
0.0 75.0 4.0 50.0 10.0 47.5 60.0 47.4 90.0 56.5100.0 75.0	000640
24 8 75.0 0.0 59.5 15.0 49.0 24.0 45.5 A4.0 44.5103.0 55.0116.0 56.0140.0 75.	000440
25 7 .10 22850.0 0.0 72.0 0.0 60.0 11.0 52.0 11.0 41.5 95.0 41.5 95.0 60.0 95.0 72.0	000470 000480
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	000690 000700 -
27 8 .01 24000.0 0.0 70.9 0.0 49.0 14.0 45.0 30.0 44.0 84.0 44.0 84.0 47.0 96.0 47.0 97.0 70.	000720
	000730
	600750 0001750
30 6 10 26900.0	000770
U.U. IV.V. V.V. 47.4 /.C. 4/.0 19.0 38.0 74.0 38.0 /5.0 /0.0	000789

# INPUT - WHTDAT/2

31	6 .10 29000	0	000790
_15	4 10 29000	0 20.0 30.0 40.0 30.0100.0 50.0100.0 10.0	000810
33	· · · · · · · · · · · · · · · · · · ·	2 70.0 38.2 70.0 70.0	000830
=10 34		,0 20+0-38+2-90+0-38+2100+0-50+0106+0-70+0	000840 000850
-18 35	•0 67.0 20.0 41 7 •10 31000	6 27.0 38.0 72.0 38.0 86.0 42.0110.0 55.0110.0 65.0	000840 000870
36	•0 65 0 0 0 51 7 •10 32910	n 27.0 43.5 43.0 42.2 53.0 35.0100.0 33.0110.0 65.0	000HA0 000H90
-2	.0 65 0 -2.0 5n	n 24.0 37.5 35.0 35.6 81.0 35.6104.0 51.0104.0 65.0	
- Å	0 65 0 4 0 50.	n 22.0-34.0 70.0-34.0100.0 43.0120.0 45.0135.0 50.0135.	+0-65+0000+20
20	·0 67.0 0.0 50	n 5.0 43.5 30.0 40.5 34.0 35.5 92.0 35.5130.0 44.0130.	•0 65•0000440
<u></u>	.0 63.0 0.0 50	0 20.0 39.0 20.0 32.0100.0 32.0100.0 39.0110.0 50.0110	•0 65.c000360
-3	.0 6.0 -3 0 50	0 20.0 38.6100.0 38.6110.0 43.5110.0 65.0	000370
ื่่ง	Q 65 0 0 0 0 50	0 23.0 38.6 23.0 30.6103.0 30.6103.0 -38.6113.0-43.5113.	+0-65+0001000
10	.0 60,0 0.0 40	.n 5 25.0 32.3 30.0 29.0 68.0 29.0 80.0 42.0 80.0 60.0	001010
0	· C 60 0 0 0 44	n 17.0 29.0 58.0 29.0 70.0 37.8 85.0 41.5 85.0 60.0	
<b>4</b>	.0 60.0 0.0 40	.0 5 15.0 36.0 27.0 26.0 70.0 27.5 85.0 38.3 85.0 60.0	001060
- 45	6 .10 40800. 0 58.0 0.0 42	0 10.0 25.5120.0 25.5130.0 42.0130.0 58.0	001070
46	4 .04 40400	0 5 73.0 24.5 73.0 58.0	001090
47	6 05 40800	0 10.0 24.5120.0 24.5130.0 42.0130.0 58.0	
4 8 D	6 .05 42800 0 58.0 0.0 35		001130 001140
4 Š	8 US 44800		
5 ğ	6 .20 46400	$\begin{array}{c} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 $	001170
5 ĭ	6 10 47920		
52 52	8 20 47920		001210
53	6 20 47920	0 0.0 21.0 0.0 10.0 70.0 10.0 10.0 21.0 45.0 34.0 45.	001530 •• = 001530
54	8 49000	0 10+0 1/+2 40+0 1/+2 40+0	001240
-10	8 40 10 0 28. 8 40 50230	n 38.0 23.0 4/.0 16.0102.0 16.0110.0 21.5118.0 21.6155.	• 0 48 • 0001760 001270
56	·0 49.0 0.0 30. 7 .49 52465.	0 10.0 18.7 38.0 10.8 89.0 11.8 94.0 20.5120.0 30.0120.	•0 48•0001280 001290
58	0 40 0 0 0 30 4 20 54700	n 20.0 8.2100.0 8.5105.0 23.5116.0 23.5116.0 48.0	001300 001310
58	0 40.0 -0.0 11 5 01 54700	0170+0 11+0170+0-48+0	
Ō,	0 48.0 0.0 17	0 10.0 17.3175.0 17.3175.0 48.0	ŎŎĨ <b>Ĵ</b> ĂŎ

# LISTING OF PROGRAM FINDIF

<u>.</u>-' -

	PROGRAM FIN	DIF	73177	UPT=0	TRACE		FTN 4.0+P3	355	01/28/74	11+35+12+	PAGE	1
		PROG	RAM FIND	IF (INPUT	.OUTPUT.PUN	CH, TAPE5=IN	PUT.TAPE6=OUTF	PUT+TAPE7	P000100		•	
		IUNCH	) NSIUN H(2	200.2).0	(200.2) .WL (2	20) . TWL (20)	.X (200) +BOT (20	00)		•		
5		DIME	NSIUN HDI	(200-2)	PEHYD(2) + 18	EDP(2)	~ '		000140		•	
	ę	THIS	PHUGRAM	RUUTES	A FLOOD USI	NG AN EXPLI	CIT FINITE DIF	FERENCE				
10	ç	TECH	NIQUE	A DICTI	ONARY OF THE	VARIABLES	FOLLOWS.		-000180			
	<u>ç</u>		I' ELEVA	TTUNS OF	BANK FULL	NORMAL FLOW	T THE SECTIONS	5	- 000210			
15	ů č		X THE MARK	AXIMUM T	THE STEP AT	BANK FULL	CONDITIONS		000230		•	
15	Č	HÎI. NRE	J) WATER	LEVELS NUMBER	AT THE SECT	IONS IN THE CHAN	NEL		000250			
	Č	NSEC GII.	S NUMHFR J] FLOW	ALES AT	IONS IN THE	CHANNEL			000270			
		GIN GOUT	THE VOL	JME OF T	HE WATER WHICH	ICH FLOWED- FLOWED OUT	OF THE CHANNE		000290-			
	ů č	TIME	I START	TIME	CIENI MANN. Me	1105 1			000320			
25	. Č	TWL (	1) THE W	I UWRATE	S OF THE POL FLS OF THE	INTS IN A S POINTS IN	TAGE DISCHARGE	E CURVE	000340			
	č	- X(I) - XTO	DISTANCE	E FROM T TUTAL L	HE OUTFALL Ength of The	IN FEET			- 000360 - 000370 -			
30	3	READ	PHURLEM	VARIABL	ES				0003R0 000390			
		READ	(5+501) (5+501)	NRLACH+D	+RC+SLOPE+DE T+TIMÉF	EPTH	-		000410			
35		READ	(5+506) 1 (5+506) 1						000430		•	
		READ	(5+5061N) (5+5061N)	SFU1+NSE SFU1+NSE	02 D2				- 000450 000460			
		READ	(5+504) ( (5+505)	10215 (al(I)•T	I(I)+I=1+NQ	PTS)			000470			
φU	č	SET	PHOHLEN N	ANIABLE	S YD(1)#PFHYD	(2)=1.			000500			
		 — G=32	1=FM0M0=(	).0					- 000520 000530			
45		Δ=Ϊ+ P=T+	DEPTH 2. PUFPTH						000540			
	C	SET	UP THE IN	NFLOW HY	E+HC+G+GMAX) DROGRAPH	1			000570			
50	18	0 01(1	1 = QI(I) + ( 10TA: /FL	JMAX JMAX JA{{NRFA	CHI				000590			
		- NSEC QIN#	S=NHFACH	•1					-000610- 000620		· · ·	
55	_	NPWI IEND	=<0 (1)=[END	(2)=IEOP	(1)=IEDP(2)	•0			000630			
	•	TIMF TIMF	20. 1=0.						000050			

### LISTING-FINDIF/2

FRO	DGRAM FINDIF 73/73 UPT=0 TRACE	FTN 4-0+P355 01/28/74	· 11+35+12+ P	AGE 2
	CEL-SURTIGEDEDTH) +OMAX/(DEPTHET)	000670		
	DIMAX#DX/CFL	000680		
60				
	C DERTYE STAGE DISCHARGE CURVE	- 000710		
	WE(1)=0.	000730		
65	TWL(1)=0.0 DO.T. I=<.NDWI	000740		
	• • • • • • • • • • • • • • • • • • •	000760		
	A=T+FLOAT(7"1)*DH P=T+2+*FLOAT(7"1)*DH	000770		
70	CALL MANNGO (A. P, SLOPE, RC, G, QX)	000790		
	7 TWL(I)=0x	000810		
	C SET INTERAL CONDITIONS AT STEADY STATE	000820		
75	CALL INTERI (OT.TI.NOPTS.TIME.Q(1.2))	000840		
	B01(1)#XT0TAL#SLOPE+100+ X(1)=0.0	000550		
		000870		
80	H(1.1)=H(1.2)==0T(1)+FLOWD	000890		
	DRCP=UX*SLOPE DU K I=Z.NSECS	000900		
	Q([1,1]) = Q([1,2]) = Q([1,2])	000920		
	X(I)=X(I+1)+DX	000240		·····
	6 H(I.I)=H(I.C)=HOT(I)+FLOWD C CALCULATE INTIM STOPAGE IN THE CHANNEL	000950		
	STCPEL ATOTAL + 1 + FLOWD	000970		
90	C CALCULATE NOMBER OF STEPS BETWEEN PRINT OUT KSTEP=TIMEF/(119,*OT),+0,5	000330		
	IF (KSIEP.EA.O) KSTEP=1 TCOUNTEKSTEP	001000 001010		
	DTHYD(1)=UTHYD(2)=FLOAT(KSTEP)+DT/3600-	001020		
95	WRITE (6+2001) ATOTAL.DEPTH.SLOPE.RC.QMAX.CE			
-	WRITE (0:2008)   BRITE (0:2008)   BRITE (0:2008)	001050		•
	WRITE (0,2006) ITEST	001070		
100	WRITE(0+2004) G(1+2) DO 100 I=1.NSEUS	001080	•	
	100 WRITE (5,2005) x(1) +BOT (1) +H(T,2)			
	$IF(NPU \cdot LF \cdot n)  WRITE(6 \cdot 2012)$	001120		
105	IF (NPU2+1E+0) WRITE (6+2013)		-	
2 V J	TIME=TIME+DT/2.	001150		
	C START CALCULATIONS OF D	001170		
110	60 CALL IN FENT (QTITINOPTSITIME QUS)		*	
110	DO 20 IZZINSECS	001200		
	HA=(H(I+2)+H(T*)+2)-BOT(I)+ROT(I=1))/2. DHDX=(H(T+2)+U(T+)+2)}/(X(T)+X(T+1))	001210		

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### LISTING - FINDIF/3

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PROGRAM FINE	IF 73/77	UPT=0 TRACE	FTN 4.0+P355	01/28/74 11.	35,12. PAGE	3
112	FAC7=(H(I=1+2 FAC7=Q(I+1)*( FAC4=1+/(G*I*	) -H(I-1,1)+H(I+2)-H( RUT(I-1)-BOT(I))/(G' H4*0I)	(I,I))/(G*T*HA*HA*4.*DT) *T*T*HA**3.)/(X(I)-X(I-1))	001240 001250 001260		
120 20	+/3.)) (Q(1.2)=(OHOX= QIN=QIN+(Q(1) GOUT=QOUT+(Q( FMCMI=FMOMI+0	0(1,1)*(FAC2+FAC4))/ 1)+0(1+2)/2+0 N>ECS+1)+0(NSECS+2)) (NSECS+2)+TIME/3600.	/(FAC1+FAC2+FAC3=FAC4=FAC5) //2+	001280 001290 001290 001310 001320		
125 C	FMCMO=+MOMO+Q TRANS+EH INFO - DO 30 I=1+NSE	INSEQ2.2) +TIME /3600 A		001330		
c <sup>3(</sup> 130 C	H(I+1) = H(I+2) INCPEMENT TIM TIVF=IIMF+DT/ DO THE CALCUL CALI INIFHI(W h(N <ecs+2)=w\sf< td=""><td>F 7 4 AIIONS FOR H 1 + TWL+NPWL+Q (NSECS+2 75+100+0 75</td><td>2),WLDS)</td><td>001360 001370 001380 001390 001400 001410 001420</td><td></td><td></td></ecs+2)=w\sf<>	F 7 4 AIIONS FOR H 1 + TWL+NPWL+Q (NSECS+2 75+100+0 75	2),WLDS)	001360 001370 001380 001390 001400 001410 001420		
135	K=NSECS+2=T J=K=1 IF(T+LQ+NSFCS IF(T+LT+NSFCS H(1+2)=H(1+2)	) DX=X(2)-X(1) ) DX=(X(K)-X(NSECS-) -UT/DX+(0(K-2)-0(1))	\${\/ <b>?</b> •	001430 001440 001450 001460 001470		
140 Č	PRINT THE VAL IF ITIME I T.O. IF (TCOUNT.NE)	HES IF APPROPRIATE 1 GO TO 80 KSTEP) GO TO-170		001480 001490 001500 001510		
145 7(	INTFR=1LND(1) ICCUNTEU HYD(INTED:1)= HYD(INTED:1)= HDP(INTED:1)= HDP(INTED:1)=	+1 (9(NSEQ1+1)+0(NSEQ1+ (7(NSEQ1+1)-80(NSEQ1+ (7(NSEQ1+1)-80(NSEQ1+ (7(NSEQ2+1)-80)(NSEQ	2))/(GMAX+2.) 2))/(GMAX+2.) 2))/(DEPIM 2))/DEPIM	001530 0015340 0015540 00015567 0015567		
150	IEND(1)=1FNU( MX=H(NSTAGF,1 QX=(Q(NSTAGE, SF=0X*0X*HC*R D0=/0(NSTAGE,	2) = IEDP(1) = IEDP(2) = ] ) = AOT(NSTAGE) 1) + Q(NSTAGE+2) + Q(NST C*(T+2.+HX) + (4./3.) C*(T+2.+HX) + (4./3.)	(EDP(2)+1 TAGE+1+1)+n(NSTAGE+1+2))/4. 1/(1+4941+49+(10,/3+)			
155	DV=nQ/(G+Hx+T CALI PLUTPT(Q IF(NPU2) 170, WRITE(7,3006)	+UT) X+HX+9) 1/0+150 HX+QX+SF+DV		001630 001640 001650 001650 001660		
C 17( 160 80 C 50	F ICCUNI=ICOUNT TRANSFER INFO DO 50 I=1+NSE 0(I+1)=0(I+2) INCREMENT TIM	TATION IN THE Q MAT	<b>TRIX</b>	001670 001680 001690 001700 001710		
165	IF (TIME .LE TI TIMFI#TIMET/3	MEF) GO TO 60	1 4 8 <del>1 1 </del>	001730		<u>.</u>

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### LISTING-FINDIF/4

PR	ROGRAM FINDIF 73/73 UPT=0 TRACE	FTN 4.0+P355	01/28/74	11:35,12:	PAGE	4
	D0 90 I=2+NSECS 90 STCREZ=STOREZ+M(I+2)-BOT(I)		001810 001820			
175	STCREZ#STOREZ+(H(1,2)-BOT(1)-H(NSE) STCREZ#STUREZ#UX#T	C5+2)+BOT (NSECS) ) /2+0	-001830			
- • •	QIN±QIN⇔nT QOUT=QOUT⊕nT	•	001850 001860			
	C CALCULATE CENTHOIDS OF THE HYDROGR	APHS	001870			
180		-100	801820			
	WRITE (6, 2003) EPROP		001910			
	IF (NPU) 120,120,110		001930			
185	110 WRITE(7+3001)DEPTH+T+SLOPE+RC IEK=IENU(1)		001940			
	00 160 1=1. TEN Hyd (1.1) - Hyd (7.1) * CMAX		001960			
100	HYD (1, 2) = HYD (1, 2) * QMAX		001940			
144	160 HOP (112) #HOP (172) *DEPTH		- 002000			
	NSED=NSED1		002020			
195	NSED=NSED1 DO 13V Jal.2		002040			
	WRITE(/*ROA2) J*BOT(J) Do 140 I±l.NPWL		002050	٥		
· · · · · · · · · · · · · · · · · · ·	A=T+FLOAT(T+1)*DH E=BOT(1)+W(1)		- 002070			
200	WRTTE(7.3003)E.A.TWL(1)		002000			
	WRITE (7.3004) J.NSEG.DTHYD(J) QMAX	+GX+IEND(J)	002110			
	WRITE(7+3005) (HYD(1+J)+I=1+IEN) WRITE(7+3008) (SED+DTHYD(J)+IEN	,	002130			
205	WRITE(/*3009) (HDP(I*J)*I=1*IEN) GX=00UT		002150			
•	NSED=NSED2		- 002160 - 002170		·····	
210	130 CONTINUE		002180			
<b>• • •</b>	CALL OUTPLY		002200			•-7
	STCP		002220			
			- 082220 -			
	C FORMAIS 500+ INPUT FORMATS		004250			
	C FORMATS 200 + OUTPUT (PRINTER) FOR C FORMATS 3000+ OUTPUT (PUNCH) FORMA	MATS TS	002270			
220	C 501 FORMAT(F10 1.F10.2.2F10.8.F10.2)		002290			
	502 FORMAT(13+2F10+2)		- 002310			
245	504 FORMAT(13)		002330			
225	506 EORMAT(213)		002350			
	2001 FORMAT(141,///+20X++FLOOD ROUTING	USING THE EXPLICIT FINITE **	#002360			

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LISTING - FINDIF/5

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230		SCONNI	TIUNS*+	//+35X++1	OTAL LENGTH	#*+F1	2.1, FTA. /- 35Y. #01.01	00238			
230		4 14X	**34*F1 1MUM FF	2.3.4 FT/	FT4,/935X,4M	ANNINGS Nº	X+4#4+F12+5+/+	35X 00240	č		
	2002	6+F17. FORMA	2 + + F T / SI	C*) 2⊐X.+PR(	BLEM VARIABL	ES*•/•		00242	Ŏ		
235		1 35×.	*NUMHER 1+* SEC	NUS	ES =49X+13+ 5X+4MAX TIME	INCREMENT	INCREMENT +5X+	NDS*) 002440	ç ç		
	2003 2004	FORMA	I (//+35 I (//+25	( THE EF	ELEVATIONS	AND INITIAL	WATER LEVELS A	T EACH 00247			
	2005	24CHAI	NAUF 9.9.	( *INVER1	ELEVATION +	13X, WATER	EVEL*1	00249	0 	······	
	2006	FORMA	T ( 354 . 4	FAT NUMP	OF HYDROGRPH	•9X•13) NO*•13•*=*	F8.2.* HRS*./.	00251	Š		
245	2008	FORMA	35X.4	TUTH OF	OF HYDROGRPH CHANNEL =	NO* 13 * #	F8 294 HRS4)	00253	8		
	2009	1/+40X	(/,J5X       0 + 0       /, 36X	NIFORM F		UNSTEADY	AT SECTION #91: "Low#) The tiow to the	002561 002561	0 0 		
250	2010	1* FI 0 2* No=	W VS TI * 14}	4F * • / • 35)	+ 000 AT SEC	TION NO=+.I	+6x, B+++ AT SE	CTION + 00256	0		
	2011	FORMA	TIME +./	+1HIS IS	A PLOT OF TH AT SECTION	HE RATIO OF NG=#114,6X1	DEPTH TO MAX DI	PTH** 00260	0		
755	2012	FORMA	* 14) T(/,J5X	HYO_HYDE	OGRAPHS WERE	PUNCHED		00263	0		
255		FORMA	1 (2x,+0)	PIH=*+F7	+2+* WIDTH=*	F7+21+ SLO	E=++F10+8++R	C=++F10002650 D02660	ő		
	3002 3003	FÖRMA FORMA	T (*STOR	HATING	CURVE ID=+,	Il+* VS=1	IN E=++F8.2)	002670	Ŏ		
260	3004	FORMA	T (*RECAL NA=1	HYD4.1	1X+#ID=#+I1+ +#PEAK=**F7*	₩ ΗΥ <u>ρ</u> Ν∩= 1,* R0=#∎	12.0. DT=+.F	<b>5</b> 00269	0		
	3005	FORMA	T(25(/+)	54.6(1)	E8.0) +1x))			27200	0		
265	3007 3007	FORMA	TI TES	NUMBER	E SEC NO=+ 1	3.* DT=*.F9	6++ HRS*+	00274	Š ·		
	3009	I + NO FORMA	PIS##,	(3) (F(.2))			•	00276	0		
	•	END						00278	0		
				<u></u>							
	····				·····						

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# LISTING OF PROGRAM KINDIF

1TAPE7=PUNCH) DIMENSION H(207), R(2(1), WL(20), TWL(20), X(200), BOT(200) DIMENSION AL(20), FUN(2), TWL2(20), FUN2(20), GINIT(200) DIMENSIONHYD(200, FUN(2), FEAK(2), TEND(2) DIMENSION HOP(2), U,2), PEHYD(2), TEOP(2) DIMENSION HOP(2), U,2), PEHYD(2), TEOP(2)	000101 000110 0000130 0000130 000150 000150 000150
DIMENSION AL(20),FUN(2),TWL2(20),FUN2(2), JINIT(200) DIMENSIONHYD(20),2),DIHYD(2),PEAK(2),IEND(2) DIMENSION HYD(2),U,2),PEHYD(2),IENP(2) DIMENSION HYD(2),U,2),PEHYD(2),IENP(2)	000120 000130 000130 000150 000150 000155
DIMENSION HOP (2) () 2) PEHYD(2), IEOP (2)	000150 000150 000155 000155
11 T ME NIS T PN - PT 1 2 9 1 . 1 1 1 2 9 1	000150 060155 000160
DIMENSION DITIIIIIIIII	630160
C THIS PROCRAM SOUTES & FLOOD USING & KINEMATIC WAVE HETHOD	0.30170
C A DICTIONARY OF THE VARIABLES FOLLOWS	666180
C BOT(I) ELEVATIONS OF THE CHANNEL POTTONS AT THE SECTIONS	666233
C CEL THE CELERITY AT BANK FULL NORMAL FLOW	000210
C DTYAR THE MAXIMUP TINE STEP AT BANK FULL CONDITIONS	000230
C H(I,J) WATER LEVELS AT THE SECTIONS	000250
C NEEDS NUMBER OF SECTIONS IN THE CHANNEL	000265
C O(I, J) FLOW PATES AT THE SECTIONS	600280
C QOUT THE VOLUME OF WATER WHICH FLOWED OUT OF THE CHANNEL	032360
C RC ROUGHNESS COEFFICIENT MANNINGS N C TIMEI STAPT TIME	000310
C TIMEF THE FINISH TIME THE THE FINISH TIME OF THE POINTS IN A STAGE DISCHARGE CHRVE	010330
C HEII THE WATER LEVELS OF THE POINTS IN STAGE DISCHARGE CURVE	0 0 7 350
C XTOTAL THE TOTAL LENGTH OF THE CHANNEL	003373
C PEAD PROPLEM VARIABLES	0 J 0 380 0 J 0 380
PFAD(5, CO1) XTOTAL, T, °C, SLOPE, DEPTH	020-00
READ(5,503) ITEST	000420
PEAD(5,506) NPU,NPU2 READ(5,506)NSED1+NSED2	000439
PEAT (5,516) NSED1, NSED2 BEAD (5,516) NSED1, NSED2	000450
PEAD(5,515) (QI(I),TI(I),I=1,NQPTS)	010475
C SET PROPLEM VARIAPLES	060493
TLT(1)=-1.0 G=32.2	000565
	0.530
CALL HANNER (A, P, SLOPE, PC, G, RMAX)	001554
C SET UP THE INFLOW HYDROGRAPH NPTS=II/FE/DI	000569
00 180 I=1, 00PTS	000580
DX=XTOTAL/FLOAT (NREACH)	010600
NSFCS=NFEACH+1 NPWL=20	000630
IE(40(1)=IEND(2)=IEDP(1)=IEDP(2)=NPTS	000660
TIMEI=0.C	030675
UEL=SIRI(G+DEPTH)+QMAX/(DEPTH+T) DTMAX=DX/CEL	000683 000693
00 193 J=1.NPTS TTYED=FLOAT(T-1)+DT	000760

#### LISTING - KINDIF/2

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CALL INTER1 (QI,TI,NQPTS,TIMEQ,QINIT(I)) 190 CONTINUE 000723 000740 g DERIVE STAGE DISCHARGE CURVE DH=DEPTH/FLOAT(NPWL-1) AL(1)=0.0 010750 011773 AL(1)=7. HL(1)=7. THL(1)=7. DO 7 I=7. DO 7 I=7. HL(1)=WL(1-1)+DH A=T\*FLOAT(I-1)+DH AL(1)=A\*DX P=T+2.+FLOAT(I-1)\*DH CALL MANNG7(A,P,SLOPE,RC,G,OX) TWM (T)=7X 0.0783 510791 000810 000359 7 THE(1)=2X 030860 č 000870 SET INITIAL CONDITIONS AT STEADY STATE CALL INTER: ()I,TI,NOPTS,TIME,Q(1)) 897(1)=XTOTAL\*SLOFF+1J(. 056980 036893 000910 X(1)=1.5 CALL INTEP1 (WL,TWL,NPWL,O(1),FLOWD) DPOPENY\*SLOPP H(1)=301(1)+FLOWD 0:0920 CALL INTEP1 (4, THL, NPHL, 0(1), FLOHD) DPC=NYSLOPE H(1)=301(1)+FLOHD DO 6 I I2, NSECS GOT(I)=POT(I)+FLOHD 6 H(I)=GOT(I)+FLOHD C GALCULAIF INTIALD STORAGE IN THE CHANNEL STOPE1=XTOTALIAT FLOHD OT4VT(1)=GTHYD(2)=DI/363C. DO 201 INETA=15 BCTA=1.(-FLOAT(IALFA-1)/4.0 OO 211 TALFA=1)/4.0 OO 211 TALFA=1)/4.0 C HVTIE THE CONTITIONS FOD THIS RUN WGITE(6,2:1) YTOTAL, DEPTH, SLOPE, RC, QMAX, CEL WGITE(6,2:1) YTOTAL, DEPTH, SLOPE, RC, QMAX, CEL WGITE(6,2:1) YTOTAL, DEPTH, SLOPE, RC, QMAX, CEL WGITE(6,2:1) YTOTAL, DEPTH, SLOPE, RC, QMAX, CEL WGITE(6,2:1) NEACH, JT, DTMAX OIN=7001=STOPE7=TTMEEFHOMI=FMOHO=0.0 PFAX(1)=PFAX(2)=PEHYO(1)=FHOHO=0.0 PFAX(1)=FAX(2)=PEHYO(1)=FHOHO=0.0 PFAX(1)=FAX(2)=FEHYO(1)=FHOHO=0.0 PFAX(1)=FAX(2)=FEHYO(1)=FHOHO=0.0 PFAX(1)=FAX(2)=DEHYO(1)=FHOHO=0.0 PFAX(1)=FAX(2)=DEHYO(2)=1.0 PFAX(1)=FAX(2)=DEHYO(2)=1.0 PFAX(1)=FAX(2)=DEHYO(2)=1.0 PFAX(1)=FAX(2)=DEHYO(2)=1.0 PFAX(1)=PAX(2)=DEHYO(2)=1.0 PFAX(1)=FAX(2)=DEHYO(2)=1.0 PFAX(1)=FAX(2)=DHYO(2)=1.0 PFAX(2)=FAX(2)=DFAX(2)=1.0 PFAX(2)=FAX(2)=DFAX(2)=DFAX(2)=1.0 PFAX(2)=FAX(2)=DFAX(2)=DFAX(2)=1.0 PFAX(2)=FAX(2)=DFAX(2)=1.0 PFAX(2)=FAX(2)=DFAX(2)=1.0 PFAX(2)=FAX(2)=DFAX(2)=1.0 PFAX(2)=FAX(2)=DFAX(2)=1.0 PFAX(2)=FAX(2)=DFAX(2)=1.0 PFAX(2)=FAX(2)=DFAX(2)=1.0 PFAX(2)=FAX(2)=DFAX(2)=DFAX 666533 000940 0u0 950 000960 050970 980 0000090 021463 001410 001016 001217 001320 001230 001040 1350 1052 031354 001356 001063 001062 001063 001064 CJ1J80 CJ1J80 331110 0 J1120 0 J1130 0 J1140 001150 0.1165 001180 001190 001200 ČÁĽĹ KÍŇRUT (FUN,TWL,NPWL,GUN,TWL2,NPWL2,QLT,TLT,1,DX,DT,NPTS,Q,H)ÖÖĪŽĪŎ

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DE 123 NET NEED 12 20 TO 22 CALL INTER1 (WL, TWL, NPWL, Q(II), HD) 23 H3P(II)=HD/3EPTH 22 IF ((I+1).NE.NSED2) GO TO 24 DO 25 II=1, NPTS CALL INTER1 (WL, THL, NPHL, H(II), HD) 25 HDP(II, 2)=HD/3EPTH HDP(11,2)=HD/JEPTH CANTINUF IF(I+HE+NSEO1) GO TO 2( NO 21 TI=',MPTS OIN=GIN+3(I) FM3M5=FLOM3 +3(I) FM3M5=FLOM3 +3(I)\*FLOAT(II-1)\*DT/360G. N32=O(I)/3MAX IF(PEAK1-GT.039) GO TO 21 TME1=FLOAT(II-1)\*DT/363(.0) PEAK1=300 24 TIME1=FLOAT(II-1)\*DT/363(.0 PEAK1=70 HY7(II,1)=773 IF((I+1).NE.NSE72) GO TO 40 DO 41 TI=1,NPTS 00'1T=700T+4(II) PM2:I=F:0~I+7(II)\*DT\*FLOAT(II-1)/3600. 071=H(TI)/7~4X IF(0EAK2.GI.377) GO TO 41 TIMT2=FLOAT(II-1)\*0T/363(.3) PEAK2=703 HY70(IL-2)=700 21 C NSED=NSED2

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831239

821348 01126 001270

011334

001340 001350

HETTE (7:3F14) ALANGERADTHYD(J), QMAX, QX, IEND(J) 821869 WPITE(7,3605) (HYD(I,J),I=1,IEN) IFINIS=-1 WRITE(7,3010) IFINIS 011830 NSED=NSED2 00185 13C CONTINUE 12C CONTINUE 210 CONTINUE 210 CONTINUE 210 CONTINUE STOP 0 . 1 9 7 691886 001982 001886 111990 CCCC FORMAT STATEMENTS FORMATS 500+ INPUT FORMATS FORMATS 2000+ OUTPUT (PRINTER) FORMATS FORMATS 3000+ OUTPUT (PUNCH) FORMATS 001340 561 FORMAT(F10.1,F1).2,2F10.8,F10.2) 502 FORMAT(I3,2F10.2) 503 FORMAT(I3) 504 FORMAT(I3) 11199 002010 515 FORMATIFIC.2,F1C+C) 516 FORMATICING 2011 FORMATICING 2011 FORMATICING 2011 FORMATICING 2011 FORMATICING 2011 FORMATICING 2011 FORMATICING 2011 FORMATICING 2011 FORMATICING 2011 FORMATICING 2011 FORMATICING 2011 FORMATICING 2011 FORMATICING 2012 FORMATICING 2012 FORMATICING 2012 FORMATICING 2012 FORMATICING 2012 FORMATICING 2012 FORMATICING 2013 FORMATICING 2014 FORMATICING 2014 FORMATICING 2015 FORMATICING 2016 FORMATICING 2016 FORMATICING 2016 FORMATICING 2017 FORMATICING 2017 FORMATICING 2016 FORMATICING 2017 FORMATICING 2 505 FORMAT(FIL.2,F13.5) 506 FORMAT(213) 002310 0)2260 0)2270 0)2280 1 VS (1F2,7,33,7,7)(C M) SCOTAGE NO ,17,00, 2+ 10=+14) 2012 FORMAT(/,35X,\*NO HYDPOGRAPHS WEPE PUNCHED\*) 2013 FORMAT(/,35X,\*THE STAGE DISCHAPGE CURVE WAS NOT PUNCHED\*) 2014 FORMAT(35X,\*ALFA\*,15X,\*=\*,F12.3,/,35X,\*BETA\*,15X,\*=\*, 612293 012295 2115 F03MAT(35X,\*HYDROGRAPH NO\*,I3,\* PEAK=\*,F6.4,\* TIME=\*,F6.3,\* HRS\*) 022296 30C1 F09MAT(2X,\*DEPTH=\*,F7.2,\* WIDIH=\*,F7.2,\* SLOPE=\*,F10.8,\* RC=\*,F10.323C0 1.8) 3672 FORMAT(\*STORE RATING CURVE ID=\*,11,\* VS=1 MIN E=\*,F8.2) 3603 FORMAT (2000,3(F14.3,10)) 3604 FORMAT(\*RECALL 490\*,110,\*ID=\*,11,\* HYO NO=\*,13,\* DT=\*,F9.6, 002320 002340

### LISTING - KINDIF/5

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 1\* 19RS
 DA=1.0\*,/,21X,\*PEAK=\*,F7.1,\*
 R0=\*,F12.0,\*
 NO
 PTS=\*
 0)2350

 30[5
 F0.MAI
 (25(/25X,6(1X,F8.(1,1X)))
 0)2370
 0)2370

 30[5
 F0.MAI
 (25(/25X,6(1X,F8.(1,1X)))
 0)2320

 30[6
 F0.MAI
 (25(/25X,6(1X,F8.(1,1X)))
 0)2320

 30[1
 F0.MAI
 (10,12)
 0)2320

 30[1
 F0.MAI
 (13)
 0)2426

 30[1
 F0.MAI
 (F12.3,F12.3)
 0)2426

 31
 F0.MAI
 (110(1X,F7.2))
 0)22420

 END
 0)22420
 0)22420

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## LISTING OF SUBROUTINE KINFUN

5	SUR DIN 1 GUS 2 FOS 8 FF	POUTINE KINF SNSTON SUS († 1 1=1,NUS (1)=BFIA*0US 7 T≈1,NUS (1)=BEIA*QUS (1)=BEIA*QUS (1)=	UN (ALF, RFTA US), QUS (NUS) (1)-ALF*SUS) (1)+(1.0-ALF	.DT.SIIS.JUS.X ),SDS(NDS),QN (T)/JT -)+SNS(T)/NT	VUS, SDS, NNS, NNS NUS, SDS, NNS, NNS NNS), GIS (NUS	, GUS, FDS) , FDS(NOS)	004180 004180 004190 004200 004210 004220 004250		P a 9r	1	<b></b>
•	<b>ร</b> พับ		-				004260	· · · · · · · · · · · · · · · · · · ·			
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# LISTING OF SUBROUTINE KINRUT

SUBROUTI	NE KINRUT	73/73	OPT=0 T	RIGE	CTN	4.°+P355	08/15/73	14.55.38.	PAGE	1	-
	SUBR	OUTINE KI	18117 (595,	105, NOS, GUS, QU	S, NUS. OLT. TLT	.NLT.DX.DT.N	TN, 003980				
	10IN - DIME 10001	NSION FDS	נאוחבָזי, פּטפ נ	NUS1 . 205 (NDS) .	LT (NLT) , QUS (	NUSI, QIN (NIN	003909 • 004000				
5	C ÓÓÚI	(1) 405T -	TE DEFINED	5,000T(1),FAC3	ING THE SUBRO	UTTNE.		<u> </u>			
	CALC DC 1	INTER1 () I=2.NIN	505,205,NU	S.OTN(1),GAC1)			004949 204052	-			
19		 F=0_0 TVT(1)_LT	074105400	0 2			004057 004979 004089				
		THTER1 T	J +DT LT,TLT,NC	T.TIME.OSIDEJ							
15	2 FAC4	E=2510E*0 =FAC3+GAC	( [+6107-00U	I(I-1)+01N(I-1)	+DSIDE		004117 004120				
	- FAC	=FAC4 =GAC2		3+FAG4+1104 111			004150 004150				
20	RETU END	¢Ň					104169 004170				
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## LISTING OF SUBROUTINE HPLOT

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SUBROUTINE MPLOT	73/73 UPT=0 TRACE	FTN 4.0+P355	01/28/74	11 <u>.</u> 35,18.	PAGE	1
c	SUBROUTINE HPI UT (OCFS+DT+IEND+P THIS SUBROUTINE PLOTS EITHER 1 DIMFNSIUN CCFS(200+2)+DT(2)+TEN	EAK.TIME) DR 2 HYDROGRAPHS ON A SET OF AXIS D(2).PEAK(2).CFS(120).DATA(2)	002790 002800 002810			
5 100	WRITE(0,100) FOFMAT(1H1) ID1=1 ID2=2 GATA //HO,01US,84 ANK,DASH,DOT/1	46.184.18 .184.18.2	002820			
10	WRTO         # 1           XMRTO         # 1.           MAX         = 11A		002870 002890 002890			
C 15 C 27 28	ARE THERE 1 OF 2 HYDROGRAPHS IF(102) 7. 27. 28 DETERMINE HIGHEST PEAK IF 2 HYD OMAX = PFAK (ID1) GO TO 3/ IF(PEAK (ID1) - PEAK(ID2)) 29. 2	ROGRAPHS 80 30	0022920 0022920 0022930 0022950 0022950			
<b>20</b> 29 30 C	GMAX = PFAX(INC) GMAX = PFAX(INI) IF > Hyurngraphs Determine Larg Hydrograph IF Necessary	EST DT AND INTERPOLATE OTHER	- 002990 002990 003000 003010			
25 31	IF(DT(IU)) = DT(ID2)) 32, 33, 3 L = IU2 GO TO 35 I = IU2	•	003020 003030 003040 003050 003050			、 
30 35	R = ΙΟΪ M = ΙΕΝΟ(L) ΤΙΟ = υΓ(K) ΤΙΟμ = υ <sub>Γ</sub> (K) DO 75 I = 2. M		003070 003080 003090 003100 003110			
35 77	ΤΊΟΗ = ΙΥΛΗ + UT(L) ΙΓ(ΤΙV - ΤΤΟΗ) 76•77•75 J = J • Ι		003120 003130 003140		•	
40 76 1	ČFS(J) = OCFS(1+L) TIO = TIO + DT(K) GU TO 75 J = J + 1 CFS(J) = OCFS(1-1+L) +((TID - T - CCFS(1-1+L)) TID = TIO + DT(K)	[DH + DT(L)) / DT(L)) + (OCFS(I,L)	003150 003160 003170 003180 003190 003200 003210			
45 75	CONTINUE IEND(L) = J DI(L) = DI(K)	· · · · · · · · · · · · · · · · · · ·	003220 003230 003240			
78 33 50 37	UO 78 1 = 2+ J OCFS(I+L) = CFS(I) IF(IFNU(ID1) = IEND(ID2)) 37+ 3 M = IEND(ID1) GO TO 39 M = IEND(ID2)	7, 38	003250 003260 003270 003280 003290 -003300 -			
55 C 39	TF(M - MAX) <sup>24</sup> 59 459 64 DETFHMING TIRF SCALE MRTO = MAX / SO TO 4		003320			

### LISTING - HPLOT/2

SUBRO	UTINE HPLOT 73/73 UPT=0 TRACE	FTN 4.0+P355	01/28/74 11.35.18.	PAGE 2
	64 M = MAX		003360	
60	C PLOT HYUROGRAPHS		003380	
	65 CFS(1) = 0ASH *	AX) . DOT .	003400 003400	
65	57 FORMATLIX+F5+P+* *,119A1)		003420	
•-	JI # 10 DO 50 J # 19 50		003440 003450	
	• IF (J = J) -7. 68. 67		003460 - 003470	
70	$\begin{array}{rcl} 69 & CF5(1) = 0 \\ 60 & T0 & 71 \end{array}$		003480	
	67 DŎ 7Ŭ 1 # 1, MAX 70 CES:11 # HIANK		003500 003510	
75	71 G2 = Q1 ~ Y5Ci		003520 003530	
	D0 51 I = 2, M K = K • MHTO	······	- 003550	,,, <u>,,</u> ,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,,
	$\begin{array}{rrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrrr$	1	003570	
80	51 CONTINUE	X) - DOT	003590	
	IF(102) 52, 57, 53		003610	
85			- 003630	······
	IF (nCFS(T+102) - Q1 1 66, 56, 5 66 IF (nCFS(T+102) - Q2) 5, 5, 5, 56		003650 003660	
90	56 CFS(K) = P(U) 60 T0, 54	• •	003670	`
	5 CIS(K) = HIANK 54 CONTINUE		003700	
	3 FORMAT(1H++11X+11BAL)		003720	
43	$59 J1 \pm J1 + 10$		003740 +	
	2 FORMAT(14++F6,4)		003760	v
100	60 FORMAT(11X.120A1) 50 CONTINUE		003780	
	CFS(1) = TTME DTT = UT(101) * 10. / XMRTO		003800	
105	C PUT TIME ARRAY IN CFS AND WRITE	TIME SCALE	003830	
	$\begin{array}{l} 61  CFS(1) = CFS(1-1) + DTT \\ & WRITE(6+62)  (CFS(1) + 1 = 1 + 12) \\ \hline \end{array}$		003850	
110	02 FURMATIOX+12F1V+3) WRITE(6+63)		003870	······
110	6 RETURN		003890	

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# LISTING OF SUBROUTINE RESVOR

	COMMON VELOWAX	ES, TWL+NPWL+DT+NPT5+Q+H)	00	4270 4280		
5		)) • IWL (201 • WL 1 (20) • G (NPJ • PP (2) • HH (200 • 2) • I 1 (2) • D	S)+H(NPIS) 00 TH(2) 00 00 00 00	4290 4300 4310 4320 4320		
10	CO 1 I=1.20 GUS(I)=TwL(I)/2.0 wL1(I)=TwL(I) FUS(I)=wL(I)/2.0+RKgWL1 CALL KINDUT(FUS)WL120.6	(]) ##XRES/D] 3US9TWL9209QLT9TLT91940X4	00 00 00 00 00 00 00	4340 4350 4360 4370 4370		
15	II(1)=II(2)=NPTSPP(1)=PP(2)=1.0DTH(1)=UTH(2)=UT/3600.0DO 2 I=1.NPTSHH(I.1)=UTI/24AX			4390 4400 4410 4420 4420		
20	TIME UN CALL HPLOT(HH.UTH.II.PP. WRITE(5,100) 100 FORMAI(350.0THIS IS THE	HYDROGRAPHS AT THE IMAG	INARY RESERVOIR++ 00	4450 4460 4470 4480		<u></u>
25	UO 3 1=1.eNPTS 3 Q (I) == H (I) RETURN END	JUIFLU##*** */	00 00 00 00	4500 4510 4520 4530		
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## LISTING OF SUBROUTINE RIVER3 (AND THE DRIVER)

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DIMENSION WED(20), DD(20), WK1(20), WK2(20), WK6(20), WK7(20), WK9(20)	
CALL CONNEC(6LOUTPUT) WHILE(6.1)	0 0 0 1 3 0 0 0 0 1 4 0
1 FURMAT(A SUPPLY NO CF SECTIONS, , 13*)	
2 FURMÁT(15) #R17E(6,3)	
3 FURNAT (* SUPPLY MAX NO OF PTS,,,I3*) READ(5,4)MAXPIS	000190
• 4 FURMAT(13) NF1=10	000210
CALL HIVERS(8, H, NPTS, X, PC, NSECS, MAXPTS, 32, 2, 1, 5, 6, NF1, AT, +ELLY, GT, NLD, GD, NN1; NK2; HK6; KK7; NK9)	000220
	000230 000240
UIMENSION B(NSECS, MAXPIS), H(NSECS, MAXPIS), X(NSECS),	0002501
READ(NEEAD, 10)I, x(I), RC(I), NPTS(I)	000270
NEAD (NEAD, 207 (B(1, J) M(1, J), J=1, N)	000290
20 FURMAT(3x,2F0,2,1x,2F0,2,1x,2F0,2,1x,2F0,2,	000310
HETUHN END	
SURHDUTINE RIVERS(8, H, NPTS, X, RC, NSECS, MAXPTS, G, NR1, NR2, NH, NF1, AT, SELEV, WT, WED, QD, WN1, NN2, WKD, WN7, NK9)	T, DODAG START OF RIVERS
01MEN310N B(N3ECS; MAXP13); M(N3ECS; MAXPT3); NPT3(N3ECS) + , x() SEUS), HC(NSECS)	
0144NS10H A1(NSECS,NF1),ELEV(NSECS,NF1),OT(NSECS,NF1) 0144NS10H ITST(2,100),B1(30),H1(30),OS(20,2),1040(20)	000390 000400
01MENSION ALD(NF1),UD(NF1),TI(20),OI(20) UI4ENSION AK1(AF1),AK2(AF1),AK0(AF1),AK7(AF1),AK9(AF1)	000410
014ENS10/14K3(200), #K4(200), #K5(200), #K8(2) DI 4ENS104 NSA(20), ALF(20), BET(20)	000430 000440
4 00 5 I=1,NStCS	
10 CALL HEADXS(H, H, NPTS, X, RC, NSECS, MAXPTS, NR1)	000490
11 FURNAT(* DU YOU WANT LIST OF CUMMANDS1, YES/NO *)	000520
12 FURMAT (A3) IF (SKIP-EW.3HNO ) GUTU 21	000540 000550
13 WRITE(NA,20) 20 FURMAT(* AFTER INVITATION TO TYPE "X=" GIVE ONE OF*,/,	000560
* THE FOLLOWING COMMANDS * // * DISCHARGE TO SPECIFY FLOW * //	000580 000590
* # D/S WLTO DEFINE DOWNSTREAM CONTROL LEVEL * //	00000
+ A OLD SECTION, TO PHINT COOKUS OF A SECTION * //	000950 000950
+ OLD CUEFFTO PRINT ROUGHNESS MEASURE + //	000640
• * CRITIC TO CUMPUTE CRITICAL DEPTH AT A SECTION	000650 000670
TABLE TABLE TO PRINT TABLE OF ALL SECTIONS DATA#7/7	

LISTING - RIVER3/2



### LISTING - RIVER3/3

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WRITE(NW,416) NF1 416 FOMMAT(* NO, OF PTS, TRUNCATED TO *,13)	001390 001400
GIG 151 GOTO 25 220 WHITE[NW,221] 221 FORMAT(* DEFINE SEC, NO, WHERE D/S WATER*,/, * * LEVEL TO BE SPECIFIED, I3 *) READ(NR),2221NDS	001410 001420 001430 001440 001450
222 FURMAT(13) 1F(VDS,LE,NSECS,AND,NDS,GT.0) GOTO 223 HILFLAND, DATE AND AND AND AND AND AND AND AND AND AND	001480
223 NPSP=NDS TFA=0	001500
HITE(NW, 412) 412 FURMAT(* DEFINE NUMBER OF POINTS IN D/S WL RATING CURVE (13)*) READ(NR2,212) NPWL IF(NPWL, LT, NF1) GO TO 441	001550 001550 001550 001550
<pre>white(va,442) NF1 442 FUMMAI(* NU, UF PTS TRUNCATED TO *,13) 441 AKITE(NW,225) 225 FUMMAI(* DEFINE D/S WL AND FLOWRATE 2F9.3**/* * * START AT THE LOWEST WATER LEVEL*) DU 413 I=1,NPWL WHITE(NW,407) I</pre>	001570 001580 001580 001600 001610 001620 001630 001640
413 HEAD(192,226) HED(1); GO(1) 226 FURMAT(2F9,3) IF2=1	001660 001660 001670
GOTU 25 230 1+(1+3,EG,1) GOTO 236 MR 1+E(Nw,231)	
231 FORMAT(+ SPECIFY RESISTANCE LAW BY TYPING+,/, + CHEZY, FANNING, STRICKLER, COLEBROOK+,/, + SCHEZY, FANNING, STRICKLER, COLEBROOK+,/,	
<ul> <li>A ROUGHNESS MEASURE (OMPATIBLE *, /, * *=*)</li> <li>GOTO 232</li> <li>THITE(N, 237)</li> <li>FUMMAT(* RESPECTERY RESISTANCE LAW*, /, * **)</li> <li>* SEE PREVIOUS NOTE *, /, * *=*)</li> <li>READ(N, 2, 233) UNAME</li> </ul>	001740       001750       001770       001770       001770       001770
<pre>233 UNALLADJ IF (UNAME, EU, OHCHEZY ) 1G=1 IF (UNAME, EU, OHMANNIN) 1U=2 IF (UNAME, EU, OHSTHICK) 1U=3 IF (UNAME, EU, OHSTHICK) 1U=4 IF (UNAME, EU, OHSTUUEH) 1U=5 IF (UNAME, EU, OHSTUUEH) 1U=5 IF (IU, UT, 0) OUTU 235</pre>	001800 001850 001850 001850 001850 001850 001860 001870 001880
234 FURMAT(* NUT HECOGNIZED.,,PLEASE RETYPE*,/,* X=*) GOTU 232	
235 1+3=1 1+6=0 6010 25	001920 001930 001940
240 #HILE[NH,241] 241 FURMAT(* UEFINE SECTION NO., IS *) REAU(NR2,244)NN	001960 001960
242 FURMAT(IS) IF (VS.LE.NSECS, AND, NS. CT. 0) GOTO 245	001490 001490
246 FURMAT (* SEC, NO. *, 15, * OUT OF RANGE 1 ++, 13) GUTU 240 OUT 040 OUT 040 OUT 040 OUT 040 OUT 05 RANGE 1 ++, 13)	002010 002020
(243 FORMAT(* SECTION NO. *, 13, * HAS*, 13, * PTS, COORDS ARE*)	0020340 
HHTIC(NH+C44)(B(N3+1)+I#1+NM)	

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# LISTING - RIVER3/4

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244	FORMAT(5(F5.1,*.*,F5.1,2X)) Gutu 25	00207	0.7.9	
(250-	WHITE (NH, 251)	-00209		
	REAU(NR2,252)NS,NP IF(NS,LE,NSECS,AND,NS,GT.0) COTO 257	00211	110	
	GUTU 250 NS, NSECS	00213		
257	IF(NP,LÉ, MAXPTS, AND, NP, GT, O) GOTO 259 MMITE(NA, 258) NP, MAXPTS	21200	150 160	
258-	FURMATIA NO.DF PTS. #, 15, # OUT OF RANGE 1+#, 13) GUTO 250		170	
252	FURMAT(213) ITST(2,NS)=1	00219		
	NPTS(NS)=NP 166=0	15500		
253	WRITE(NW,253) Fuhmai(a give coords in Pairs, 2F6.1#)	55500 15500	230 240	
	- OU 250-1=1,NP	22500	250	
255	FUHMAT(* PT,*,I3,** *) REAU(N42,256)B(NS,I),M(NS,I)	00221	270 280	
256	FUHMAT(2F6,1) Cuviinue	00229	290	
260	GUTU 25 MKIIE(NM, 261)	00231	310 520	
261	REAU(NR2,262) NS	00233 00234	330 340	
262	IS (NS LE , NSECS, AND NS GT. 0) GOTO 264	00235	350 360	
24.0	GUIU 260	00237	370	
263	FURMAT(* ROUGHNESS AT SECN, *, 13, * IS*, F6, 3)	00239	490 490	
370	WHILE (NW, 271) Frummatic (W, 271) Frummatic (GEFINE SECTION NO. + COFFE IX + FE X AN	00241		
272	READ(A42,272) NS,RC(NS)	00244		
1	IF (NS, LE, NSECS, AND, NS, GT, 0) GOTO 273	00245	450 470	
275-	GOIO 270 ITST(1,NS)=1	00248	480	
	IF 6 = 0 GU TO 25	00250	500	
280	CONTINUE	00252	520	
	IF(IF2.E0.0) GOTO 220	00254	540 550	
	1F(1F3,E0,0) GOTO 230	-00259	570	
2/8	+ HY SEC, NU, (I3)*)	00258	580 590	
279	FURMAT(15) Torang te and and and ct as core or	00260	510	
	WHITE (NW, 246) NUS, NDS	00263	530	
-281-	FURMAT (///, - SEC. DISTANCE DISCHARGE WATER ENERGY		550 560	
293	ANITE (NA, 417) FURMATIA SPECIFY D/S STARTING POINT OF PROFILES -/-	00267	570 580	
}	• • 6Y SEC, NO, (I3)•) READ(NR2,212)NSDS	00259	590 700	7.1
	IF(NSDS,LE,NDS,AND,NSDS,GT.0) GO TO 418 White(NW,246) NSDS,NDS	0027	710 720	.42
<b></b>	-60 <sup>-</sup> 10 <sup>-</sup> 293 <sup>-</sup>	-00273	/30	!Y

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418 IF (NSUS, GE, NPSP) GO TO 419 WRITE (NW, 420) NSUS, NPSP, NOS 420 FURMATIA STATE NAME DELATION NOT DEFINEDA. 4	- 002740
+ * SECTION NU, *,13,* RELATION DEFINED AT*,/,	002770
+ * SECTIONS *,13,* * *,13)	002780
GU TU 293	002780
419 NPSP=NUS	002810
1F(NSOS,NE,NDS) GO TU 421	002810
UU 422 [=1,NP#L	002820
422 MK2(1)=M(1)	002830
422 MK2(1)SUD(1)	002840
425 MK2(1)SUD(1)	002850
GU 10 433 421 DU 424 1=1, NRPTS MS1(1)=FLEV(NSDS, 1)	002860 002870
424 mR2(1)=UT(NSDS,1)	002900
NRK1#VRPTS	002910
433 VRK2=NRPTS_1	002910
IF 7=1 F (VmK2;EU;0) - NmK2=1 IF 0=1 IF 0=1	002430 002430
00'300 IJ=1,NRPTS DJ 274 J=1,NQS JJ=NDS+1=J	
IF(NDS,GE,IOFQ(JJ)) GO TO 275	C Č Ž 4 8 Č
274 CONTINUE	C C 2 4 9 Q
275 V=(US(JJ,2)+VS(JJ,1))*FLOAT(IJ=1)/NWK2+QS(JJ,1)	C C 3 4 9 Q
CALL INTERI(HRI; AKI; AKI; U; ALDS)	0 0 3 0 1 0
USAL=ALDS	0 3 0 2 0
CALL SELSEC(B.H.NSECS, MAXPTS.NSD3.BI.HI)	0 0 3 0 2 0
CALL PROPS (B1,H1,NPTS(NSDS),DSWL,A,T,P,AY) At(NSUS,IJ)=A ELEV(NSDS,IJ)=NLDS	
US124US1120 Ex====================================	003070 003075 003080
DU 300 12=\US1,NSD3	003090
1=\\S\S+\US-12	003100
QU 276 J=1,NUS	003110
JJ=VdS41-J	003120
IF(I GE,10FJ(JJ)) GOTO 277	003130
276 CUNIANUE	003149
27/ U= (US(J); 2)=US(JJ,1))+LUAT(IJ=1)/NMK2+US(JJ,1) 	003150 003160 003170
284 CALL EFRA(B, H, NPTS, X, HC, NSECS, MAXPTS, I, WLDS,	003180
+ 0, G, GHELYU, ALUS, EN)	003190
GOTU 291	003209
CALL EZMA(B, M, NPTS, X, MC, NSECS, MAXPTS, I, WLDS,	003210
+ U, G, MANNGU, WLUS, EN)	113220
GOTU 291	13220
286 CALL EZRA(8,H,NPTS,X,HC,NSECS,MAXPTS,I,HLDS,	003250
+ OJ(G,SIHICU,HLUS,EN)	003250
287 CALL EZRA(U,H,NPTS,X,HC,NSECS,MAXPTS,I,WLDS, + 0,G,COLEG, WLUS,EN)	003280 003280 003290
280 CALL EZRA(B,H,NPTS,X,HC,NSECS,MAXPTS,1,HLDS,	003310
289 CALL EZHA (B, M, NPTS, X, RC, NSECS, MAXPTS, I, WLDS,	003340
, G, G, RDUGHG, WLUS, EN)	003350
326 CALL SELSEC(B,H,NSEC3,MAXPTS,I,BI,HI) GOTU 199 HITELON 137	003370 003370 003380 003380
327- FORMAT(1X,* SPECIFY HEADLOSS-COEFF	
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# LISTING - RIVER3/6

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399 CL#1.0	003430
CALL CONTROLBI, MI, NPTS(1), MLDS, EN, G, G, CL, MLDS, EUS)	03450
	003460
CALL SELSEC(B, H, NSECS, MAXPTS, I, BI, HI) CALL PHOPS(BI, HI, NPTS(I), HLUS, A, T, P, AY)	003480 003490
	003500
300 wLOS= +LUS	003520 003530
GU TO 25	003540 003550
310 CUNTINUE 317 WHITE(NW, 311)	003560
311 FURMAT(* SPLCIFY SECTION NO,I3*) READ(NH2,312) ICRIT	003580 003590
IF (ICHIT, LE, NSECS, AND, ICHIT, GT. 0) GOTO 313	003600
GUTU 317	003620 003630
313 WRITE(NW,310) 316 FURMAT(* ENTER DISCHARGE F9,3*)	003640
READ (NR2, 318) Q 318 FUHMA1 (F9, 3)	0 / Ša 6 U 0 / Sa 7 U
CALL SELSEC(8,H,NSFCS,MAXPTS,ICHIT,BI,HI) CALL CHITIC(BI,HI,NPTS(ICHIT),Q,G,YCR,HCR)	003680 003690
CALL BUTTOM (HI,NPTS(ICRIT),ROT,W(MAX) MRITE(NW,314) ICRIT,Q,YCH+BOT,HCR+BUT	003700 003710
314 FURMAT(* AT BEC*, I3,* WITH G#*, F10, 3,/, + A CRITICAL WATER LEVEL#*, F8, 3, * AND*, /,	003720
<pre>* * CRITICAL ENERGY LEVEL#*,F8,3)</pre>	001740 a 003750
320 IF4=0	003760 003770
ff (10, E0, 12) white (NW, 329)	003780 003790
00 321 1=1, NSECS	003800
wHITE(NW, 323)1,X(I), AC(1), NPTS(I)	003820 *
322 IF (IC, EG, 11, AND, ITST (2, I), EQ, C) GO TO 321	003840
* WRITE(NH, 324)I, (B(I,J), H(I,J), J#1, NP)	003860
323 FURMAT(13, F10, 1, F6, 4, 13)	00380
325 FURMAT(* NU CHANGES*)	003910
330 FURMAT(* CURRENT DATA 18*)	
321 CONTINUE	003950
	003970
400 WHITE (NW, 401)	
	004010
NPHY()=20	
467 FURMAT (* NO, OF PTS, TRUNCATED TO 20+)	004050
402 FORMAT ( ARE THE UNITS HOURS OR SECONDS +)	
ITURO CELEGI COLLO	004090 · · · · · · · · · · · · · · · · · ·

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IF(UNITS,EQ,6HMOURS) ITU=1 IF(UNITS,EQ,6HSECOND) ITU=2	004100
( V C C C C C C C C C C C C C C C C C C	004120 004130 004140
403 WRITE(NW,405) 405 Furmat(* After Each "_" Enter time and FlowRate(2F9.3)*)	004150 004160
407 FOHMAT(A PT, +, 13, + _ +)	004180 004180 004190
$\begin{array}{c} 406 & \text{FURMAT}(2F9,30) \\ 416 & \text{FURMAT}(2F9,3) \\ 411 & \text{mHIE}(Nn,409) \end{array}$	004200 004210 004220
409 FORMAT(* SPECIFY THE SEC NO. AT WHICH INFLOW IS DEFINED(I3)*) Read (NR2,212) NSINF If(state.dt.nsfcs DR.Nsinf.lt 1) go to 410	
	004260 004260
60 TO 25 410 MRILE (NM, 246) NSINF, NSECS	004280 004280 004290
425 1F(1F7,E0,1) GO TO 426	004300 004310 004320
427 FURMAT(* PROFILES ARE NOT COMPUTED VET*) 60 TU 25	004330 004340
426 FORMAT(* WARNING, A CHANGE MAS BEEN MADE AND *,/, + * NOTHING HAS BEEN RECOMPUTED *)	004350 004350 004370
432 WRITE(NW, 429) 429 FURMAT(& SPECIFY U/S SEC, ND, DF PRINTOUT (13)*) BEAN(W22212) NUSER	
IFINUSPR, GE, NUS, AND, NUSPR, LE, NDS) GO TO 430 HITE (NW, 431) NUS, NDS	
431 FORMAT(A SEC, ND, MUST BE IN THE RANGE A, 13, 4 • 4, 13) 60 TU 432 430 mm (TE (Nm, 482)	004430
462 FURHAT(* SPECIFY D/S SEC. ND, OF PRINTOUT (13)*) READ (NH2,212) NUSFR	004460 094470
WRITE( N., 431) NUS, NDS 60 TO 430	004480 004490 ° 004500
436 1F(NDSPR;G5,NUSPR;G0-T0-434 wkite(Nw,435)NUSPR;NUSPR 435 FORMAT(+ D/S SFC, +,IS, * 18 SPFC1FTFD H/S OF H/S SFC, +,IS)	004510 004520 004530
GU TU 430 434 WRITE (Nr, 281)	004550
EN=ELEV(I,1)+UT(I,1)**2.0/(2.0*G*AT(I,1)**2.0) mile(N*,439) I,x(I),UT(I,1)*ELEV(I,1),EN	00450 004570 004580
00 438 11=2, % PTS Exetty(I,II)+01(I,II)**2,0/(2,0*G*AT(I,II)**2,0) 438 MRUTE(Na,440) DT(I,II)+F(F(I)I)=F(I,II)**2,0)	004590 004600 00460
437 CONTINUE 439 FURMAT(/,I4,E10,1,3F10,3)	004630
440 FUMMAL(147,5F10,5) ARTE(Nh,292) GO TU 25	004640 004650 004650
-475 IF(IF7.NE,0) GO-TO-483	-004680
483 IF(IF5,NE,0) GO TO 480 MKLIE(NW,484)	004690 . 004700 004710
484 FURMAT(* THE INFLOW HYDROGRAPH IS NOT DEFINED*) GU TO 25 480 TO 15 THE AN CO TO 400	
MAITE(NW, 428)	004750

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443 FURMAT(* DO YOU WISH TO CONTINUE WITH THIS COMMAND (YES/NO)*)	604770
READ(N2,12) CONTU	004780
GO TO 25	004800
490 IF (NSINF. GE, NPSP, OR, NFRP. GE, NPSP) GO TO 492	004810
491 FUR TAT (* PROFILES ARE ONLY COMPUTED UP TO SEC. *13)	004820
GO TO 25	004840
492 IF (NSINF, LE, NUS, AND, NFRP, LE, NDS) GO TO 481	004850
493 FURMAT(* PROFILES NOT COMPUTED DIS OF SEC. #113)	004860
GO TO 25	004880
481 ALFA=0,5	004890
444 ARTIE(NW,446)	004900
446 FOHMAI(* IME WEIGHT FACTORS ARE ASSUMED TO BE 0.5*//	004910
+ * DO YOU WISH TO CHANGE ANY_ (YES/NO)*)	004920
NCHNGS=0	004930
READ(482,12) CONTU	004940
IF(CONTULED.3HND-) GO TO 447	014950
RITE(NW, 522) 522 FURMAT(* HUM MANY SECTIONS ARE TO BE CHANGED, (13)*) READ("R2,212) NCHNGS LE(NCHNGS.IE.20) EQ TO 523	004970 004980 004990
494 FORMAT(* NUMBER OF CHANGES TRUNCATED TO 20*)	005010 005020 005020
523 WRITE(NW, 495)	005030
495 FURMAT(* SPECIFY U/S SEC.NO. (13),*//	005040
* ALFA AND BETA (2F9.3) FOR EA. SECTION*)	005060
DU 496 I=1.NCMNGS	005060
ARITE(NA, 497) I	005080
497 FORMAT(* CHANGE NO, *, I3)	005090
496 FEAD (NR2, 498)NSA(I), ALF(I), BET(I)	005100
498 FUHMAT(I3, 2F9, 3)	005110
447 WRITE (NW, 449)	005130
449 FORMAT (* DEFINE U/S SEC, NO, WHERE ROUTING STARTS (I3)*)	005140
READ(N+2,212) NSR	005150
IF ("SR.GE.NPSP.AND.NSR.LT.NDS) GO TO 450	005160
#FITE(NH,431) "PSP.NDS	005170
GO TO 407	005180
450 IF(NST, EU, NERP OR NSR EU, NSINF) GO TO 451	005100
ARTIE(Na, 452)NSINF, NERP	005200
452 FURMAT(* THE ROUTING MUST START AT SEC. NU.2*,//	005210
+ 1% IS,* UK *,IS)	005220
GO TU 447	00530
451 ARITE(NW,453)	005260
453 FORMAT(* DEFINE D/S SEC. NO. WHERE ROUTING FINISHES (13)*)	005260
Read(NR2,212) NFR	005262
1F(NFR,GE,NDR) 60 TO 530	005262
ARITE(NW,455) NFR,NSR	005264
GO TO 451 530 IF (NFR.GE.NSINF.AND.NFR.LE.NDS) GO TO 454	005265
WRITE(VA,455) NFR,NSINF,NDS	005280
455 FURMAT(* SEC. ND. *,15,* OUT OF RANGE *,13,* * *,13)	005290
454 ARTIE (N. 456) 456 FURMAT (* DEFINE TIME STEP FOR COMPUTATION, USE *./, *.*.SAME UNITS DE TIME AS USED FOR HYDROGRAPH (F9,4)*)	005300 005310 005320 005330
457 FURMAT (F9, 4)	005350
458 FORMAT(* ENTER START TIME WITH CONSISTENT UNITS (F9.3)*)	005370
REAU(NR2,459) TIMES	005380
459 FORMAT(F9.3)	005380
533 WRITE (N, 462)	005410
402 FURMAT (* ENTER FINISH TIME WITH CONSISTENT UNITS, (F9,3)*)	005410
PLAD (ND), HON TIME	005410
IF (TIMEF, GE, TIMES) GO TO 531	005422
532 FURMATIAN FINISH TIME *, F9.3,* IS EARLIER THAN*,/,	005426 W 005427

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# LISTING - RIVER3/9

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GO TO 533	0054 <b>28</b>
531 TT#TIMES WKA(U)=WK9(1)=m1.0	005430
IF (NSH NE NS INF) GO IO 466	005450
	005451
528 FORMAT (+ THE INFLOW HAS BEEN ROUTED THROUGH+,/,	005453
• # A RESERVOIR DO YOU WISH TO ROUTE THE #,//	005455
REAU (NH2,12) CONTU	005457
I IE (LONTULED. 3HNO ) GO TO 466	005458
527 NN=0	005459
00,461, 1=1,200	005470
	005480
1F(17,GT,TIMEF) GO TO 463	005500
	005510
461 CUNTINUE	00520
GU TU 463	005540
00 465 I=1,200	005550
ÇALL INIERI (WR3, WK3, NN, TT, NK4(1))	0°557°
TTELLADELT (IL OF TIMEF) GO IO 485	005580
NNN=NNN+1	005500
465 CUNTINUE	
TT=TIMES	005620
DU 470 1=1.NN	005040
	005650
470 WK5(I) #WK4(I)	005670 -
( 463 D)=DELTA TECTUED 2) GO TO 467	005680
DIELELA + 3600	005700
	005710
NAK22NRPTS	005730
IF(4SR-1,EQ,NFR) GO TO 529	005732
DU 468 I=NSR.NFR	005740
	005760
IF(CHIVES.EU.O) GO TO 501	005765
00 499 11≛1,ŇCHNGS	005780
499 CUNTINE (VSA(11), 20, 147) CU TU SUU	005790
	<b>ុ</b> ំប៉ន្ត៍ខ័រ្ម័
AFTA=ALF(11)	005820
501 0x=x(1)=x(1+1)	005830
$1 \qquad 1F(0x, Eq, 0, 1), GO, TO 468$	005850
	005670
Whili = BETA+OT (I + II) + ALFA+STOR	0.05480
	005400
469 WK0(1])=GT(],[])	005910
WK4(1)=WK5(1) CALL KINRIT(WK2.WK6.NRPTS.WK1.WK7.NWK2.WK8.WK9.1.DV.DT	005920
( +,NN,WK3,WK4)	005940
	005950
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468 CONTINUE 529 WKITE(NW,472) NFR	005970
472 FORMAT (//, * HYDROGRAPH AT SEC. NO. * 13,/)	005990
UU 4/3 1=1,NN TT=FLUAT(I=1)*DELTA+TIMES	006010
474 FOHMAT (1X, F11, 3, 1X, F11, 3)	006030
502 JF(JF5, NE, 0) GO TO 503	006060
GU TU 25 503 1F(1F1,NE,0) GO TO 520	006090
S21 FUHMAT (* SPECIFY THE DISCHARGES *)	006100
520 MKITE (NW, 504) NSINF, NFRP	000120
+ A AT SEC, NO, W, I3; + OR R, I3; PEAD(D2, 512) NOL	006150
IF (VAL, EU, NSINF, DR, NRL, EQ, NFRP) GO TO 505 HAITE (VM, 506) NSINF, NFRP	005170
506 FURMAT(* THE RESERVUIR MUST BE AT *,/, + * SEC.NU. *,13,* UR *,13)	006190
GU TU 503 505 HHITE(NH,507)	006210 006220
READ (NR2,508) HK,RX	006230
500 FURAL(2F9.5) MHITE(NH,456) MADING (2F) DELTA	006250
WRITE(NR, 458) HEAD(NR2, 458) TIMES	006280
534 WHITE(NW, 462) 	006300
IF (11MEF.GE, TIMES) GO TO 535 WRITE (NN, 532) TIMEF, TIMES	006312 006314
535 TI=TIMES	006316
IF (NRL, NE, NSINF) GO TO 509	006330
1F8=1 00 511 1=1.200	006355
CALL INTERI(QI,TI,NPHYD,TT,WK3(I)) **5(1)=TT	006370 006380
IF (IT, GT, TIMEF) GO TO 512	005390
511 CUNTINUE	006420
509 NANEO DU 513 1=1,200	008430
CALL INTERI (WK3, WK5, NN, TT, WK4(1)) IF(TT.GT.TIMEF) GO TO 514	006450 006460 006470
	006480 006490
513 CONTINUE 514 NN=NNN	006500
DU 515 I=1,NN	000250
	006540 006550
512 D1=0LLTA 1 F(1)U_F(0,2) G0 T0 515	005570 Q
DT#DELTA+3640	006590

DU 525 I=1,NUS JJ=NQS+1=I	006601
525 CUNII/UE (10/0/0/0/0/0/0/0/0/0/0/0/0/0/0/0/0/0/0/	006603 006604 006605 006605
D0 517 1=1,NRP15	006610
W=(us(JJ,2)-GS(JJ,1))*FLOAT(I=1)/NWK3+QS(JJ,1)	006615
STOK=HK*(U*AHX	006620
WK1(1)=0,5*U	006630
<pre>WK2(I)=0,5+0+STOR/DT</pre>	00650
517 wh6(I)=x7(I)=0	00650
wh4(I)=xk5(I)	00660
CALL KINRUT(wK2,WK6,NRPTS,WK1,WK7,NWK2,WK8,WK9,1,1,0DT,	006670
+NF. WK3, WK4)	006680
MHITE(NW,518) NRL	006590
518 FURMAT (//, * HYDRUGRAPH AFTER RESERVOIR AT SEC. *, I3,/,	006700
+ * TIME FLOWRATE *)	006710
DU 519 I=1,KN	006720
TI=FLOAT(I=1) #DELTA+TIMES	006730
wK3(I)=*K4(I)	006740
519 mNIIE(Nm,474) TT,WK3(I)	006750
290 \$TOP END	006760 006765 006780

# TYPICAL OUTPUT FROM FINDIF

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FLOOD ROUTING USING THE EXPLICIT FINITE DIFFERENCE	METHOD	•
CHANNEL PROPERTIES AND FLOW-CONDITIONS		•
TOTAL LENGTH = 50000.0 FT MAXIMUM DEPTH = 20.0 FT SI OPF = 00020 FT MANNINGS N = 16675 7 CU CFLEPITY AT MAX = 33.72FT/ WITH OF CHANNEL = 100.0 FT	VFT SECs SEC	
PROBLEM VARTABILS NUMBER OF REACHES = 25 TIME INCREMENT = 29.7 SE WAX TIME INCREMENT = 59.3 SE TEST NUMBER = 150	CONDS CONDS	,
INVERT ELEVATIONS AND INITIAL WATER LÉVELS AT INTIIAL FLOW RATE 3335.1 CU CHAINAGE INVERT ELEVATION 0.0 110.0 2000.0 109.6 4000.0 109.2	EACH SECTION SECS WATER LEVEL 0 117.00 0 16.60 0 16.20	
8000.0       104.4         10000.0       104.4         12000.0       104.4         12000.0       104.4         12000.0       104.4         12000.0       107.2         14000.0       107.2         16000.0       106.4         18000.0       106.4         18000.0       106.4	1         5         6           0         1         15         6           0         1         15         6           0         1         14         6           0         1         14         6           0         1         13         6           0         1         13         6           0         1         13         6	· ·
$\begin{array}{c} 220000 \\ 24000 \\ 24000 \\ 26000 \\ 26000 \\ 0 \\ 104.6 \\ 3000 \\ 104.6 \\ 3000 \\ 104.0 \\ 32000 \\ 0 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 0 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 103.2 \\ 36000 \\ 1000 $		۲.
3800.0       102.0         4000.0       102.0         42000.0       101.6         46000.0       101.2         46000.0       101.2         46000.0       100.8         48000.0       100.8         30000.0       100.4	109.40 109.60 108.60 108.70 108.70 107.80 107.80 107.80 107.80	
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a service a service of the service o

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OUTPUT-FINDIF/4

THE FHROR IN THE VOLUME WAS= CENTRUID OF HYDROGRPH NO 1= CENTRUID OF HYDROGRPH NO 21=	.022 PER CENT 2.64 HRS 3.19 HRS		•	
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		•		
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OUTPUT-FINDIF/5

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# TYPICAL OUTPUT FROM KINDIF

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		E METHOD	D ROUTING USING THE KINEMATIC WAY	FLOOD R
<u>_</u>		1TTONS	CHANNEL_PROPERTIES_AND_FLOW COND	Сн
-		5000000 000000 001497 16675.7 3.721 100.0	TOTAL LENGTH = NAXTHUM DEPTH = SLOPE = NANNTHISS N = NAXIMUM FLOW = CELEPITY AT MAX = WIDTH OF CHANNEL =	
		2n0-n 74-2 220 9-nun 1-nun	Devoluted Application Contraction Contract	P 9
	CTION WATER LEVEL 117.00 116.50	TFR LEVELS 3335-1 EPT ELEVAT	INVERT ELEVATIONS AND TRITIAL WA INITIAL FLOW PATE CHAINSS INV 250 - 4	I۷
	116.07 115.59 115.07 114.59 114.07 114.67	10 10 10 10	760°.0 770°.0 1000°.0 1250°.0 1500°.0	
-	113.00 112.50 112.50 112.50 111.50 111.50 111.50		2 jún ° ° ° 225 j ° ° ° 225 j ° ° ° 25 j ° ° ° 25 j ° ° °	
	110.57 110.67 119.57 139.57 139.57 158.57 158.57	10 10 10 10	1250 3590 3750 37550 470 4250 50 4250 50	•
	107.50 107.00	ÎŬI 101 RVF W4S NO	47645 5000°+0 The stage discharge cu	
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## OUTPUT-KINDIF/2

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OUTPUT - KINDIF/4

- ----•. NOTE: THIS IS ONLY ONE OF TWENTY-FIVE SIMULATIONS PRODUCED BY A SINGLE EXECUTION OF "KINDIF." . ÷ . . . 353

# TYPICAL OUTPUT FROM RIVER3

SUPPLY NO OF SECTIONS... IS 11 SUPPLY MAX NO DE PTS...IS DO YOU WANT LIST OF COMMANDEN...YEE/ND YES AFTER INVITATION TO TYPE ":-" GIVE DNE DF THE FOLLOWING COMMANDO.... DISCHARGE....TO SPECIFY FLOW D/S WL.....TO DEFINE DOWNSTFEAM CONTPOL LEVEL INFLOW.....TO DEFINE INFLOW HYDROGRAPH RESISTANCE... 10 SET FLOW RESISTANCE LAW OLD SECTION..TO PRINT COORDS OF A SECTION NEW SECTION. TO REDEFINE COUPLS OF A SECTION OLD COEFF....TO PRINT ROUGHNESS MEASURE NEW COEFF....TO PEDEFINE POUGHNESS MEACURE CRITIC.....TO COMPUTE CRITICAL DEPTH AT A SECTION CHANGES.....TO PRINT CHANGES OF COOPDS OF ROUGHNESS TABLE.....TO PRINT TABLE OF ALL SECTIONS DATA PROFILES..... TO PRINT OUT SURFACE PROFILES COMPUTE.....TO COMPUTE PROFILES ROUTE.....TO ROUTE THE FLOOD RESERVOIR.... TO ROUTE THROUGH A RESERVOIR RESTART.....TO BEGIN AGAIN HELP.....FOR COMMAND OPTIONS STOP.....TO TERMINATE RESISTANCE : --SPECIFY RESISTANCE LAW BY TYPING CHEZY, MANNING, STRICKLEP, COLEBROOK SMOOTH OP ROUGH....P.S. IS YOUR ROUGHNESS MEASURE COMPATIBLES :--MANNING :--DISCHARGE ENTER MOLOF SUB-REACHES (13) WITH DIFFEPENT DISCHAPGES 1 SUPPLY UPSTREAM SEC.ND. (13) AND DISCHARGES (2F9.3) FOR EACH SUB-REACH SUPPLY THE LOW VALUE FIRST REACH NO. 1 1 3000. 16675. ENTER NO. OF PTS IN FUNTION CURVES (13) 5 COMPUTE :--DEFINE SEC. NO. WHERE DVS WATER LEVEL TO BE SPECIFIED. 13 11 DEFINE NUMBER OF FOINTS IN DAS WE PATING CUPVE (13) DEFINE D/S WL AND FLOWPATE 2F9.3 START AT THE LOWEST WATER LEVEL PT. 1 7 107. 3000. PT. 2 7 113. 10000. 3.7 PT. 129. 16675.

## OUTPUT - RIVER3/2

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:- COMPUTE SPECIFY UPSTREAM LIMIT OF PROFILE BY SEC. ND. (13) 3 SPECIFY DVS STAPTING POINT OF PPOFILES BY SEC. ND. (13) 9 STAGE DISCHARGE PELATION NOT DEFINED AT SECTION ND. 9 RELATION DEFINED AT SECTIONS 11 - 11 SPECIFY D/S STAPTING POINT OF PPOFILES BY SEC. NO. (13) 11 COMPUTE :---SPECIFY UPSTREAM LIMIT DF PROFILE BY SEC. ND. (13) 1 SPECIFY D/S STARTING PDINT DF PROFILES BY SEC. ND. (I3) 3 :- PRDFILES SPECIFY U/S SEC. ND. DF PRINTDUT (13) 1 SPECIFY D/S SEC. ND. DF PRINTDUT (13) 11

SEC. ND.	DISTRACE	DISCHARGE	E WATER LEVEL	ENEPGY LEVEL
1	0.0	3000.000 1	116.564	116.888
	ł	6418.750 1	20.627	121.193
	Ģ	3837.500 1	23.977	124.746
-	1.	3256.250 1	27.846	127.985
•	10	6675.000 t	130.000	131.079
2	5000.0 0	3000.000 1	15.565	115.890
	i i	6418.750 1	19.618	120.185
•	0	9837 <b>.</b> 500 t	122.953	123.725
	10	3256.250 1	126.025	126.967
	- 1 6	6675.000 1	129.000	130.079
3	10000.0	3000.000 1	14.568	114.892
	6	5418.750 1	18.606	119.174
	0	9837.500 1	121.921	122.697
	10	3256.250 1	25.000	125.945
	1 8	6675.000 i	128.000	129.079
4	15000.0	3000.000 1	13.573	113.896
	6	6418.750 1	17.588	118.159
		9837.500 1	20.881	121.661
	10	3256.250 1	23.970	124.917
	16	6675.000 1	27.000	128.079

# OUTPUT - RIVER3/3

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5	20000.0	3000.000 6418.750 9837.500 13256.250 16675.000	112.581 116.563 119.828 122.922 126.000	112.904 117.137 120.614 123.884 127.079
6	25000.0	3000.000 6418.750 9837.500 13256.250 16675.000	111.595 115.529 118.759 121.885 125.000	111.916 116.106 119.553 122.842 126.079
7	30000.0	3000.000 6418.750 9837.500 13256.250 16675.000	110.618 114.480 117.669 120.828 124.000	110.937 115.062 118.474 121.791 125.079
8	35000.0	3000.000 6418.750 9837.500 13256.250 16675.000	109.658 113.410 116.550 119.756 123.000	109.973 114.000 117.368 120.728 124.079
à	40000.0	3000.000 6418.750 9837.500 13256.250 16675.000	108.725 112.308 115.389 118.667 122.000	109.034 112.910 116.228 119.649 123.079
10 ,	45000.0	3000.000 6418.750 9837.500 13256.250 16675.000	107.832 111.159 114.170 117.556 121.000	108.132 111.779 115.036 118.551 122.079
11	50000.0	3000.000 6418.750 9837.500 13256.250 16675.000	107.000 109.930 112.861 116.415 120.000	107.285 110.579 113.769 117.428 121.079

### OUTPUT-RIVER3/4

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:--POUTE THE INFLOW HYDROGRAPH IS NOT DEFINED · : : INFLOW HOW MANY POINTS DEFINE THE INFLOW HYD?..(I3) Э ARE THE UNITS HOURS OF SECONDS? SECONDS. AFTER EACH "?" ENTER TIME AND FLOWRATE..(2F9.3) PT. 1 ? 3335. 1500. 2 ? ΡT. 6500. 16675. PT. 3 î 11500. 3335. SPECIFY THE SEC NO. AT WHICH INFLOW IS DEFINED...(I3) 1 :- ROUTE THE WEIGHT FACTORS ARE ASSUMED TO BE 0.5 DO YOU WISH TO CHANGE ANY? (YES/ND) NO DEFINE U/S SEC. NO. WHERE POUTING STARTS (IS) 1 DEFINE D/S SEC. NO. WHERE POUTING FINISHES (13) 9 DEFINE TIME STEP FOR COMPUTATION. USE 200. SAME UNITS OF TIME AS USED FOR HYDROGRAPH (F9.4) ENTER START TIME WITH CONSISTENT UNITS (F9.3) 0.0 ENTER FINISH TIME WITH CONSISTENT UNITS, (F9.3)15000.

HYDROGRAPH AT	SEC. NO.	9
TIME	FLOWRATE	
0.000	3335.000	
200.000	3335.000	
400.000	3335.000	
600.000	3335,000	
800.000	3335.000	
1000.000	3335.000	
1200.000	3335.000	
1400.000	3335.000	
1600.000	3335.796	
1800.000	3327.287	
2000.000	3359.838	
2200.000	3317.736	
· 2400.000	3282.170	
2600.000	3389.122	
2800.000	3417.854	
3000.000	3321.493	
3200.000	3209.806	
3400.000	3230.446	
3600.000	3339.303	
3800.000	3462.927	
4000.000	3520.652	
4200.000	3486.220	
4400.000	3374.778	
4600.000	3246.841	
4800.000	3147.967	
5000.000	3067.595	

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5200.000	3030.879
5400.000	3000.000
5600.000	3000.000
5800.000	3000.000
6000.000	3226 843
6200.000	3848 215
6400.000	4647 202
6400.000 6600.000	4041.200 EEEA 332
6600.000	2000.326
5000.000 7000 000	04(C.C(0 7405 071
7000.000	7400.371
7200.000	0676.111
7400.000	0741.74C 0405 425
7000.000	242J.46J 10207 104
(000.000	10307.126
0000.000 0000 000	11/100.271
6200.000 9400 000	11412.010
0400.000	11/42.2/2
0000.000	10005 000
8800.000 8800.000	12783.708
9000.000	13878.784
9200.000	14707.358
9400.000	10074.040
7600.000	10801.727
9800.000 10000 000	10000.710
10000.000	15046.010
10200.000	15556 072
10400.000	15179 402
10200.000	14752 209
11000.000	14297 967
11200.000	13816.157
11400.000	13270.195
11600.000	12733.183
11800.000	12221 965
12000.000	11759.379
12200.000	11241.673
12400.000	10788.580
12600.000	10411.226
12800.000	9854.774
13000.000	9194.578
13200.000	8653.719
13400.000	8267.276
13600.000	7981.947
13800.000	7720.165
14000.000	7437.375
14200.000	7127.168
14409.000	6768.503
14600.000	6352.529
14800.000	5926.439
15000.000	5468.961

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## OUTPUT-RIVER3/6

2---RESERVOIP SPECIFY LOCATION OF RESERVOIR (13) 1 DR AT SEC. NO. 9 9 ENTER K.X, DESCRIBING THE RESERVOIR (2F9.3) 10000. 1.1 DEFINE TIME STEP FOR COMPUTATION. USE 200. SAME UNITS OF TIME AS USED FOR HYDROGRAPH (F9.4) ENTER START TIME WITH CONSISTENT UNITS (F9.3) 1800. ENTER FINISH TIME WITH CONSISTENT UNITS, (F9.3)15000.

#### HYDPDGRAPH AFTER PESERVOIR AT SEC. 9

TIME FLOWRATE 3327.287 1800.000 3327.414 2000.000 3327.503 2200.000 2400.000 3327.288 2600.000 3327.353 3327.946 2800.000 3000.000 3328.271 3200.000 3327.784 3326.945 3400.000 3600.000 3326.617 3327.198 3800.000 3328.480 4000.000 4200.000 3329.843 4400.000 3330.627 4600.000 3330.472 3329.436 4800.000 3327.709 5000.000 5200.000 3325.540 5400.000 3323.124 3320.607 5600.000 3318.109 5800.000 3316.515 6000.000 6200.000 3318.237 3325.477 6400.000 6600.000 3339.291 6800.000 3360.106 3387.985 7000.000 7200.000 3422.736 3463.202 7400.000 7600.000 3508.038 3557.841 7800.000 8000.000 3613.138

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8200.000 3	3672.337	
8400.000 3	3733.976	
8600.000 3	3798.373	
8800.000 3	3867.082	
9000.000 3	3941.519	
9200.000 4	4022.080	
9400.000 4	\$108.008	
9600.000 4	4197.724	
9800.000 4	1289.343	
10000.000 4	\$381.007	
10200.000 4	4471.140	
10400.000 4	4558.660	
10600.000 4	1642.863	
10800.000 4	1723.283	
11000.000 4	1799.644	
11200.000 4	4871.759	
11400.000 4	1939.310	
11600.000 5	5002.116	
11800.000 5	5060.350	
12000.000 5	5114.337	
12200.000 5	5164.085	
12400.000 5	5209.665	
12600.000 5	6251.655	
12800.000 5	5289.681	
13000.000 5	5322.672	
13200.000 5	5350.727	
13400.000 5	5374.952	
13600.000 5	5396.372	
13800.000 5	5415.494	
14000.000 5	5432.346	
14200.000 5	5446.757	
14400.000 5	5458.451	
14600.000 5	5467.036	
14800.000 5	5472.274	
15000.000 . 5	5474.030	
':- ROUTE		
THE WEIGHT FACTO	JFS ARE ASCUMED TO BE 0.5	
DO YOU WISH TO C	HANGE ANY? (YESZNO) NO	
DEFINE UVS CEC.	ND. WHERE POUTING STRPTS (I3) 9	
DEFINE DVS SEC.	ND. WHEPE POUTING FINISHES (I3) 11	
DEFINE TIME STEP	P FOR COMPUTATION. USE	
SAME UNITS OF TI	IME AS USED FOR HYDROGRAPH (F9.4)	200.
ENTER STAPT TIME	E WITH CONSISTENT UNITS (F9.3) 6400.	
ENTER FINISH TIM	1E WITH CONSISTENT UNITS, (F9.3)15000.	

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# OUTPUT-RIVER3/8

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HYDROGPAPH AT	SEC. ND.	11
TIME	FLOWRATE	
6400.000	3325.477	
6600.000	3327.819	
6800.000	3321.889	
7000.000	3317.980	
7200.000	3319.483	
7400.000	3327.469	
7600.000	3342.786	
7800.000	3366.003	
8000.000	3395.695	
8200.000	3430.487	
8400.000	3471.075	
8600.000	3517.936	
8800.000	3569.884	
9800.000	3625.135	
9200.000	3682.877	
9400.000 0400.000	3743.678	
9600.000 0000 000	3808.780 2878 882	
7800.000 10000 000	0017.70E 00E7 171	
10000.000	3737.171	
10200.000	4040.078	
10400.000	4161.316	
10500.000	4217.004	
10800.000	4307.244	
11000.000	4400.103 AA09 915	
11200.000	4407.010	
11400.000	4658 855	
11800.000	4738.455	
12000.000	4813.923	
12200.000	4884.808	
12400.000	4951.156	
12600.000	5013.130	
12800.000	5070.385	
13000.000	5123.317	
13200.000	5172.834	
13400.000	5218.654	
13600.000	5259.752	
13800.000	5295.676	
14000.000	5326.839	
14200.000	5354.114	
14400.000	5378.271	
14600.000	5399.791	
14800.000	5418.942	
15000.000	5435.648	

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## OUTPUT - RIVER3/9

FROUTE THE WEIGHT FACTORS ARE ASSUMED TO BE 0.5 DO YOU WISH TO CHANGE ANY? (YES/ND) NO DEFINE U/S SEC. NO. WHERE ROUTING STARTS (I3) 1 DEFINE D/S SEC. NO. WHERE ROUTING FINISHES (I3) 11 DEFINE TIME STEP FOR COMPUTATION. USE SAME UNITS OF TIME AS USED FOR HYDROGRAPH (F9.4) ENTER START TIME WITH CONSISTENT UNITS (F9.3) 1400.0 ENTER FINISH TIME WITH CONSISTENT UNITS, (F9.3)15000.

200.

HYDROGPAPH AT	SEC. NO.	11
TIME	FLOWRATE	
1400.000	3335.000	
1600.000	3335.135	
1800.000	3333,147	
2000.000	3344.816	
2200.000	3312.288	
2400.000	3344.725	
2600.000	3382.342	
2800.000	3288.979	
3000.000	3272.541	
3200.000	3364.068	
3400.000	3441.552	
3600.000	3379.674	
3800.000	3267.358	
4000.000	3201.949	
4200.000	3237.249	
4400.000	3344.666	
4600.000	3463.018	
4800.000	3519.677	
5000.000	3486.291	
5200.000	3410.329	
5400.000	3298.021	
5600.000	3207.041	
5800.000	3120.187	
6000.000	3067.402	
6200.000	3028.419	
6400.000	3010.347	
6600.000	3000.000	
5800.000 7000 000	3000.000	
7000.000	31/7.5/4	
7200.000	4025.357	
7400.000	H717,004 E002 021	
7800.000 7800 000	7705.001 7705 275	
2000.000	7776 079	
8280 880	8429 149	
0000.000	ショレン・エキン	

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### OUTPUT - RIVER3/10

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8400.000
                9241.573
               10001.907
   8600.000
               10486.739
   8800.000
               10824.222
   9000.000
   9200.000
               11162.359
   9400.000
               11659.443
   9600.000
               12359.981
   9800.000
               13206.846
               14065.501
  10000.000
  10200.000
               14847.776
               15453.160
  10400.000
               15829.071
  10600.000
  10800.000
               15987.072
               15960.140
  11000.000
               15784.936
  11200.000
  11400.000
               15497.727
               15174.797
  11600.000
  11800.000
               14765.795
  12000.000
               14289.758
  12200.000
               13744.334
               13259.507
  12400.000
  12600.000
               12741.053
  12800.000
               12156.478
               11746.403
  13000.000
  13200.000
               11369.349
  13400.000
               10814.897
  13600.000
               10143.846
  13860.000
                9459.336
                8892.297
  14000.000
  14200.000
                8483.380
  14400.000
                8155.858
  14600.000
                7865.744
                7581.413
  14800.000
  15000.000
                7280.501
       CRITIC
. : --
SPECIFY SECTION NO,13
                               11
ENTER DISCHAPSE
                  F9.3
                             5000.
AT SEC 1 WITH 0= 5000.000
CRITICAL WATER LEVEL= 114.266 AND
CRITICAL ENERGY LEVEL= 116.399
:---
       DLD SECTION
DEFINE SECTION NO. . IS
                               3
              3 HAS 4 PTS....CODRDS ARE
SECTION NO.
               0.0.108.0 100.0,108.0 100.0,128.0
 0.0,128.0
:--
       STOP
     STOP
      6.791 CP SECONDS EXECUTION TIME
```

#### APPENDIX H

### SUMMARY OF NOTATION

This appendix provides a dictionary of the notation used in this thesis. It should be noted that several variables serve more than one purpose. The appropriate definition can be inferred from the context and from the description provided with the variables. Where there is more than one definition of a variable, the location of the less frequent definition(s) is identified.

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Α	=	area of the cross section		
Ā	=	an error, expressed as a term of a Fourier Series		
Ao	=	a coefficient on a term of a Fourier Series		
a	=	a coefficient of regression (equation 4.12)		
B	=	invert elevation at a section		
b	=	a coefficient of regression (equation 4.12)		
С	=	celerity of a wave relative to water (Chapter 2)		
	=	kinematic wave velocity		
Cl	=	a stability condition variable		
C2	Ξ	a stability condition variable		
с	=	a coefficient of regression (equation 4.12)		
D	=	a coefficient of an error term		
е	æ	an exponent (equation 5, 26)		
Fr	=	Froude Number	1	

- g = acceleration of gravity
  - = a displacement on the distance axis from the point about which a Taylor Series is written (Appendix B, C)
- H = water surface elevation, (usually associated with a position in the space-time diagram)
- HORZ= a variable discribing the horizontal distance from an arbitrary axis to a coordinate describing the channel cross section
- h = water surface elevation
  - = a displacement on the time axis from the point about which a Taylor Series is written (Appendix B, C)
- I = a subscript to denote a position on a space-time diagram

= inflow (equations 3.21 - 3.24)

i = √-1

J = a subscript to denote a position on a space-time diagram
 K = a subscript to denote a position on a space-time diagram
 (Chapter 2, Appendix A)

- a parameter in the equation describing storage in an
   imaginary reservoir or an elementary reach
- = a coefficient (Appendix F)
- KN = Kinematic Courant Number

	k	=	a coefficient in the equation relating attenuation and the		
			number of reaches in a simulation (equation 5.26)		
		=	a coefficient in the equation relating "reduced"		
			attenuation to relative live storage (equation 5.29)		
	L	=	a subscript to denote a position on a space-time diagram		
		=	total length of the channel (Chapter 5)		
	m	=	a coefficient (Appendix A)		
	Ν	=	number of reaches in the total length of the channel		
	NL	=	a parameter which describes the "bow" in the relation-		
			ship between storage and flow rate (defined in		
			figure 6.4)		
	n	Ξ	Mannings roughness coefficient		
		=	a coefficient (Appendix A)		
•	0	=	outflow .		
	Ρ	=	wetted perimeter of a cross section (Chapter 2, figure 3.1,		
			equations 3.42 - 3.47, equations 4.1 - 4.5 Appendix F)		
		=	peak outflow divided by peak inflow or full bank flow rate		
	Q	±	flow rate		
	đ	=	flow rate per unit width		
	ą	=	rate of lateral inflow		
,	R	Ξ	hydraulic radius		
	S	1	slope (Chapter 2, figure 3. 1, equations 3.42 - 3.46,		
			Appendix F)		

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01		skew factor, a parameter which describes the slope		
		of the line from the base flow rate and the time of the		
		centroid of the hydrograph to the peak outflow and the		
		time of peak outflow (Defined in figure 6, 3)		
Sf	11	slope of the friction line		
So	=	bed slope		
ST	=	Storage in an elementary reach or reservoir		
Т	8	time (usually associated with a finite difference step)		
Ть	=	time base of the inflow hydrograph		
Τc	=	time of the centroid of the outflow hydrograph measured		
		from the time of the centroid of the inflow hydrograph		
Тр	=	time of peak outflow measured from the time of peak		
		inflow		
Τw	=	surface width		
VER:	Γ=	a variable describing the horizontal distance from an		
VERT	Γ=	a variable describing the horizontal distance from an arbitrary axis to a coordinate describing the channel		
VER!	Γ=	a variable describing the horizontal distance from an arbitrary axis to a coordinate describing the channel cross section		
VER!	Γ= =	a variable describing the horizontal distance from an arbitrary axis to a coordinate describing the channel cross section flow velocity under steady state conditions		
Ver:		a variable describing the horizontal distance from an arbitrary axis to a coordinate describing the channel cross section flow velocity under steady state conditions an exponent in the equation describing storage in an		
Ver: Vo w		a variable describing the horizontal distance from an arbitrary axis to a coordinate describing the channel cross section flow velocity under steady state conditions an exponent in the equation describing storage in an imaginary reservoir		
VER Vo w	Γ= = =	a variable describing the horizontal distance from an arbitrary axis to a coordinate describing the channel cross section flow velocity under steady state conditions an exponent in the equation describing storage in an imaginary reservoir distance (usually associated with a finite difference step)		
VER Vo w		a variable describing the horizontal distance from an arbitrary axis to a coordinate describing the channel cross section flow velocity under steady state conditions an exponent in the equation describing storage in an imaginary reservoir distance (usually associated with a finite difference step) distance		
VER Vo w		a variable describing the horizontal distance from an arbitrary axis to a coordinate describing the channel cross section flow velocity under steady state conditions an exponent in the equation describing storage in an imaginary reservoir distance (usually associated with a finite difference step) distance a parameter used in Muskinghum Flood Routing		

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- Y = vertical depth of water
- y = vertical depth of water
- yo = vertical depth of water under steady state conditions
- Z = Courant Number (Chapter 2)
  - = side slope of a triangular cross section (Chapter 3)
- α parameter which identifies the location of the nucleus
   of a molecule in the x axis. Referenced to the upstream,
   lowest time level of the molecule
- $\beta$  = a parameter which identifies the location of the nucleus of a molecule in the  $\dagger$  axis. Referenced to the upstream, lowest time level of the molecule
- **c** = error term due to finite difference approximation.
- $\gamma$  = a wave number (Appendix A)
- 2 = a parameter which identifies the location of the nucleus of a molecule on the x axis. Referenced to the centre of the molecule
- $\sigma$  = a wave number (Appendix A)
- $\Phi$  = a parameter which identifies the location of the nucleus of a molecule on the  $\dagger$  axis. Referenced to the centre of the molecule

#### APPENDIX I

To provide further clarification of the procedure that could be followed in calibrating a kinematic flood routing model. the following example is provided. The numerical values used in this example are taken from tables 3.3, 3.4 and 3.5. However. it will be assumed that these results are available only after a simulation of the particular physical system. The channel employed has been described previously in Chapter 2 as System 1. The waterway is 50,000 feet long, 100 feet wide, with a depth of 20 feet. A symetrical triangular inflow hydrograph was employed as a description of the time variant inflow. A peak outflow of 0.840 times the full bank flow was assumed. Figure 4.2 shows both the inflow hydrograph and various outflow hydrographs predicted using a numerical solution of the momentum and continuity equations. It is intended that this example will show the calibration of a kinematic model to emulate the peak outflow. A discussion of the wave shapes will be made in the conclusions of this appendix.

The procedure follows the description presented in Chapter 5 on pages 153 to 155.  $\cdot$ 

1. To describe the hydrograph adequately, a time step.

 $\Delta T$  , of 200 seconds was assumed. It should be

noted that this results in the truncation of the peak inflow to 0.984 times that specified in the inflow hydrograph.

- 2. From the results presented in Appendix F, a value of C, the kinematic wave velocity, of 12.3 feet per second is assumed. This is equivalent to the full bank flow kinematic wave velocity. (With the hindsight provided by results presented in Chapter 5. this is an accurate estimate of the effective kinematic wave velocity.)
- 3. An estimate of the size of the distance step,  $\Delta X$  can be obtained from setting

### $\Delta X \neq C \Delta T$

The first estimate of  $\Delta X$  subject to the constraint that  $\frac{L}{\Delta X}$  is an integer is:

### $\Delta X = 2500$ FEET

Performing a simulation with  $\mathbf{a} = 0.0$  and  $\mathbf{\beta} = 1.0$ yields a peak outflow of 0.843 times full bank flow. Performing the calculations with the nucleus located in the upper right hand corner of the molecule defined in Figure 3.2, as specified above, results in the maximum amount of attenuation for the given  $\Delta \mathbf{X}$  and  $\Delta \mathbf{T}$  with a nucleus inside the molecule. 4. Because the peak outflow does not correspond with the objective. it is necessary to modify the value of  $S\Delta X$ . To provide an estimate of the next value of  $S\Delta X$  to use in an attempt to achieve the desired attenuation, the following steps may be taken. First, calculate the value of  $S\Delta X$  for the completed simulation.

$$S = (2\alpha - 1) + (1 - 2\beta) C \frac{\Delta T}{\Delta X}$$
$$= (-1) + (-1) 12.3 \frac{200}{2500}$$
$$= -1.984$$

Thus:

### SAX = -1.984 × 2500 = - 4960 FEET

Assuming that no attenuation occurs with  $S\Delta X = 0.0$ , a straight line may be plotted through the points

P = 1.0,  $-S\Delta X = 0.0$ 

and

 $P = 0.843, -S\Delta X = 4960$ 

This straight line may be extended to obtain the next estimate of  $S\Delta X$ , approximately 5,000 feet.

It is not possible to increase the value of  $S\Delta X$  by varying  $\alpha$  and  $\beta$  subject to the constraint that the n ucleus must remain inside the molecule with the previously defined parameters. Thus, it is necessary to vary the values of  $\Delta X$  and/or  $\Delta T$ . For a system of arbitrary geometry with sections defined at regular or irregular intervals, modification of  $\Delta X$  may be impractical due to the fact that the system definition will be significantly altered. Thus, the variation of  $\Delta T$  may prove to be the most fruitful alternative in altering S. The limiting constraints on the selection of  $\Delta T$  (and  $\Delta X$ ) are the discretization error and the stability requirements.

5. For the second interation, the simulation was repeated with  $\alpha = 0.0$ ,  $\beta = 1.0$ ,  $\Delta X = 5,000$  feet and  $\Delta T = 200$  seconds. This results in a peak outflow of 0.804 times the full bank flow. Clearly this is an over estimate of  $S\Delta X$ . For this simulation:

 $S\Delta X = ((2\alpha - 1) + (1 - 2\beta) C \frac{\Delta T}{\Delta X}) \Delta X$ =  $((-1) + (-1) 12.3 \frac{200}{5000}) 5000$ 

= -7460 FEET

6. The above simulation defines another point on the which could be used as a guide for further refinements. The third simulation is attempted with  $S\Delta X = -5,000$  feet. Again it should be noted that there are four variables that may be used to modify  $S\Delta X$ :

### α, β, ΔΧ, ΔΤ

а	β	ΔΧ	ΔΤ
0,00	0.5	5, 000	variable*
0.246	1.0	5.000	200
0.25	0.5	10, 000	variab <b>le</b> *
0.374	1. 0	10, 000	200

For example a few of the points are:

Two of the above points, each identified with an asterisk, are shown in tables 3.4 and 3.5. Interpolation to obtain the 0.840 contours on these tables will give an indication of the interelation of the various parameters which define  $S\Delta X$ . The choice of the appropriate values of  $\alpha$  $\beta$ ,  $\Delta X$  and  $\Delta T$  remains at the discretion of the user who must observe the constraint of stability and discretization error. Figure 5.8 shows a comparison of the kinematic simulation performed with

$$\alpha = 0.0$$
  
$$\beta = 0.5$$
  
$$\Delta X = 5,000 \text{ feet}$$
  
$$\Delta T = 200 \text{ seconds}$$

and a dynamic solution. The two hydrographs have the same general shape. The peak of the dynamic simulation is slightly higher than the peak of the kinematic solution. Plotting the first simulation presented in this appendix on the graph may yield a result which agrees more closely with the dynamic solution.

In certain instances, there may be significant differences in the shapes of the two hydrographs. For such cases further research to determine methods of modifying the shape of the hydrograph predicted by the kinematic solution would be required. Preliminary results presented in Chapter 3 on figures 3, 5 and 3, 6 indicate that the hydrograph is insensitive to the variables used to define  $S\Delta X$ 

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