

EFFECT OF FLOOR SLABS AND FLOOR BEAMS
ON STATIC AND DYNAMIC BEHAVIOUR OF
SHEAR WALL STRUCTURES

by

JAYANTA K. BISWAS, M.E.

A Thesis

Submitted to the Faculty of Graduate Studies
in Partial Fulfilment of the Requirements
for the Degree
Master of Engineering

McMaster University

November 1970

MASTER OF ENGINEERING (1970)
(Civil Engineering)

McMASTER UNIVERSITY
Hamilton, Ontario

TITLE: Effect of Floor Slabs and Floor Beams on Static
and Dynamic Behaviour of Shear Wall Structures.

AUTHOR: Jayanta K. Biswas, M.E. (Calcutta University)

SUPERVISOR: Dr. W.K. Tso

NUMBER OF PAGES: vi, 170

SCOPE AND CONTENTS:

This thesis studies the effect of floor slabs on the static and dynamic behaviour of the shear wall structure. A single component has been analysed using the 'Matrix Transfer' technique along with Vlasov's thin walled elastic beam theory. Experimental verification was done on a small scale plexiglas eight storey model in the form of a channel section for both static and dynamic loading.

The thesis also deals with the analysis of the non-planar shear walls coupled through floor beams subjected to static loading. The continuum approach along with Vlasov's theory has been used in the analysis. Experimental verification was done on a small scale plexiglas model in the form of two equal angles connected by eight floor beams at equal spacing.

ACKNOWLEDGEMENTS

I wish to express my profound gratitude to Dr. W.K. Tso for his continued guidance and encouragement in carrying out this thesis.

I would also like to thank Dr. A.C. Heidebrecht for arranging discussions among shear wall group in order to exchange views and ideas.

I thank Dr. A. Ghobarah and Mr. R.K. Raina for their co-operation. Thanks are also due to Mrs. S. McQueen for neatly typing this thesis.

This investigation was made possible through the financial assistance of the National Research Council, to whom I extend my sincere thanks.

TABLE OF CONTENTS

Chapter		Page
1	INTRODUCTION	
	1.1 Description of Shear Wall	1
	1.2 Shear Wall Project	4
	1.3 Review of Part Work	4
	1.4 Present Investigation	8
2	STATIC STUDY OF SHEAR WALL WITH FLOORS	
	2.1 Summary	10
	2.2 Matrix Transfer Method	10
	2.3 Application of Matrix Transfer Method	14
	2.3.1 State Vector	16
	2.3.2 Field Transfer Matrix	16
	2.3.3 Action of Floor Slab	20
	2.3.4 Point Transfer Matrix	22
	2.3.5 Boundary Conditions	23
	2.4 Computer Program	25
	2.5 Experiment	25
	2.6 Results and Discussions	31
3	DYNAMIC STUDY OF SHEAR WALL WITH FLOORS	
	3.1 Summary	49
	3.2 Matrix Transfer Method	49
	3.3 Theoretical Analysis	50
	3.3.1 Field Transfer Matrix	50
	3.3.2 Point Transfer Matrix	55
	3.3.3 Boundary Conditions	58
	3.4 Computer Program	60
	3.5 Experiment	60
	3.6 Results and Discussions	69
4	STATIC FORMULATION OF NON PLANAR COUPLED SHEAR WALL	
	4.1 Summary	83
	4.2 Theoretical Analysis	83
	4.2.1 Notations Used	85
	4.2.2 Geometric Relations	87
	4.2.3 Displacement Consideration	89
	4.2.4 Force Equilibrium Conditions	91
	4.2.5 Differential Equations	96
	4.2.6 Boundary Conditions	99
	4.2.7 Solution	100
	4.2.8 Special Configurations	103
	4.2.9 Effect of Neglecting Axial Deformation of Piers	106

Chapter		Page
4	4.3 Computer Program	108
	4.4 Experiment	108
	4.5 Results and Discussions	113
5	CONCLUSIONS AND RECOMMENDATIONS	
	5.1 Conclusions	125
	5.2 RECOMMENDATIONS	127
Appendix A	Vlasov's Theory of Thin Walled Beam	130
Appendix B	Computer Programs	
	Static Analysis of Shear Wall with Floors	134
	Dynamic Analysis of Shear Wall with Floors	139
	Non Planar Coupled Shear Wall	146
Appendix C	Experimental Data	151
BIBLIOGRAPHY		167

NOTATIONS

The following symbols are used throughout this thesis without further definition. Other symbols are defined when used.

E	Modulus of Elasticity
G	Modulus of rigidity
ν	Poisson's ratio
X, Y	Orthogonal axis
Z	Vertical axis
ω	Principal Sectorial co-ordinate
A	cross sectional area
I_X, I_Y	Moment of Inertia about X and Y axes
J	Torsional rigidity
I_p	Polar moment of inertia about shear center
I_ω	Sectorial moment of inertia
θ	Rotation about Z axis
θ'	Rate of change of rotation about Z axis
u, v	Displacement of shear center in x and y directions
u', v'	Slope in zx and zy planes
V	Shear Force
M	Moment
H	Torque
B	Bimoment
N	Axial Force
a_x, a_y	Co-ordinate of shear center
ρ	Mass density of material
w	Natural frequency
$(\dot{})$	Differentiation with respect to time
$()'$	Differentiation with respect to space.

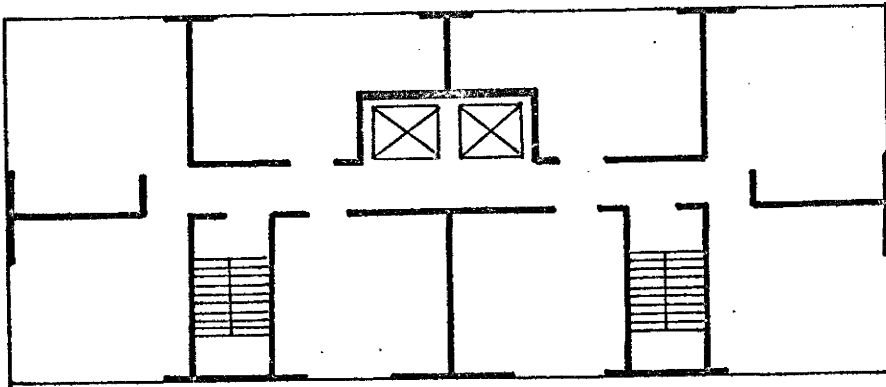
CHAPTER I
INTRODUCTION

1.1 Description of Shear Wall

As buildings increase in height, it becomes necessary to ensure adequate lateral stiffness. This stiffness may be achieved in various ways of which the use of shear wall is very common and popular.

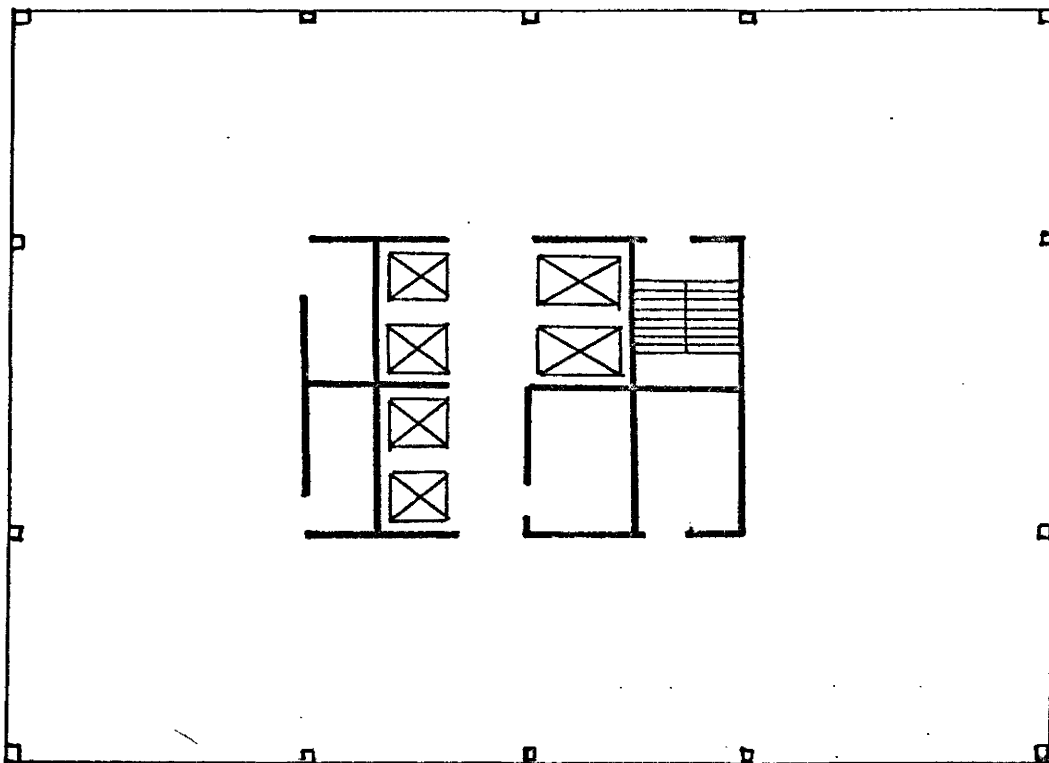
Although shear walls can be arranged in a building in innumerable ways, they can be broadly classified into two basic types. In an apartment building, shear walls are used alone and located on both sides of the corridor as shown in Fig. 1.1.1. In an office building, shear walls are located in the center to form a service core for staircases, elevators etc. This core is surrounded by a structural framing which are interconnected as shown in Fig. 1.1.2. In both the above types, shear walls serve the multipurpose function of supporting vertical and lateral loads, acting as partition walls and serving other useful functions.

Shear walls are normally interconnected by floor slabs at each floor level. These floor slabs act as highly rigid diaphragm in their own plane and bend and twist out of plane. Therefore, the slabs transmit and distribute lateral loads among the walls and also provide some resistance



Typical Apartment Building With Shear Wall

FIG. 1.1.1



Typical Office Building With Shear Walls
in the central Core Surrounded by
Structural Framing

FIG. 1.1.2

to the deformation of the walls. The effect of overall interaction between the walls and the floors is to increase the lateral stiffness of the building and to reduce stress level in the walls.

Very often shear walls are pierced to provide openings for doors, windows or corridors. The arrangement may be thought of as two or more sets of walls connected by beams. These beams resist the deformation of the wall and increase the stiffness of the assembly.

Frequently the section of shear walls are in the form of open thin walled sections. Such beams are distinguished from solid beams by experiencing longitudinal stress as a result of torsion due to warping. Appropriate theory should be considered for dealing with such sections.

For these reasons, a complete analysis of a building as shown in Figs. 1.1.1 and 1.1.2 is a most complex problem encountered in structural engineering practice. The complexity is due to various interacting elements. The dynamic analysis is even more complex. Approximate design methods, neglecting complex interaction can be used for proportioning elements which often under estimates the stiffness of building. Therefore, more sophisticated techniques of analysis are required.

The general purpose of research on shear wall structure is firstly to understand fully the behaviour of different elements and secondly, to develop more realistic methods of

analysis.

1.2 Shear Wall Project

The Canada Emergency Measures Organization is sponsoring an extensive program into behaviour of shear wall building. This project is conducted in the Department of Civil Engineering and Engineering Mechanics at McMaster University. The experimental part of the project consists of building small scale shear wall structures and studying their response due to static and dynamic lateral loadings. The theoretical part of the project consists of developing theories to explain the behaviour of shear wall structures, comparing theoretical results with experiments and developing simplified design method.

Tests have been carried on an eight feet model having E shaped section and made of non-reinforced micro-concrete. Afsar(1), Quareshi(2), Speirs(3), Raina(4) and Swift(5) studied different aspects of behaviour of shear wall structures.

1.3 Review of Past Works

Coull and Smith (6) compiled a comprehensive summary of the published literature concerning shear wall buildings.

Winokur and Gluck (7) developed a method to form lateral stiffness matrix of asymmetric building by combining lateral stiffness matrix of each element. Transverse stiffness of slab and warping torsional stiffness of individual

elements has been neglected.

The phenomenon of warping has been known to the aeronautical engineers for a long time but its application to shear wall structure is rather recent. Vlasov (8) developed the theory of thin walled open beams. Zbirohowski-Koscia (9) presented Vlasov's theory in simpler way and with an aim to make it usable by practicing engineers.

Afsar (1) has outlined various analytical and experimental approaches used in the shear wall study.

Quareshi (2) analysed the shear wall with rows of opening by frame analogy method and also conducted experiments on small scale models.

Speirs (3) studied the behaviour of floor slabs introduced in shear wall structure. His theoretical analysis is mainly based on the initial parameter approach of Vlasov.

Raina (4) studied response of shear wall structure under dynamic loading.

Swift (5) developed computer program based on matrix method, to solve asymmetric coupled shear wall. He also developed a program to analyse shear wall with floors.

Qadeer (10) and Qadeer and Smith (11) discussed the interaction between walls and slabs in a cross wall structure. Curves are given for equivalent width of slab. Experimental

work on a model was done for verification of the theory.

Taranath (12) studied open section with and without floors. Finite element treatment for floor slab is used. Multiple open section core structure coupled through floor slab is also examined for the case of static loading.

Beck (13), Rosman (14) analysed plane coupled shear wall by continuum method. The connecting beams are replaced by independently acting laminae. Coull and Choudhury (15), (16) developed design curves for different types of loading based on the continuous method of analysis.

Choudhury (17) discussed the solution single shear wall with openings by continuous method, equivalent frame method and finite element method. The behaviour of walls interconnected through floor slab is also examined. A method of complete analysis of shear wall/frame buildings taking into account their three-dimensional behaviour is presented.

Michael (18) made torsion analysis of a core wall consisting of two equal channels tied by beams at equal spacing by the continuum approach.

Jenkins and Harrison (19) analysed tall building with shear walls under bending and torsion. His bending analysis is based on stiffness matrix approach and torsion analysis is based on the theorem of minimum potential energy. The

warping stiffness of the open sections are neglected. Experiments are carried out on small scale plexiglas model in different stages.

Holmes and Astill (20) conducted experiments on a small scale shear wall structure under simulated wind load. Comparison of experimental values are made with theoretical consideration of simplified structure using Rosman's (14) theory.

Rosman (21) presented analysis of pierced torsion boxes subjected to torsion loading, arbitrarily distributed along the height. Treatment for two channel box and four angle box is done. Determination of approximate fundamental period of torsional vibration is also included.

Gluck (22) presented a lateral load analysis by three dimensional continuous method for structures consisting of simple or coupled, prismatic or non-prismatic, shear walls and frames arranged asymmetrically in floor plan. Connecting beam on the shear wall is replaced by an 'elastic media' of known stiffness properties. Treatment for thin walled open section is included. Differential equations are obtained for three-generalised displacements. In his derivation of stiffness matrix for 'elastic media', slight inconsistency of the use of 'thin walled beam theory' was noticed. Modification of few elements of matrix has been suggested by Biswas and Tso (24) in a discussion of Gluck's (22) paper.

Macleod (23) commented on the limitation of the use of continuum method when the bending stiffness of the wall approach that of connecting beams. A criterion is developed for assessing when this effect may be important. Comparison is made with more accurate frame analysis.

Coull and Irwin (28) presented a method for the analysis of the distribution of load amongst the shear walls of a three dimensional multistorey building subjected to bending and torsion. The method is based on the continuum approach.

1.4 Present Investigation

In the second and third chapters of this thesis, particular interest is given on shear wall with floors. A shear wall structure consisting of channel section with floor slabs is analysed for static and dynamic loading using the 'Matrix Transfer' method. To the best of the author's knowledge, this method has not been used to solve similar problems before. Experimental study was conducted on a small scale plexiglas model. This model was subjected to lateral loading at different floor levels when the recorded deflections and strains were studied. It was then subjected to lateral vibration to determine the resonant frequencies. The relative strain distribution at resonance is also studied.

In the fourth chapter, particular interest is on the nonplanar shear walls coupled by floor beams where warping

due to torsion of piers is taken into account. Differential equations are obtained using the continuum approach. The present formulation is applicable to two shear walls connected by one row of beams and subjected to forces and torques distributed along its height. The experimental study was performed on a small scale plexiglas model consisting of two equal angle sections connected by beams at equal spacing. It was subjected to a force and a torque at top. The resulting strains and deflections are then analysed. A comparison of the experimental results with the theory is made in all three chapters.

CHAPTER II

STATIC STUDY OF SHEAR WALL WITH FLOORS

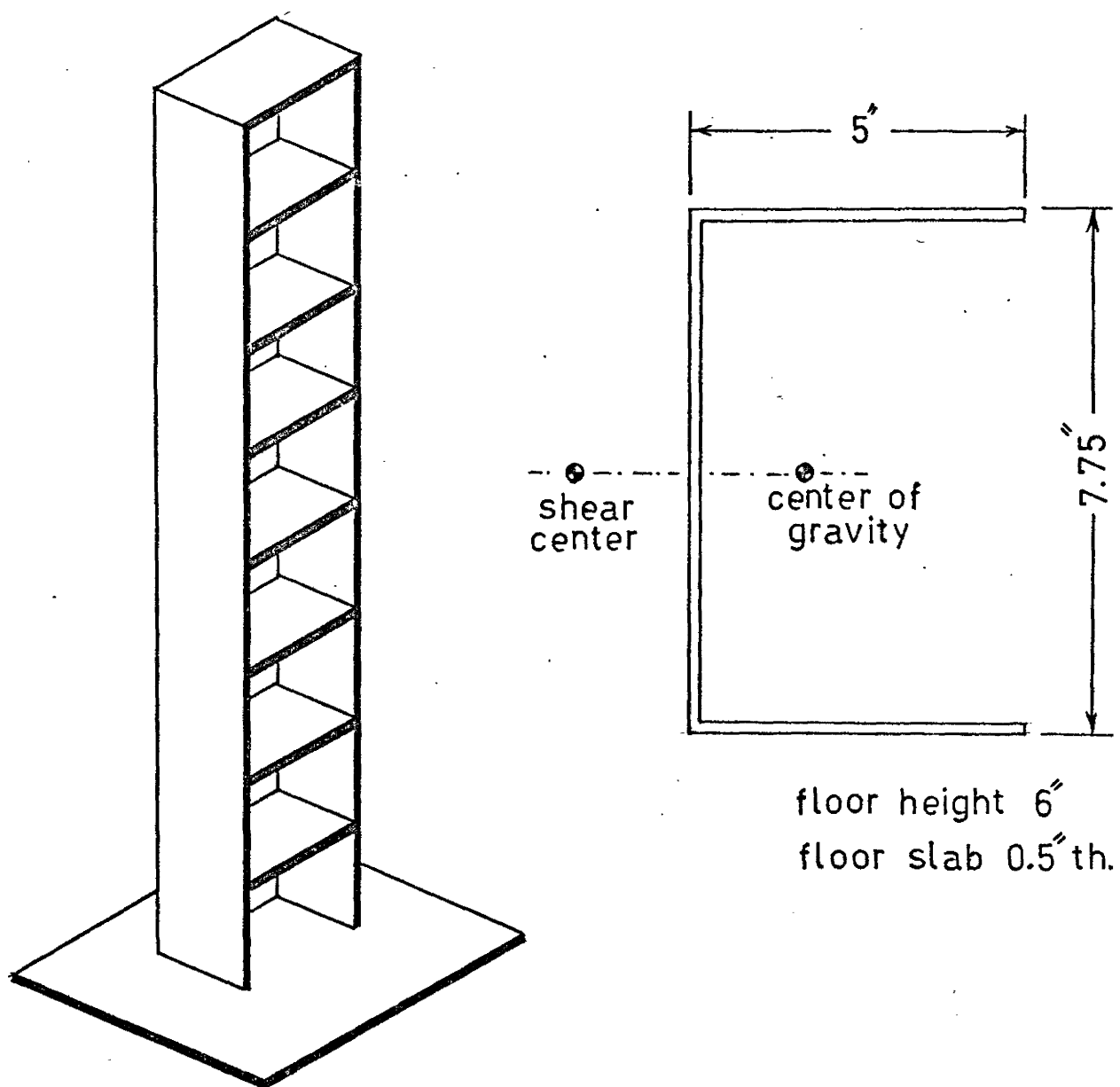
2.1 Summary

In this chapter, a shear wall structure consisting of a channel section with floor slab is analysed for static lateral load. The 'Matrix Transfer' technique is used. An experiment performed on a small scale plexiglas model Fig. (2.1.1) is described and the experimental results are compared with theoretical predictions.

2.2 Matrix Transfer Method

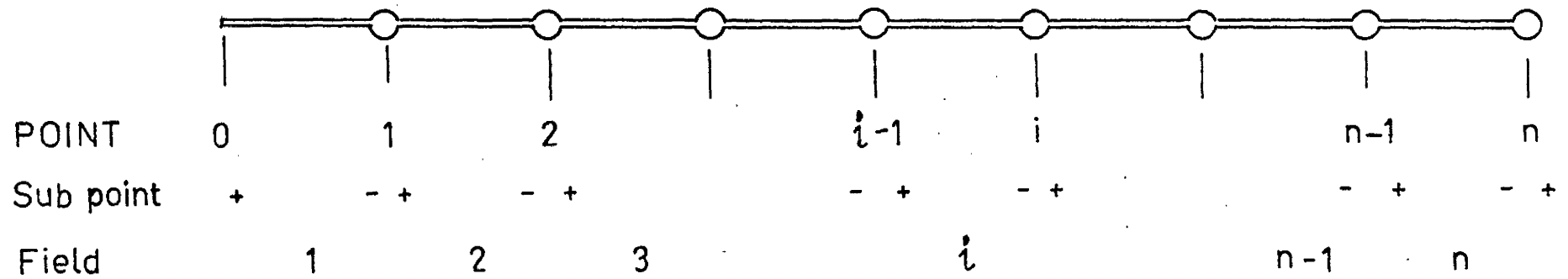
This method was originally developed by Holtzer(25) for treating torsional vibrations of shafts with lumped system. Myklestad(26) used a similar method for study of beam vibration problems. It was modified by Thomson(27) to extend its applicability to more general problems. Application of such method to static problem is less common.

Consider a system with n points and n elements along its length as shown in Fig.(2.2.1). For any point i there are two sub-point $(i)_-$ and $(i)_+$ denoting position before and after the i th point. Generalised force and displacement quantities of a sub-point is assembled in a column matrix called state vector $\{Z\}$. The part of the structure between $(i)_-$ and $(i-1)_+$ is defined as i th field and that between $(i)_-$ and $(i)_+$ is defined as i th point. Field transfer matrix $[F_i]$ relates the state vectors of two subpoints in i th field and is obtained from



TEST STRUCTURE

FIG. 2.1.1

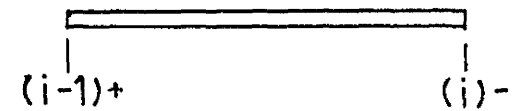


$\{Z_i^+\}$ State Vector at Point i & Subpoint +

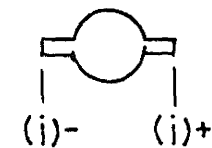
$[F_i]$ Field Transfer Matrix at i th Field

$[P_i]$ Point Transfer Matrix at i th Point

$[L_i]$ Load Vector at i th Point



i th field



i th point

FIG. 2.2.1

Simplification of this matrix equation will yield a set of linear simultaneous equation which can be solved. The solution will give the values of boundary state vectors $\{z_n^+\}$ and $\{z_0^+\}$. The state vector at other points follow from the eq. 2.2.1 and 2.2.2 as,

$$\{z_1^-\} = [F_1]\{z_0^+\}$$

$$\{z_1^+\} = [P_1][F_1]\{z_0^+\} + \{L_1\}$$

$$\{z_2^-\} = [F_2][P_1][F_1]\{z_0^+\} + [F_2]\{L_1\}$$

$$\{z_2^+\} = [P_2][F_2][P_1][F_1]\{z_0^+\} + [P_2][F_2]\{L_1\} + \{L_2\}$$

$$\{z_n^-\} = [F_n][P_{n-1}][F_{n-1}] \dots [P_1][F_1]\{z_0^+\}$$

$$+ [F_n][P_{n-1}][F_{n-1}] \dots [P_2][F_2]\{L_1\}$$

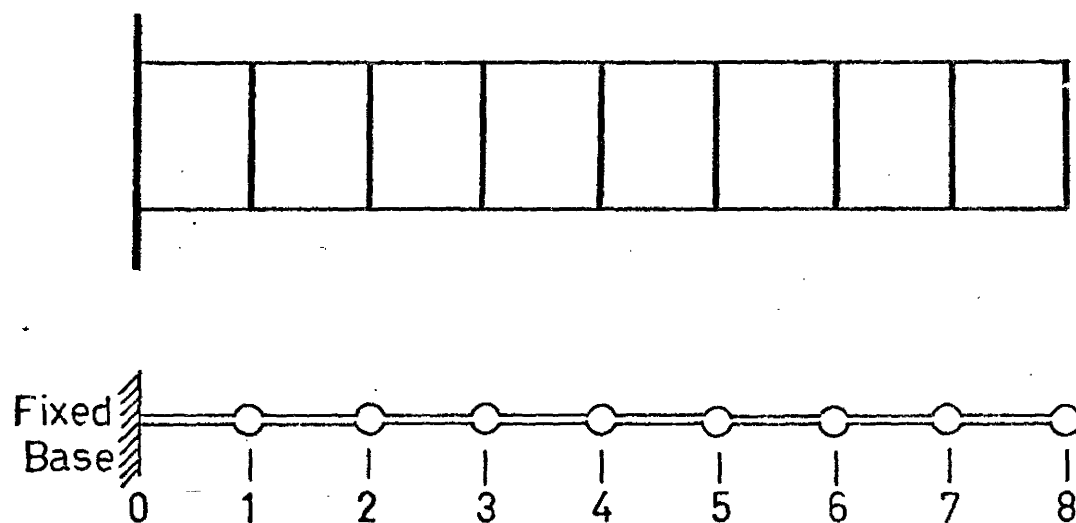
$$+ [F_n][P_{n-1}][F_{n-1}] \dots [P_3][F_3]\{L_2\}$$

$$\dots \dots \dots + [F_n]\{L_{n-1}\}$$

(2.2.4)

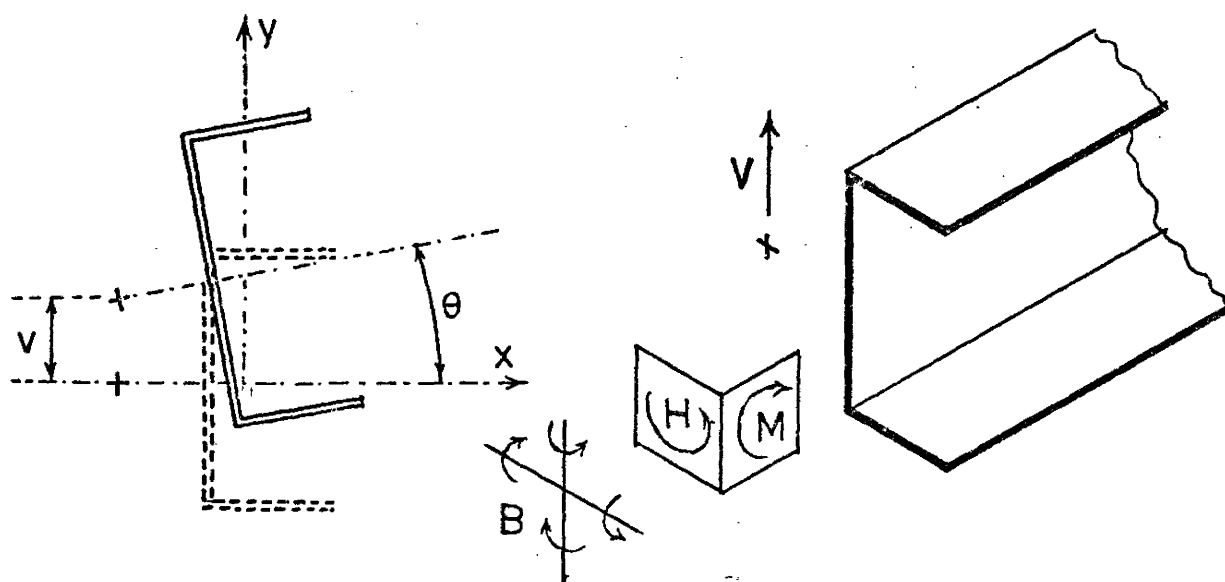
2.3 Application of Matrix Transfer Method

The structure considered consists of a prismatic mono symmetric section with equally spaced slabs. The floor slab represents 'point' and the part of the beam in between floor slabs represents 'field' as defined earlier. The simplified model to be used in Matrix Transfer method is shown in Fig. 2.3.1. In this case n equals eight and all field transfer matrices and points transfer matrices are identical.



SIMPLIFIED MODEL

FIG. 2.3.1



FORCE & DISPLACEMENT QUANTITIES

FIG. 2.3.2

2.3.1 State Vector

The state vector is an eighth order column matrix consisting of the following terms

$$Z = \begin{Bmatrix} v \\ v' \\ M \\ V \\ \theta \\ \theta' \\ B \\ H \end{Bmatrix}$$

The notations are explained and illustrated in fig. 2.3.2.

2.3.2 Field Transfer Matrix

The field is a prismatic thin walled beam of length l and its transfer matrix is obtained from the solution of the differential equations. When referred to principal axes, the uncoupled differential equations for bending in y-direction and rotation are (Eq. A.4, Appendix-A):

$$EI_x v^{IV} = 0 \quad (2.3.1)$$

$$EI_\omega \theta^{IV} - GJ\theta'' = 0 \quad (2.3.2)$$

The solution of the first equation yields the following expressions for displacement, slope, moment and shear.

$$\left. \begin{aligned} v(Z) &= D_1 Z^3/6 + D_2 Z^2/2 + D_3 Z + D_4 \\ v'(Z) &= D_1 Z^2/2 + D_2 Z + D_3 \\ M(Z)/EI_x &= v'' = D_1 Z + D_2 \\ V(Z)/EI_x &= -v''' = -D_1 \end{aligned} \right\} \quad (2.3.3)$$

Where D_1, D_2, D_3 and D_4 are constants of integration determined from the boundary condition at $Z=0$ namely

$$v = v(0), v' = v'(0), M = M(0) \text{ and } V = V(0)$$

The state vector at $Z = \ell$ can be expressed in terms of state vector at $Z = 0$ by the relation,

$$\left. \begin{aligned} v(\ell) &= (0) + \ell v'(0) + \frac{\ell^2}{2EI_x} M(0) - \frac{\ell^3}{6EI_x} V(0) \\ v'(\ell) &= v'(0) + \frac{\ell}{EI_x} M(0) - \frac{\ell^2}{2EI_x} V(0) \\ M(\ell) &= M(0) + \ell V(0) \\ V(\ell) &= V(0) \end{aligned} \right\} \quad (2.3.4)$$

The solution of the second equation yields the following expressions for rotation, warping, bimoment and torque.

$$\left. \begin{aligned} \theta(Z) &= C_1 + C_2 Z + C_3 \sinh KZ + C_4 \cosh KZ \\ \theta'(Z) &= C_2 + C_3 K \cosh KZ + C_4 K \sinh KZ \\ B(Z)/EI_\omega &= -\theta'' = -C_3 K^2 \sinh KZ - C_4 K^2 \cosh KZ \\ H(Z) &= -EI_\omega \theta''' + CJ\theta' = C_2 \cdot GJ \end{aligned} \right\} \quad (2.3.5)$$

$$\text{Where } K = \sqrt{\frac{GJ}{EI_\omega}}$$

The constants C_1, C_2, C_3 and C_4 are constants of integration determined from the boundary condition at $Z = 0$ namely

$$\theta = \theta(0), \theta' = \theta'(0), B = B(0) \text{ and } H = H(0)$$

The boundary condition at $Z = \ell$ are

$$\theta = \theta(\ell), \theta' = \theta'(\ell), B = B(\ell) \text{ and } H = H(\ell)$$

They can be expressed in terms of $\theta(0)$ $\theta'(0)$, $B(0)$ and $H(0)$ by the following expressions

$$\left. \begin{aligned}
 \theta(l) &= \theta(0) - \frac{1}{K} \sinh Kl \theta'(0) - \frac{1}{GJ} (1 - \cosh Kl) B(0) \\
 &\quad + \frac{1}{GJ} \left(l - \frac{1}{K} \sinh Kl \right) H(0) \\
 \theta'(l) &= \cosh Kl \theta'(0) - \frac{K}{GJ} \sinh Kl B(0) \\
 &\quad + \frac{1}{GJ} (1 - \cosh Kl) H(0) \\
 B(l) &= - \frac{GJ}{K} \sin Kl \theta'(0) + \cosh Kl B(0) \\
 &\quad + \frac{1}{K} \sinh Kl H(0) \\
 H(l) &= H(0)
 \end{aligned} \right\} (2.3.6)$$

Eq. 2.3.4 and eq. 2.3.6 can be combined in a single matrix eq. as

$$\{Z(l)\} = [F]\{Z(0)\} \quad (2.3.7)$$

where

$$\{Z(l)\} = \begin{Bmatrix} v(l) \\ v'(l) \\ M(l) \\ V(l) \\ \theta(l) \\ \theta'(l) \\ B(l) \\ H(l) \end{Bmatrix} \quad \text{and} \quad Z(0) = \begin{Bmatrix} v(0) \\ v'(0) \\ M(0) \\ V(0) \\ \theta(0) \\ \theta'(0) \\ B(0) \\ H(0) \end{Bmatrix}$$

Field transfer matrix is an eighth order square matrix

$$[F] = \begin{bmatrix} f_{11} & f_{12} & f_{13} & f_{14} & 0 & 0 & 0 & 0 \\ 0 & f_{22} & f_{23} & f_{24} & 0 & 0 & 0 & 0 \\ 0 & 0 & f_{33} & f_{34} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & f_{44} & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & f_{55} & f_{56} & f_{57} & f_{58} \\ 0 & 0 & 0 & 0 & 0 & f_{66} & f_{67} & f_{68} \\ 0 & 0 & 0 & 0 & 0 & f_{76} & f_{77} & f_{78} \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & f_{88} \end{bmatrix} \quad (2.3.8)$$

Where the non zero elements are

$$f_{11} = 1, f_{12} = \ell, f_{13} = \ell^2/2EI_x, f_{14} = -\ell^3/6EI_x,$$

$$f_{22} = 1, f_{23} = \ell/EI_x, f_{24} = -\ell^2/2EI_x,$$

$$f_{33} = 1, f_{34} = \ell, f_{44} = 1,$$

$$f_{55} = 1, f_{56} = \sinh K\ell/K,$$

$$f_{57} = (1 - \cosh K\ell)/GJ, f_{58} = (\ell - \frac{1}{K} \sinh K\ell)/GJ$$

$$f_{66} = \cosh K\ell, f_{67} = -K \sinh K\ell/GJ$$

$$f_{68} = (1 - \cosh K\ell)/GJ, f_{76} = -GJ \sin K\ell/K$$

$$f_{77} = \cosh K\ell, f_{78} = \sinh K\ell/K, f_{88} = 1.$$

2.3.3 Action of Floor Slab

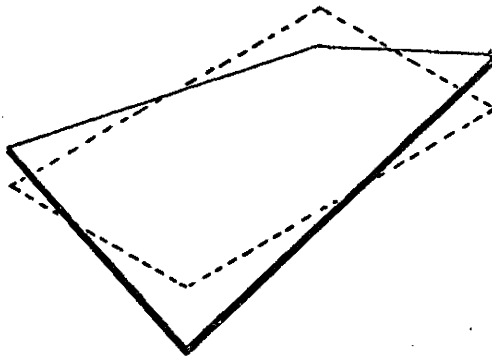
Vlasov (1) considered the effect of a diaphragm on the behaviour of a thin walled beam. In a shear wall structure, floor slab is equivalent to diaphragm. Vlasov assumed that the diaphragm acts as a plate in torsion and derived the following relationship for the bimoment applied to the shear wall by the action of the slab (Fig. 2.3.4(a)).

$$B_t = \frac{Et^3bd}{6(1+\nu)} \theta' \quad (2.3.9)$$

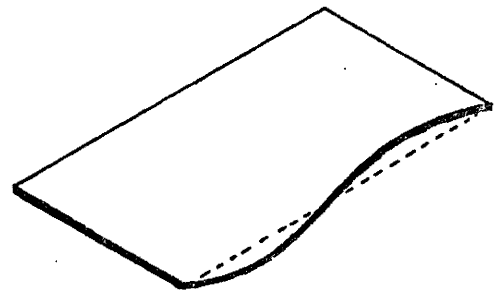
Where

- b = width of slab
- d = length of slab
- t = thickness of slab
- ν = Poisson's ratio
- E = modulus of elasticity of slab
- θ' = warping of the shear wall at the level of slab.

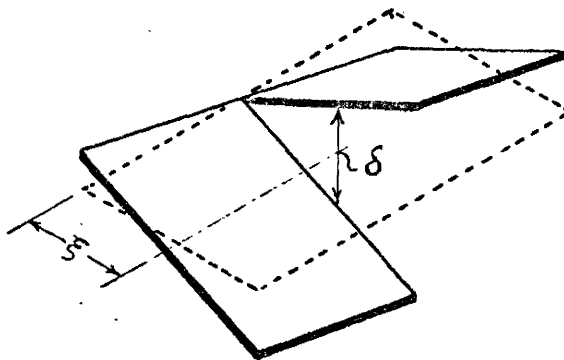
In this derivation Vlasov neglected the effect of bending of slab (Fig. 2.3.4 (b)) due to fixity of walls. This can be considered by treating the slab as a series of beams running between the flanges. The center line of the beams is the locus of point of contraflexure for each beam. If a cut is made along that line, there will be relative displacement to the left and right of the cut. (Fig. 2.3.4(c)). Shear force q will develop along the line to maintain continuity (Fig. 2.3.4 (d)). For an element of beam at a distance ξ from the wall, the sectorial areas at the center



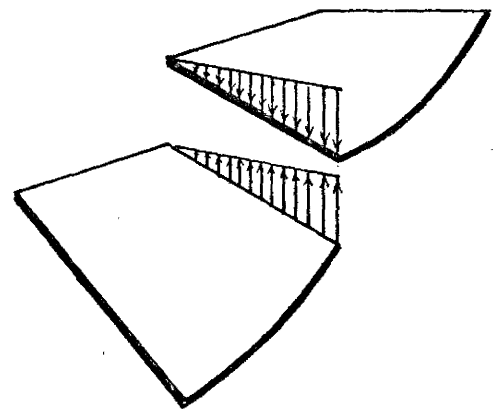
(a) Torsion of slab



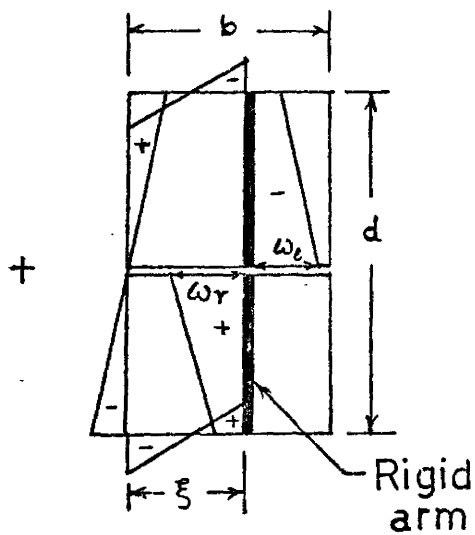
(b) Bending of slab



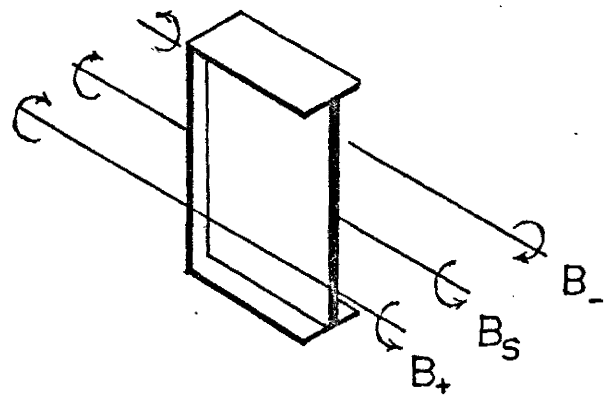
(c) Slab cut at center



(d) Shear force



(e) Sectorial area



(f) Bimoment from slab

(Fig. 2.3.4(e)) considering rigid arms attached to thin walled beam are:

$$\omega_r = \xi d, \omega_l = -\xi d.$$

Discrepancy of displacement is

$$\delta = (\omega_r - \omega_l) \theta' = 2\xi d \theta'$$

Shear force develops to maintain continuity considering bending deformation only

$$q = \frac{2Et^3\xi}{(1-\nu^2)d^2} \theta'$$

The bimoment due to the shear force is

$$dB = q(\omega_r - \omega_l) d\xi = 2 q \xi d d\xi$$

The total bimoment due to bending is obtained on integration

$$B_b = \int_0^b dB = \frac{4Et^3b^3}{3d(1-\nu^2)} \theta', \quad (2.3.10)$$

The combined bimoment due to torsion and bending is

$$B_s = B_t + B_b = D \theta' \quad (2.3.11)$$

Where

$$D = \left(\frac{Et^3bd}{6(1+\nu)} + \frac{4Et^3b^3}{3d(1-\nu^2)} \right) \quad (2.3.12)$$

From Fig. 2.3.4 (a), the bimoment contribution from slab is related to the bimoments in the walls immediately above and below the floor slab.

$$B_+ = B_- - B_s = B_- - D \theta' \quad (2.3.13)$$

2.3.4. Point Transfer Matrix and Load Vector

The Point transfer matrix [P] is a square matrix of order eighth. It is obtained from the consideration of

equilibrium and compatibility:

$$[P] = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -D & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix}$$

The load vector is a column matrix of order eight.

$$\{L_i\} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ -P_i \\ 0 \\ 0 \\ 0 \\ -Q_i \end{pmatrix}$$

Where $i = 1, 2, \dots, 8$

and P_i is applied load

and Q_i is the applied
torque at i th level.

2.3.5 Boundary Conditions

The shear wall is fixed at the base and free at top.
Therefore, the state vectors at the base and the top
can be written as

Here 'a' denotes the elements of matrix [A] as defined in eq. 2.2.3 (a). The solution is obtained by inversion

$$\{Y\} = [B]^{-1}\{N\} \quad (2.3.15)$$

Knowing $\{Y\}$, $\{Z_0^+\}$ and $\{Z_8^+\}$ can be formed. State vector at other points are obtained from expressions in eq. 2.2.4.

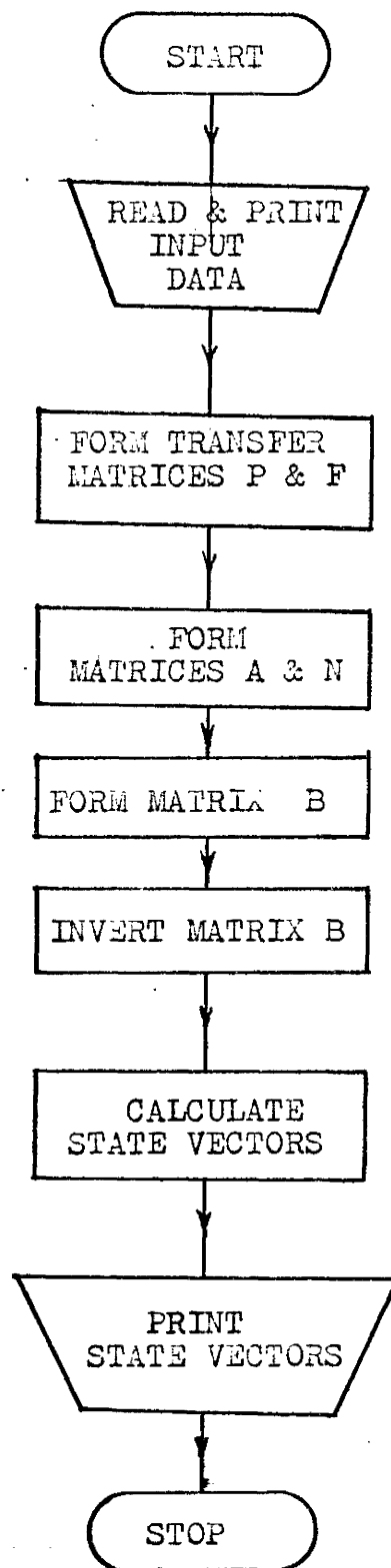
2.4 Computer Program

A computer program based on the above analysis has been written. The input data are the geometric and elastic properties and loading of the structure. The output quantities are the state vectors at all floor levels. The present program is for identical floor slabs, equal storey heights and prismatic section. Extension for stepped cases or different storey heights and floor slabs can be made with little modification.

The flow chart is given in Fig. 2.4.1 and the computer program is included in Appendix-B.

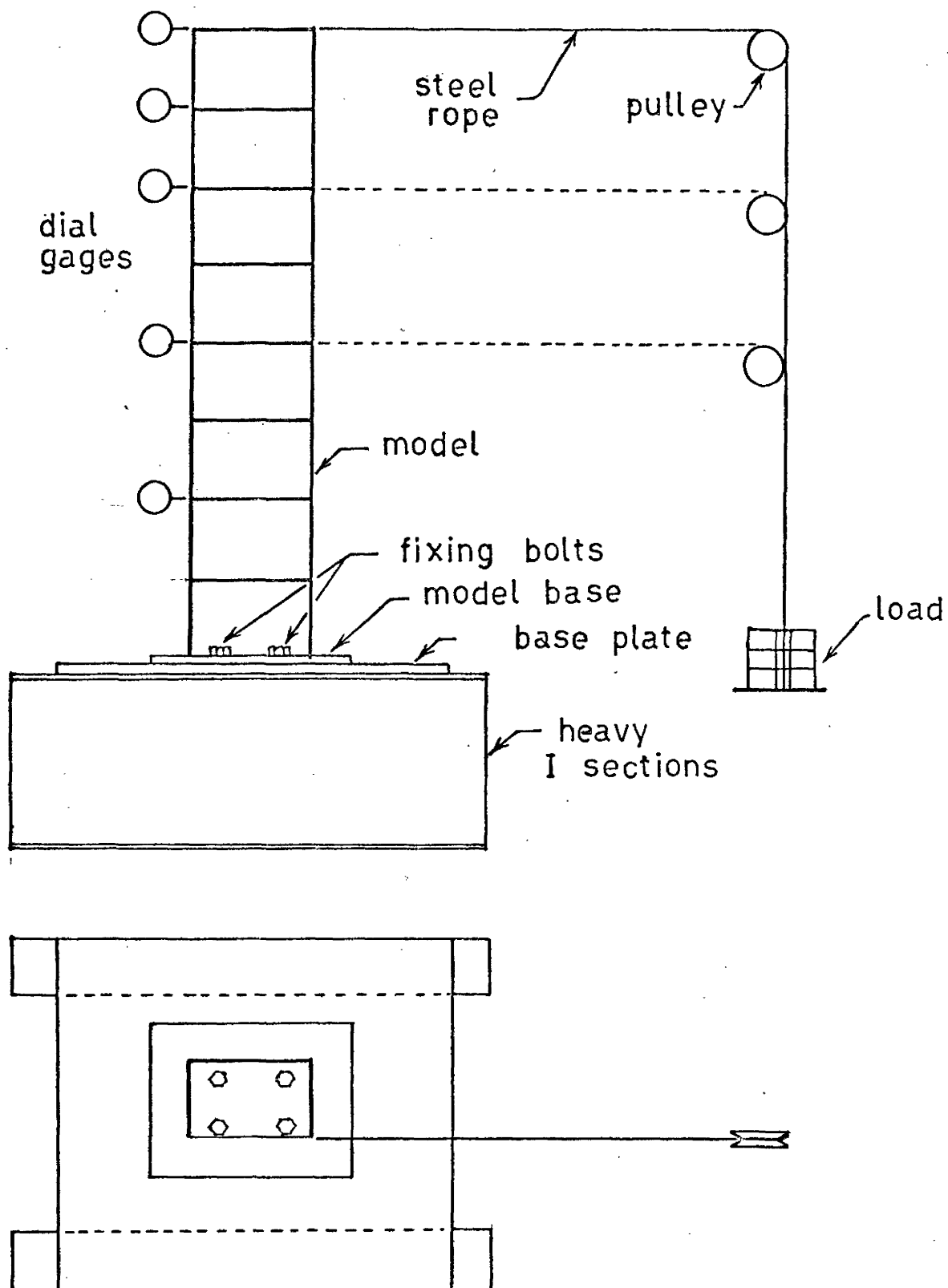
2.5 Experiment

An experiment was done on a small scale plexiglas model (Fig. 2.1.1). It was made by assembling different components representing walls and floors. The base of the model was connected to a thick base plate which in turn was fixed to two heavy I sections to achieve fixity (Fig. 2.5.1). It was loaded at the 8th, 6th and 4th floor respectively, one floor at a time, by hanging weights over



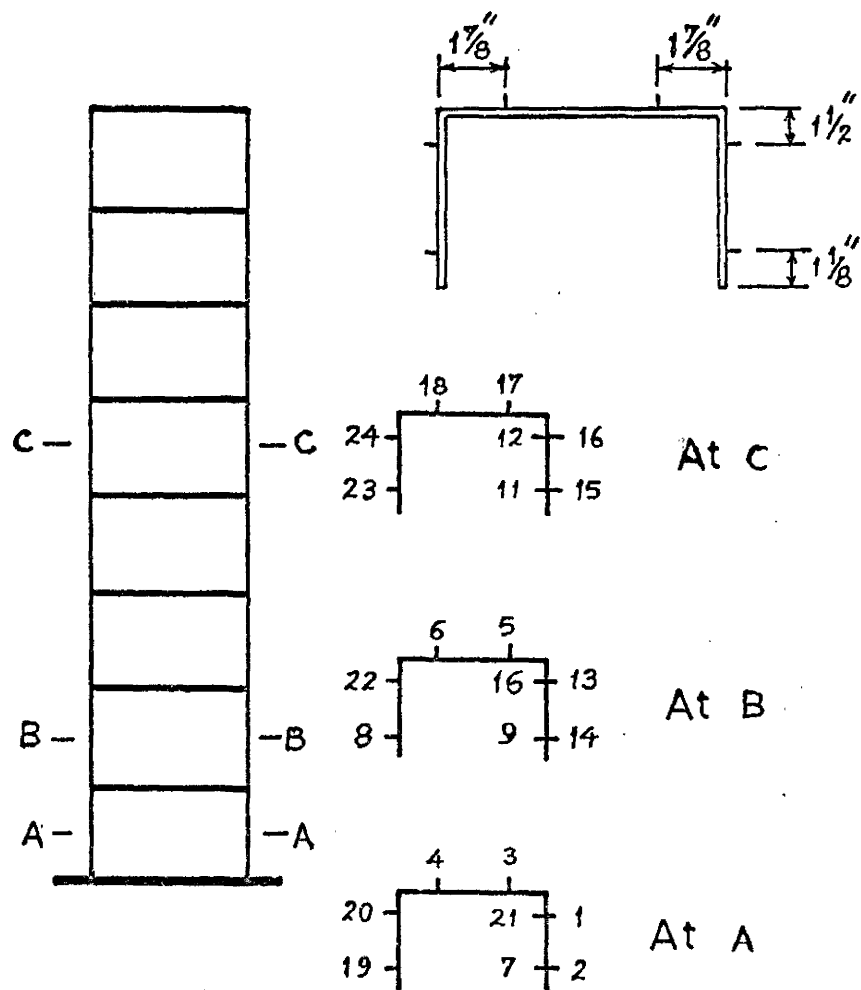
FLOW CHART

FIG. 2.4.1

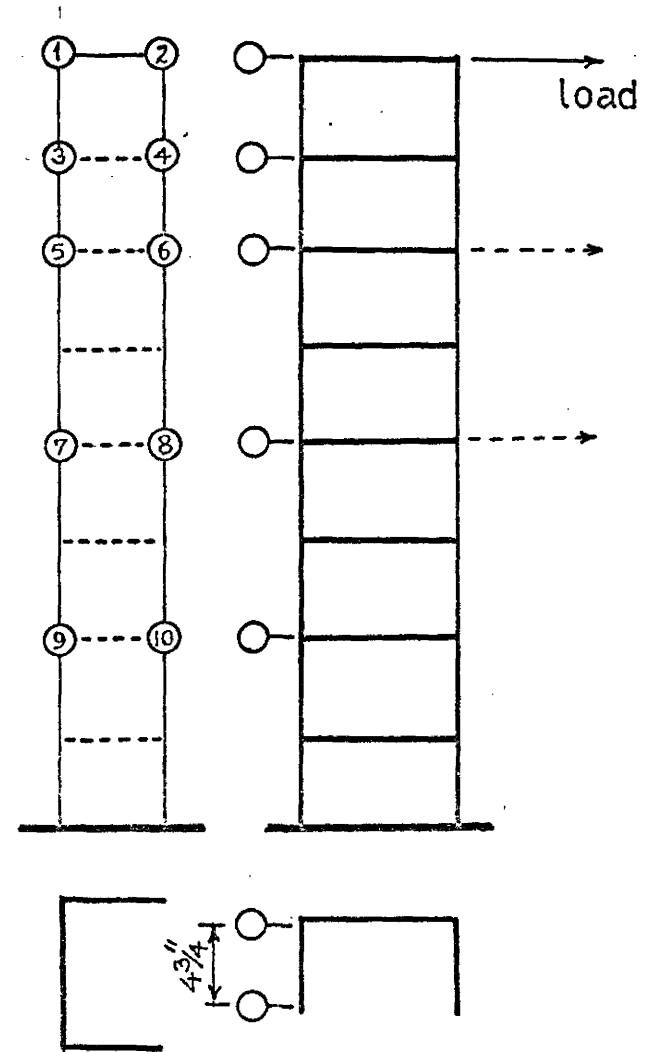


EXPERIMENTAL SET UP

FIG. 2.5.1



(a) Strain Gauge Positions



(b) Dial Gauge Positions

FIG. 2.5.2

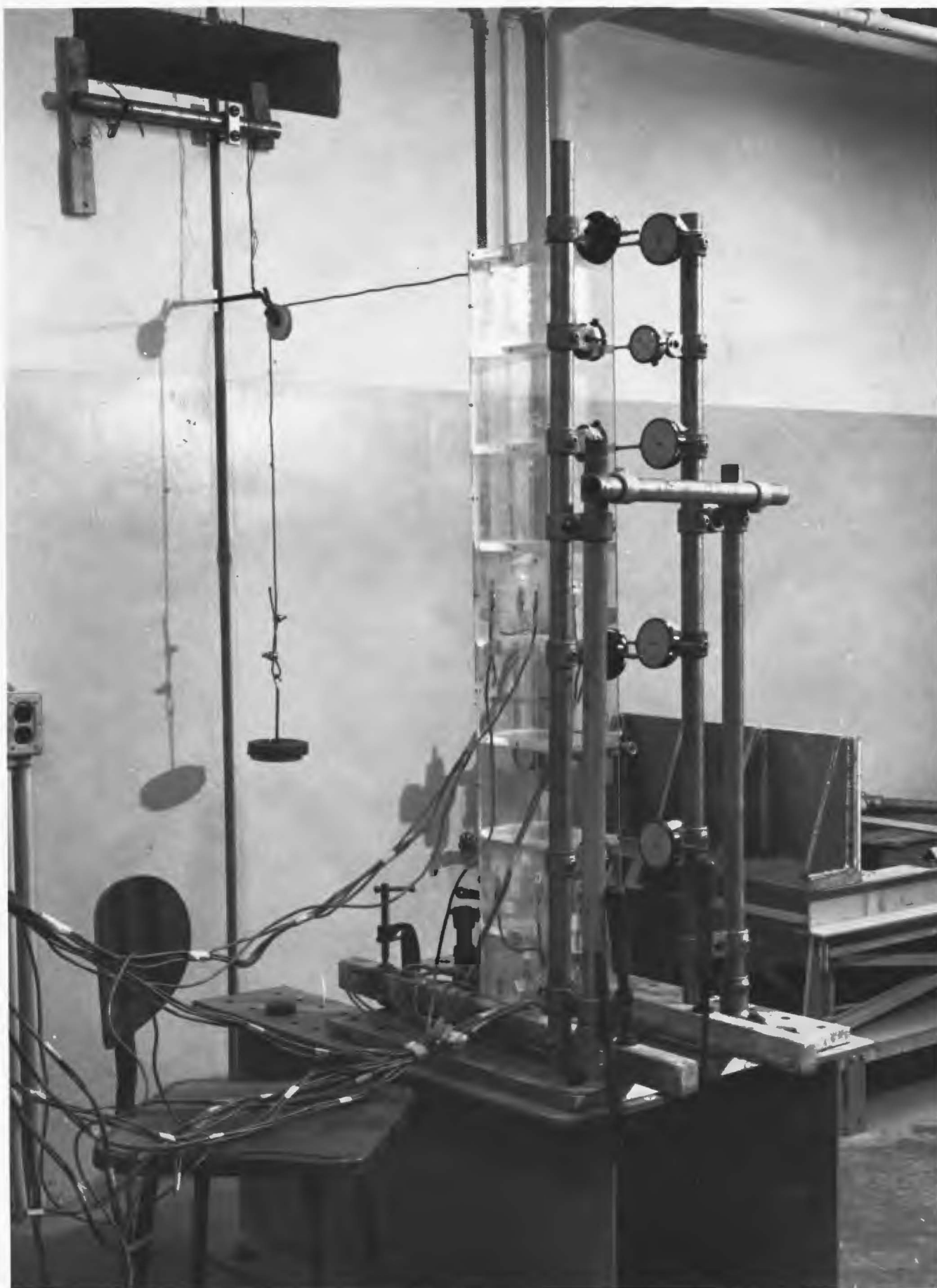


Fig. 2.5.3
EXPERIMENTAL SET UP FOR STATIC TEST

a pulley. Strain gauges were attached at the middle of 1st, 2nd and 5th storey. Leads from the strain gauges were hooked up to a strain indicator through switch boxes and strain readings at every increment of loading were taken. Deflections are measured from readings of dial gauges mounted at different points of the structure (Fig. 2.5.2). The strain gauge and dial gauge readings are tabulated in Appendix C. Fig. 2.5.3 shows the experimental set up for the case with loading at top.

The following is a list of equipment and materials used in this experimental work.

- A. Model Material: Plexiglas
Elastic properties: $E = 0.40 \times 10^6$ psi, $\nu = 0.35$
- B. Electric Resistance Strain Gauges
Make: Micro Measurement
Type: EA-41-25086-120
suitable for plastic
Resistance: $120 \Omega \pm 0.15\%$
Gauge Factor: $2.01 \pm 0.5\%$
- C. Dial Gauges:
Make: Baty
Reading: .001 in and .0001 in
- D. Strain Indicator:
Make: Budd Corporation
Reading: Directly calibrated to strain in $\mu\text{in/in}$
Range: $\pm 40,000 \mu \text{ in/in}$

2.6 Results and Discussion

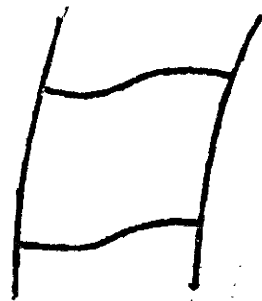
The linearity of the test structure is checked in fig.

2.6.10. The comparison of theoretical and experimental data as plotted in fig. 2.6.3 to 2.6.11 shows that the structure is not so stiff as predicted by considering both torsional and flexural stiffness of floor slabs. If only the torsional stiffness of floor slab is taken, the theoretical analysis gives a mathematical model which is more flexible than the actual structure. The difference between theory and experiment attributed to the local bending of the wall section at the joint of the floor slab as shown in fig. 2.6.1.

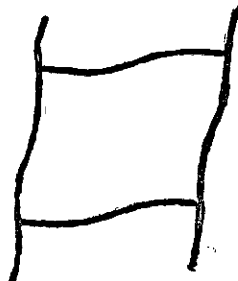
As a result of this bending, the joint is not rigid which in turn reduces the shear force q at the centerline of slab (Fig. 2.3.4d). To allow for this effect, the bimoment contribution from flexure of the floor slab is modified by a factor K . The effect of the floor is then expressed as.

$$D = \left[\frac{Et^3bd}{6(1+\nu)} + K \frac{4Et^3b^3}{3d(1-\nu^2)} \right] \quad (2.6.1)$$

An approximate method to assess the value of K is given below. Consider a one bay multistory frame as in fig. 2.6.1. Assuming points of contraflexure are at the center of storey height. Let M be the moment induced at the end of the slab strip if the joints do not rotate locally. Due to the local bending of the joints the final moment is KM where K is obtained as



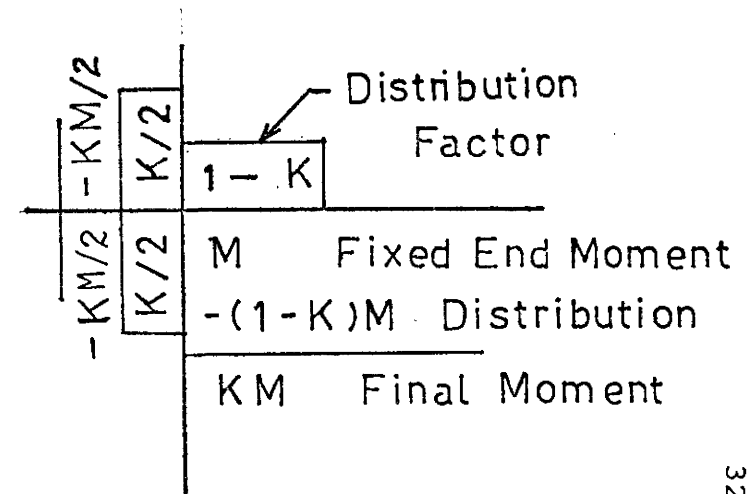
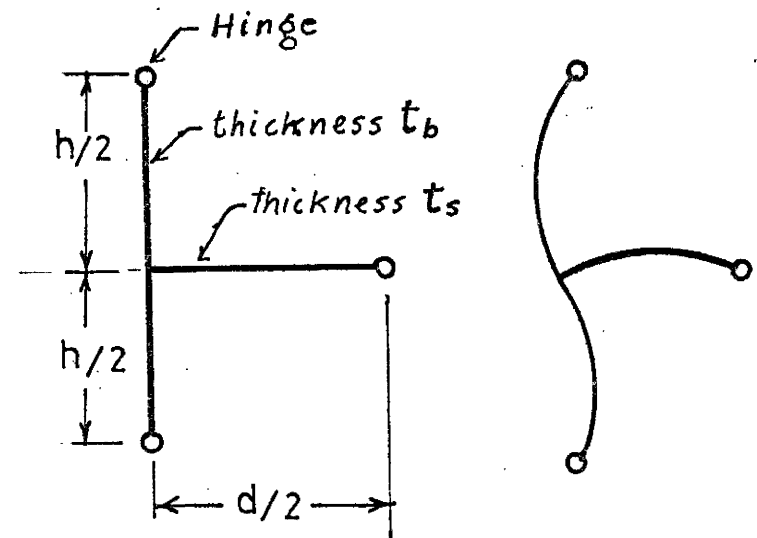
Deflected Shape in Theory



Actual Deflected Shape



FIG. 2.6.1



MOMENT DISTRIBUTION

FIG. 2.6.2

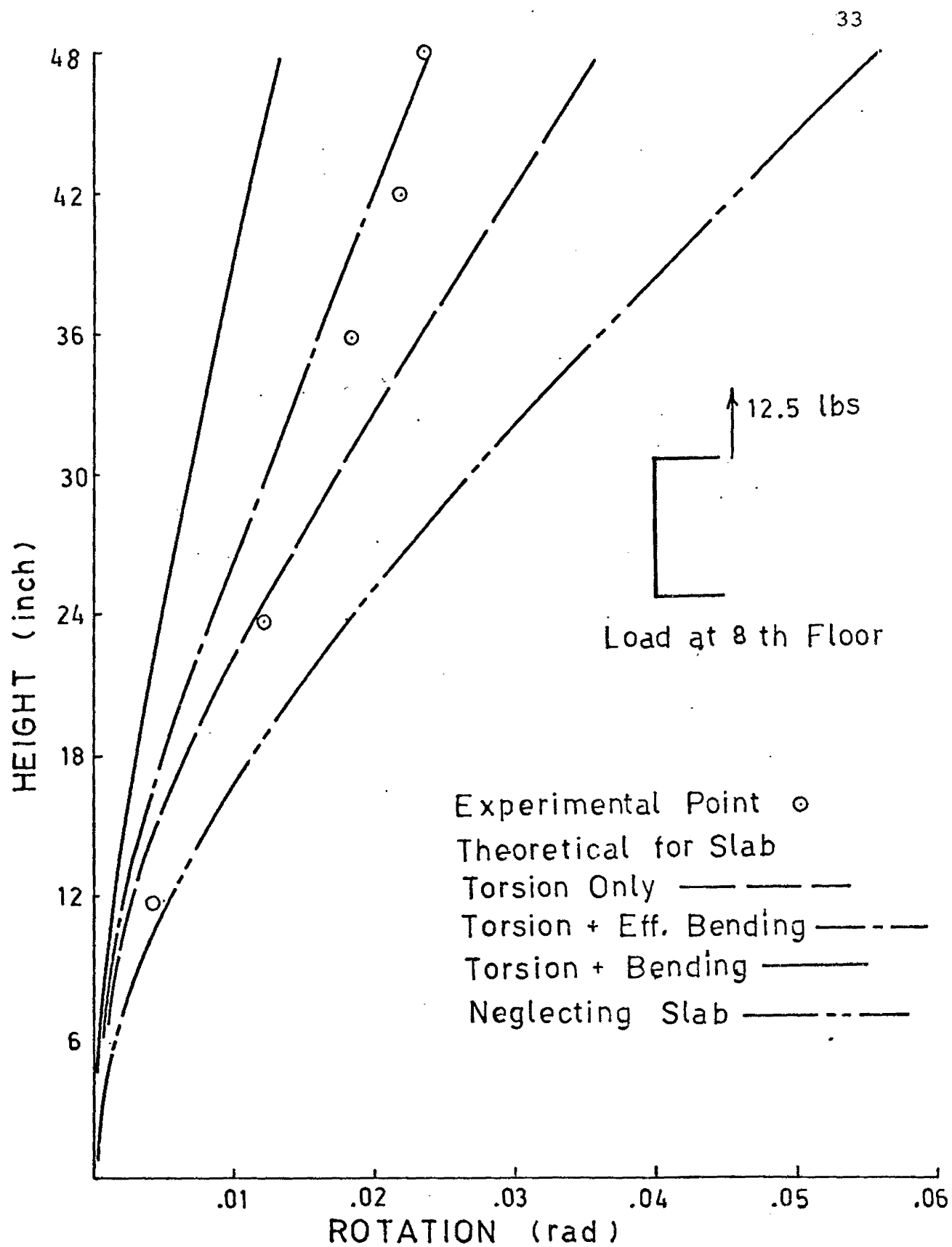


FIG. 2.6.3

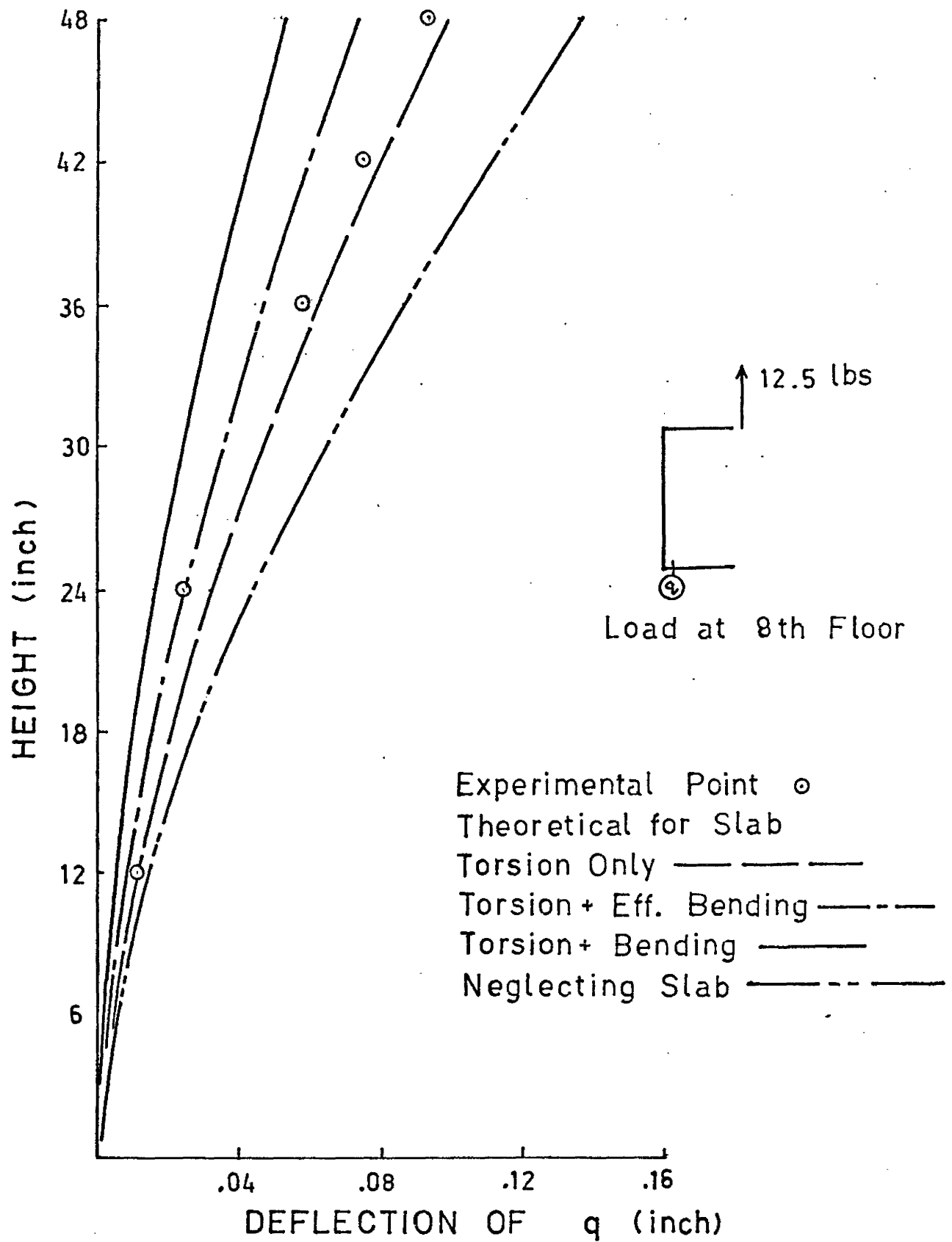


FIG. 2.6.4

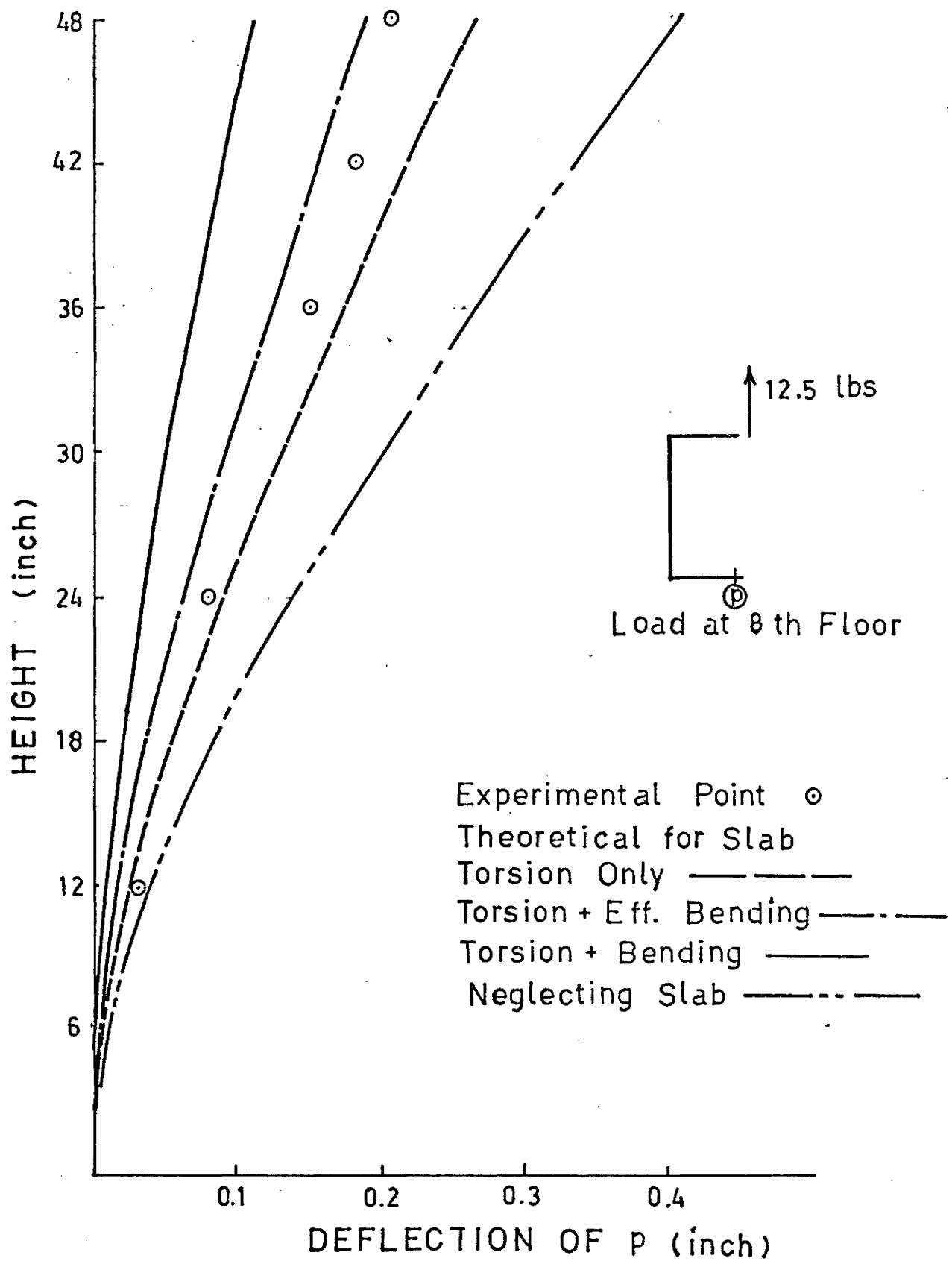


FIG. 2.6.5

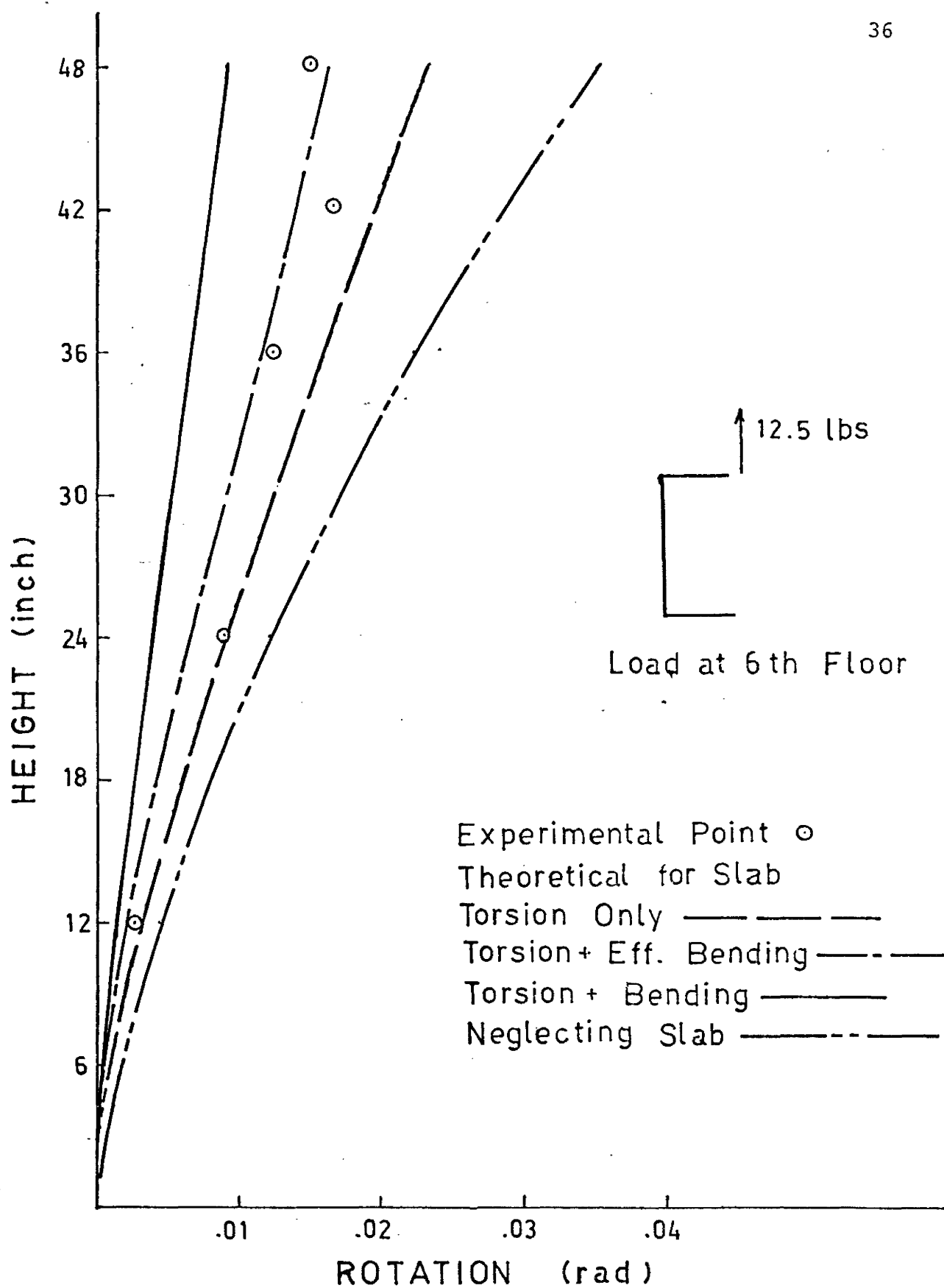


FIG. 2.6.6

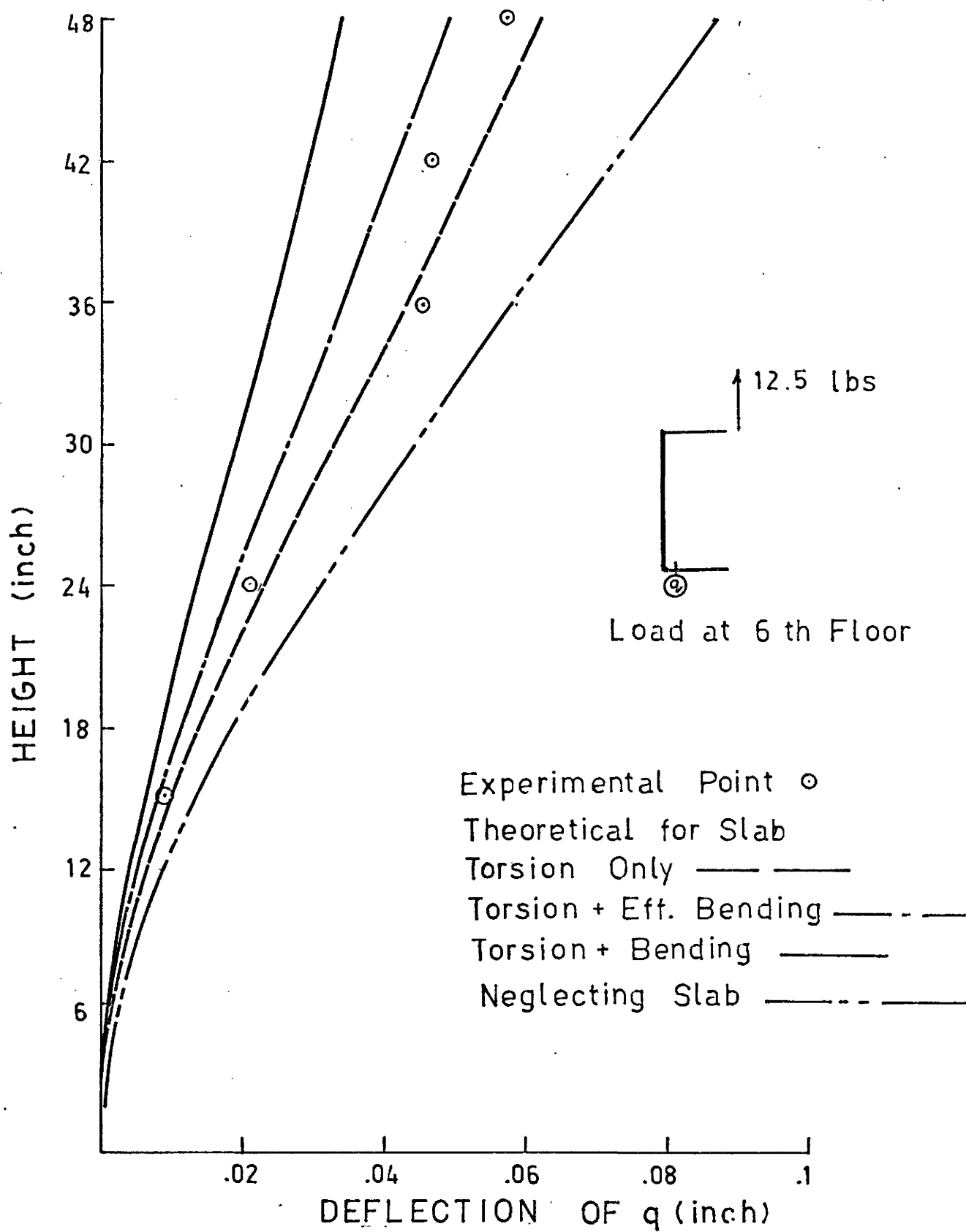


FIG. 2.6.7

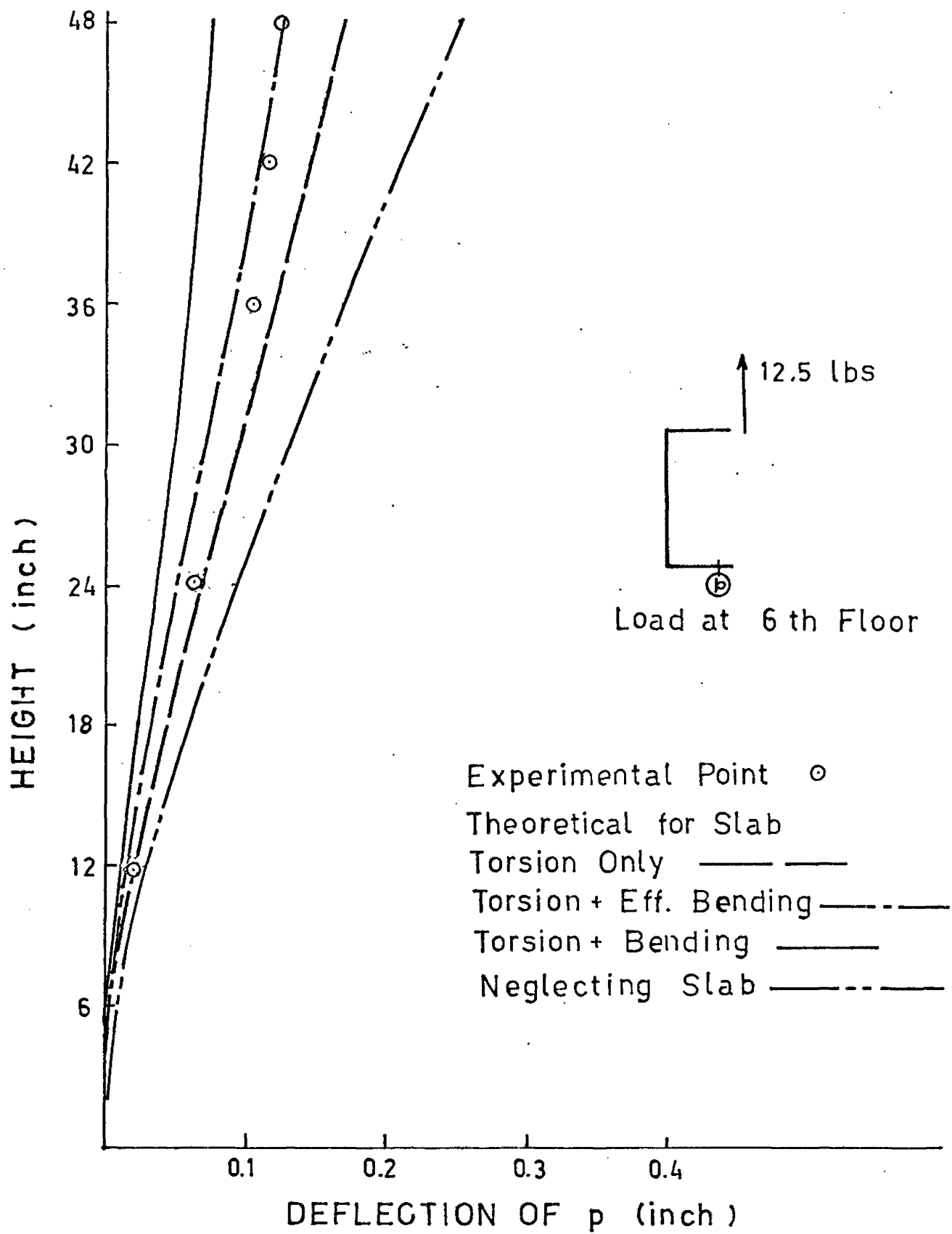


FIG. 2.6.8

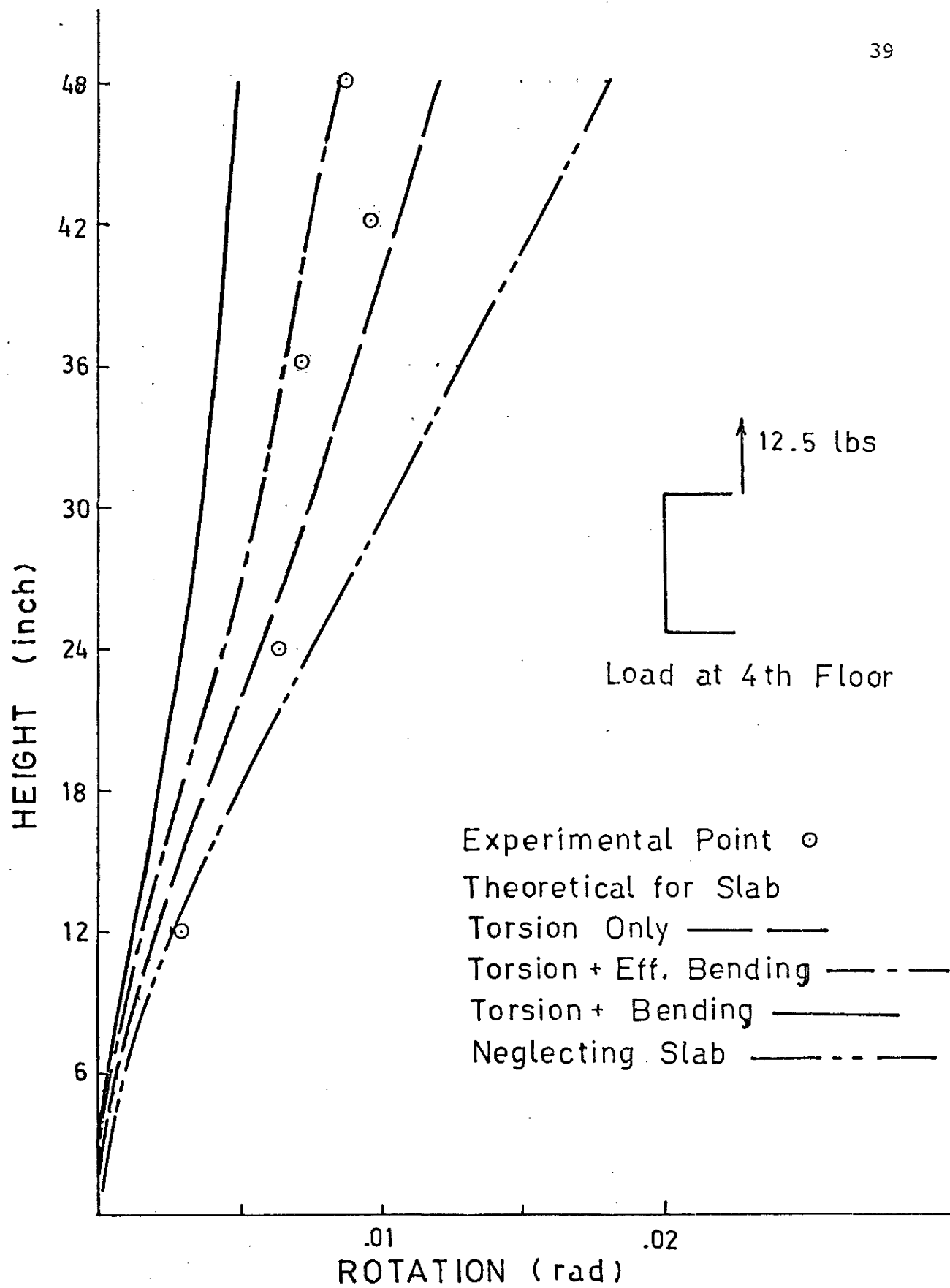


FIG. 2.6.9

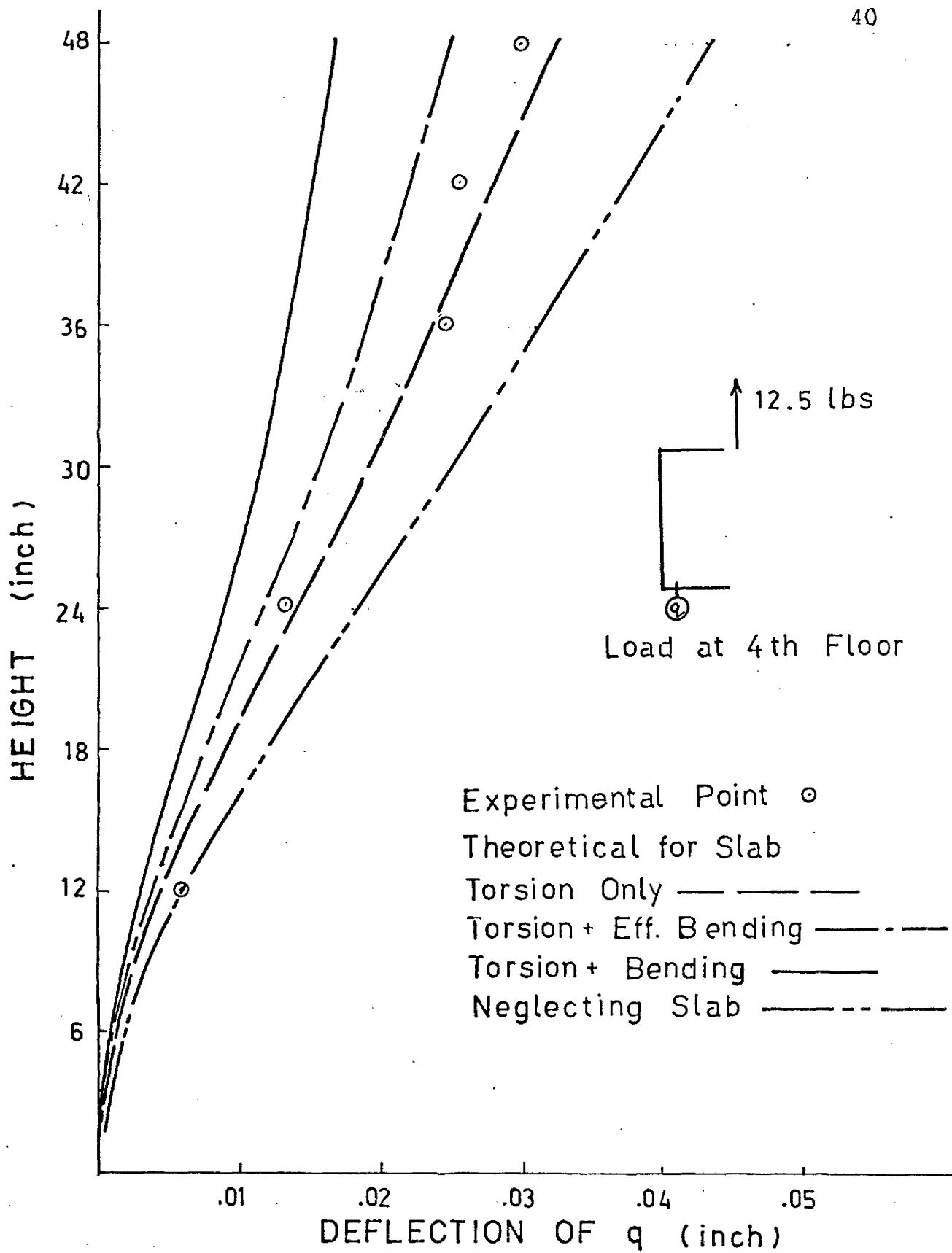


FIG. 2.6.10

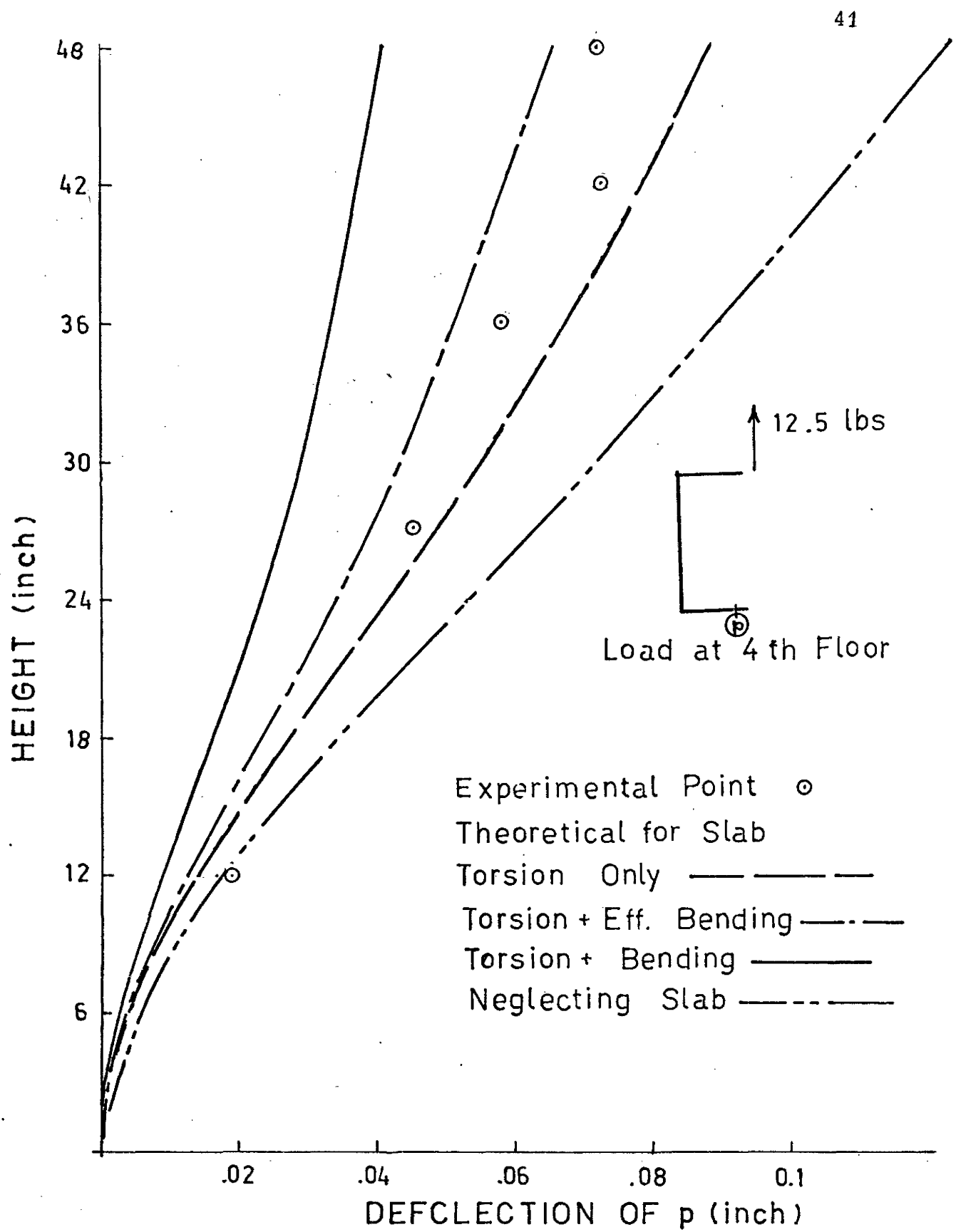
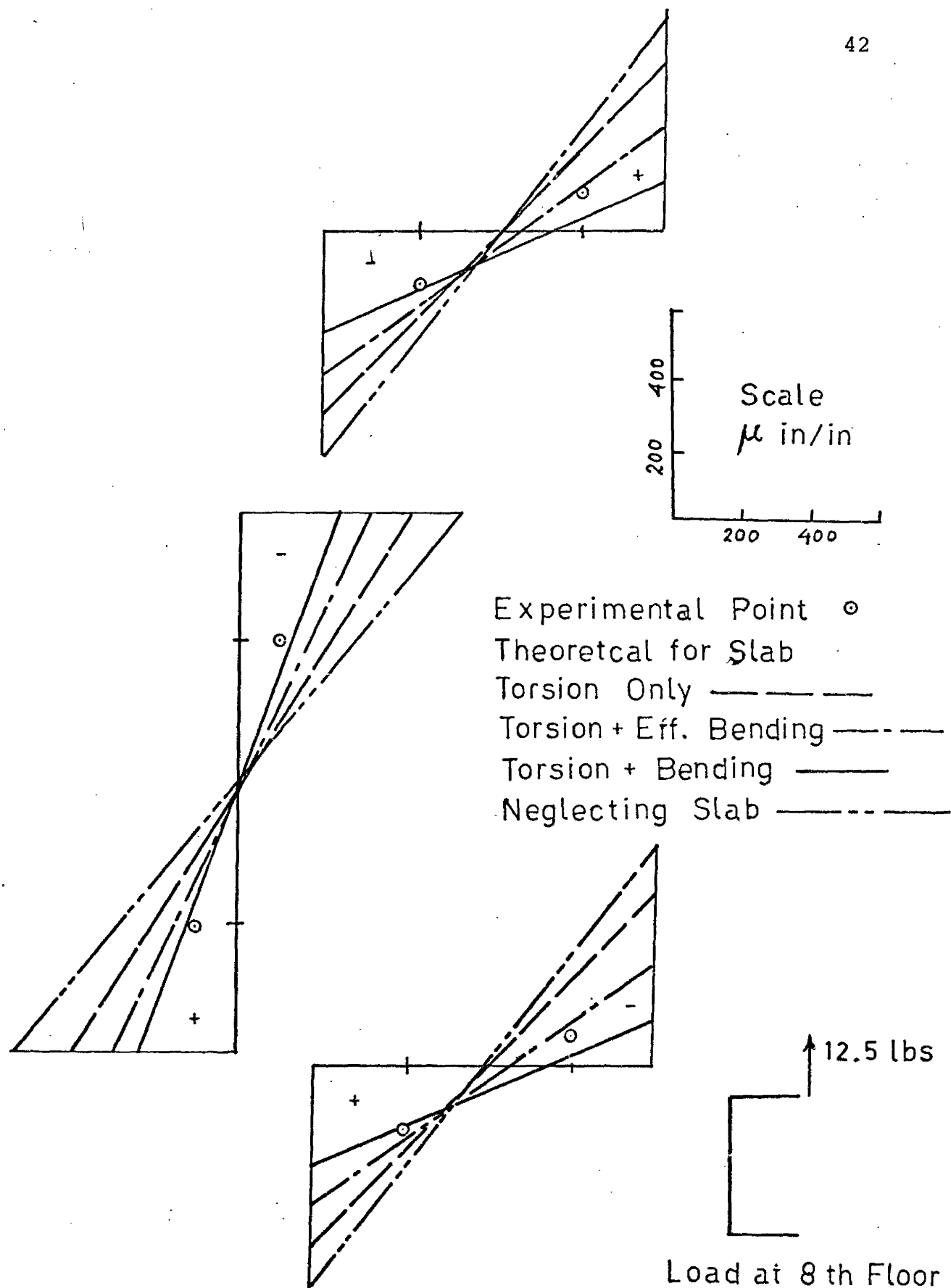
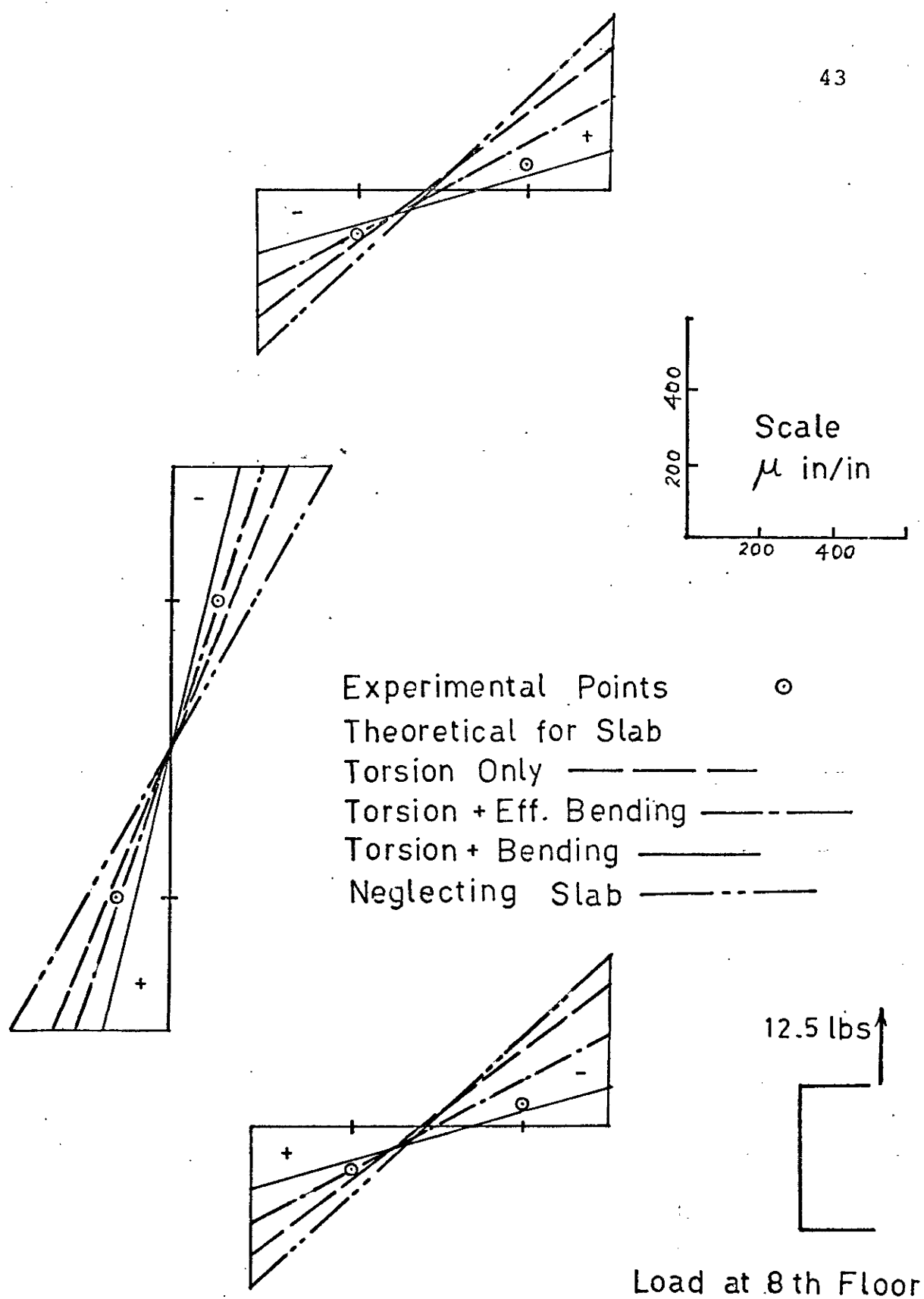


FIG. 2.6.11



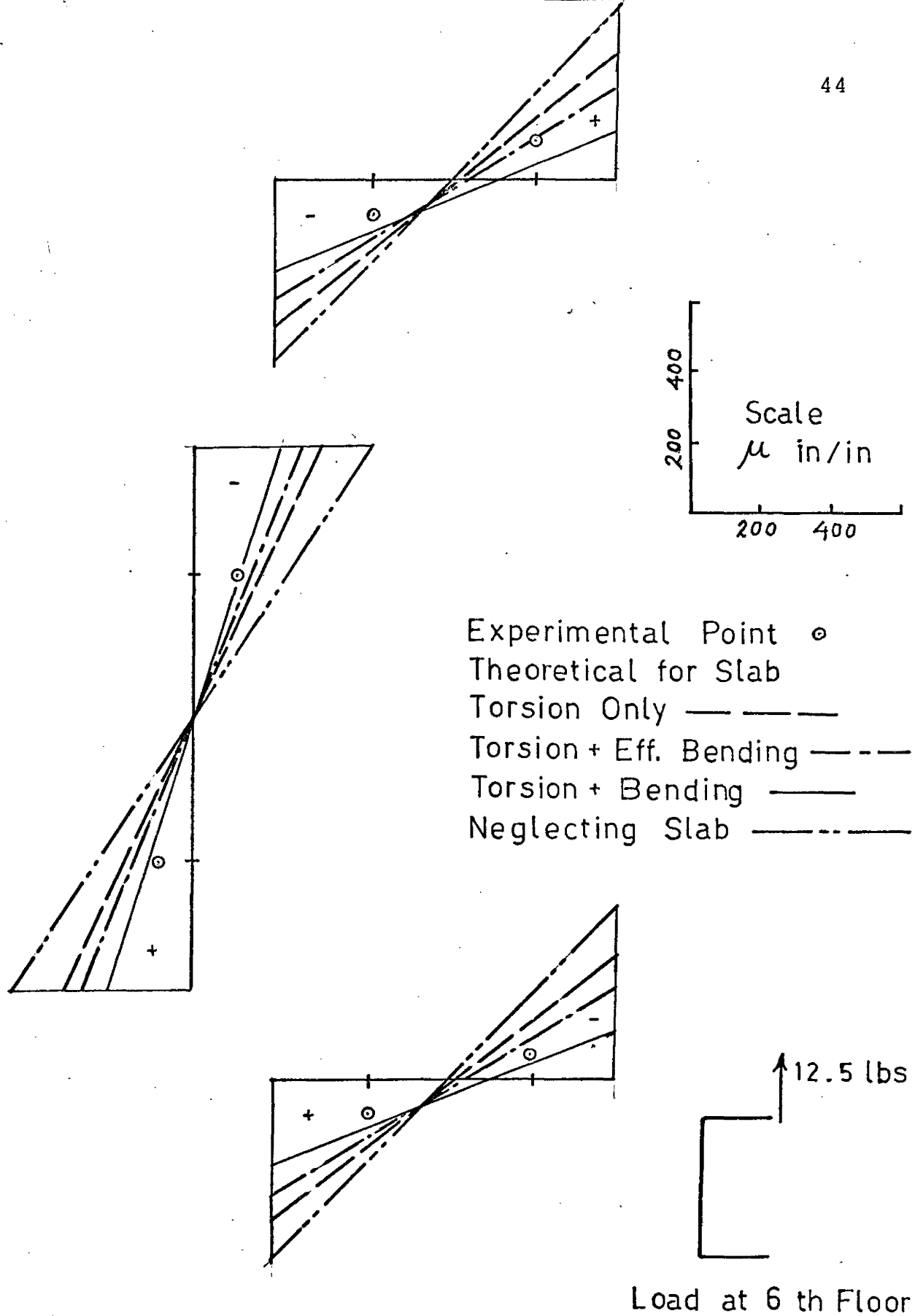
STRAIN DISTRIBUTION AT LEVEL A

FIG. 2.6.12



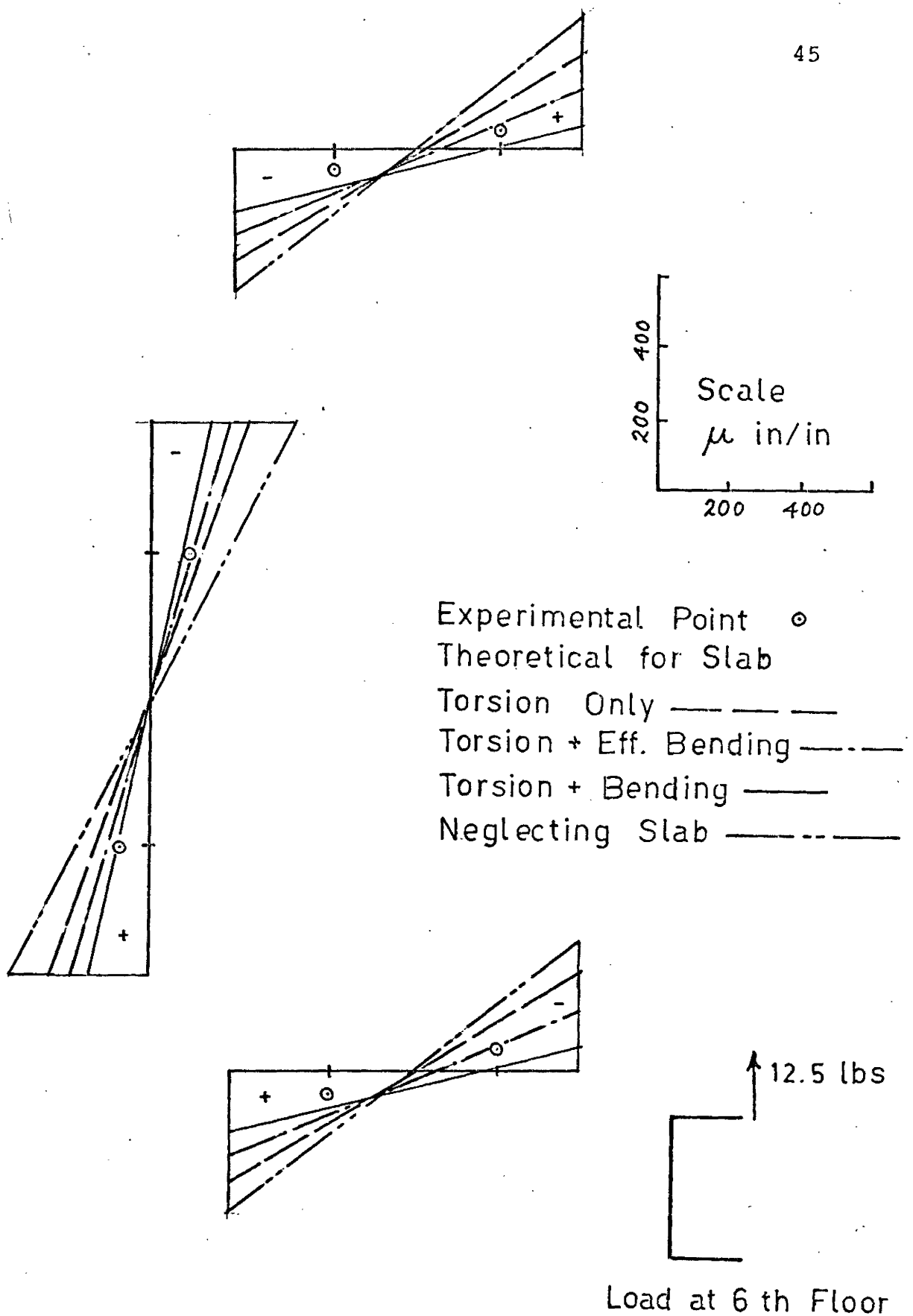
STRAIN DISTRIBUTION AT LEVEL B

FIG. 2.6.13



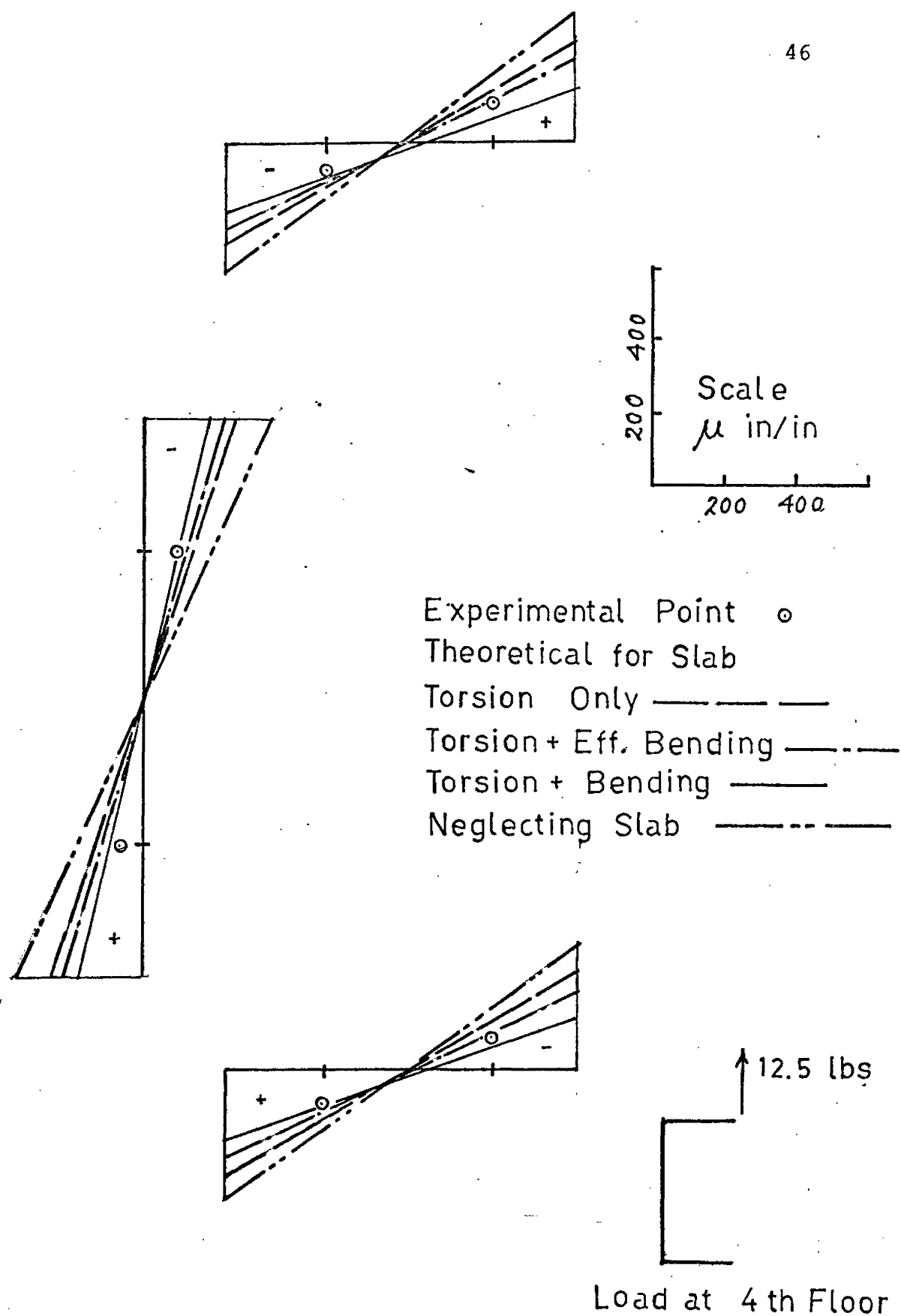
STRAIN DISTRIBUTION AT LEVEL A

FIG. 2.6.14



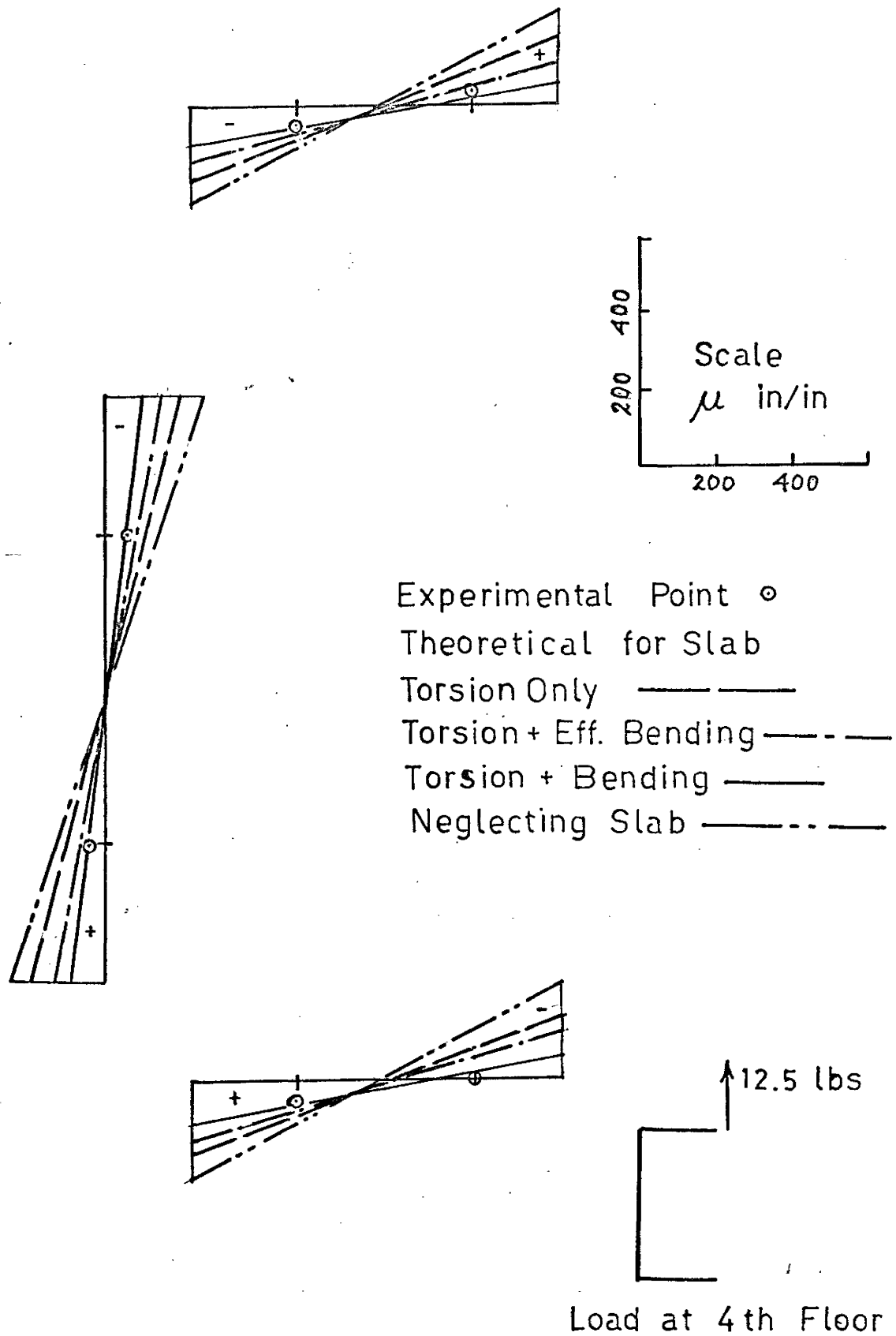
STRAIN DISTRIBUTION AT LEVEL B

FIG. 2.6.15



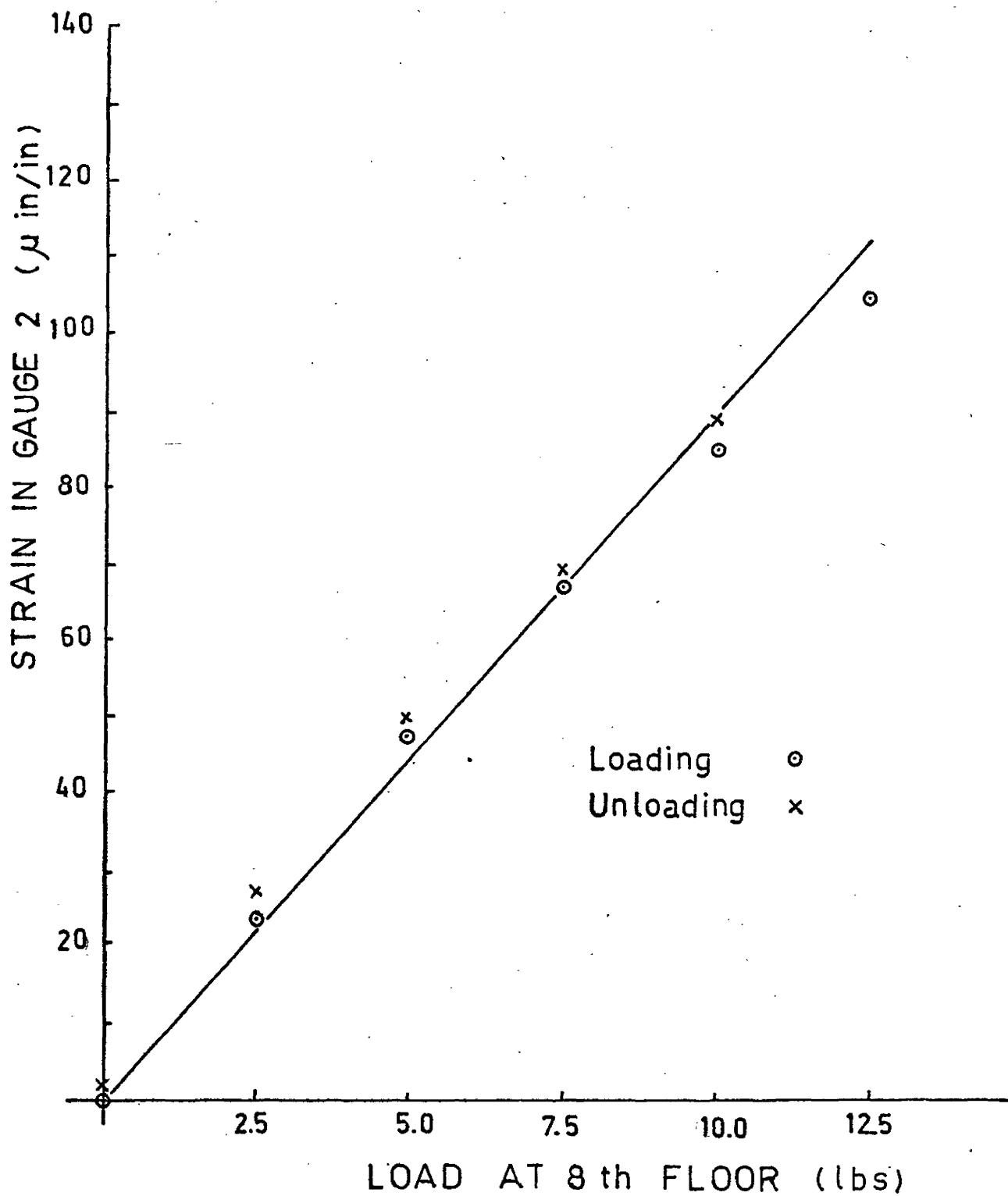
STRAIN DISTRIBUTION AT LEVEL A

FIG. 2.6.16



STRAIN DISTRIBUTION AT LEVEL B

FIG. 2.6.17



LOAD vs. STRAIN DIAGRAM FOR MODEL WITH FLOORS

FIG. 2.6.18

$$K = \frac{1}{1 + \left(\frac{t_s}{t_b}\right)^3 \frac{h}{2d}} \quad (2.6.2)$$

Where t_s and t_b are thicknesses of slab and beam flange respectively. The procedure of obtaining K is by a moment distribution scheme as shown in fig. 2.6.2.

If the ratio t_s/t_b is very small, the value of K is unity and the total flexural stiffness of the slab is effective. On the other hand if t_s/t_b is large, there will be large local bending of the flange which will violate the hypothesis of non deformable section of Vlasov. Therefore in order to use Vlasov's theory, the expression of K is valid only for $t_s/t_b < 3$.

For the model structure the value of K is 0.2545 and the displacement/rotation plots are found to have reasonable agreement with the experiment.

The strain distribution are plotted in fig. 2.6.12 to 2.6.17 and found to have reasonable agreement with the experiment.

CHAPTER III
DYNAMIC STUDY OF SHEAR WALL WITH FLOORS

3.1 Summary

In this chapter, the same shear wall structure treated in Chapter 2 is analysed, for dynamic loading. The 'Matrix Transfer' method is used in the analysis. A dynamic test was carried out to determine natural frequencies. The strain distribution at resonance was also determined. The experimental values are compared with the theoretical predictions.

3.2 Matrix Transfer Method

As in static case, field transfer matrix $[F_i]$ relates the state vector $\{z\}$ of two subpoints in i th field. The point transfer matrix $[P_i]$ relates the state vector $\{z\}$ of two subpoints in i th point. The application of the method is the same in the dynamic case except the external loading is replaced by inertial forces which appear in the point transfer matrix $[P_i]$. Consideration of equation for vibration is necessary for obtaining field transfer matrix $[F_i]$. The relation between state vectors can be expressed as

$$\{z_i^-\} = [F_i]\{z_{i-1}^+\} \quad (3.2.1)$$

$$\{z_i^+\} = [P_i]\{z_i^-\} \quad (3.2.2)$$

Where $i = 1, 2, \dots, n$

From these relations, the state vectors at inner points can be eliminated and the relation between state vectors of extreme

points $(0)_+$ and $(n)_+$ can be expressed as:

$$\{z_n^+\} = [A]\{z_0^+\} \quad (3.2.3)$$

Where $[A]$ is the combined transfer matrix and is defined as

$$[A] = [P_n][F_n] \dots [P_2][F_2][P_1][F_1] \quad (3.2.3a)$$

Substituting boundary conditions in the boundary state vectors namely $\{z_n^+\}$ and $\{z_0^+\}$ in eq. 3.2.3, a set of homogeneous linear simultaneous equations are obtained. They can be expressed in matrix form as

$$[R]\{X\} = 0 \quad (3.2.4)$$

Where $\{X\}$ is a vector formed by collecting non zero terms of state vector $\{z_0^+\}$. The matrix $[R]$ is obtained from matrix $[A]$ depending on boundary condition.

For non trivial solution, the determinant of R must vanish. Thus

$$|R| = 0 \quad (3.2.5)$$

The eq. 3.2.5 is the condition to determine natural frequency of vibration of the structure.

3.3 Theoretical Analysis

The state vector is the same as used in Chapter 2. Fig. 2.3.1 and Fig. 2.3.2 are referred to for the simplified model and the illustration of notations.

3.3.1 Field Transfer Matrix

The field transfer matrix for a thin walled beam of length l is determined from the solution of differential equation of free vibration. For mono-symmetric section, the

equations are obtained by substituting $a_y = 0$ in the eq. A.5 (Appendix A). The first and the third equations are uncoupled and represent independent extensional vibration and flexural vibration in x-direction. The remaining equations representing coupled torsional and flexural vibration in y-direction are:

$$E I_x v^{IV} + \rho A v'' - \rho I_x v'''' - \rho A a_x \theta'' = 0 \quad (3.3.1a)$$

$$E I_\omega \theta^{IV} - GJ\theta'' + \rho I_p \theta'' - \rho I_\omega \theta'''' - \rho A a_x v'' = 0 \quad (3.3.1b)$$

Assuming periodic solution of the form

$$\begin{Bmatrix} v \\ \theta \end{Bmatrix} = \begin{Bmatrix} y \\ \phi \end{Bmatrix} e^{i\omega t} \quad (3.3.2a)$$

and substituting eq. 3.3.2 in eq. 3.3.1, there is obtained

$$E I_x y^{IV} - \omega^2 \rho A y + \omega^2 \rho I_x y'' + \omega^2 \rho A a_x \phi = 0 \quad (3.3.3a)$$

$$E I_\omega \phi^{IV} - GJ\phi'' - \omega^2 \rho I_p \phi + \omega^2 \rho I_\omega \phi'' + \omega^2 \rho A a_x y = 0 \quad (3.3.3b)$$

Expressing in terms of y from eq. 3.3.3a there is obtained

$$\phi = B_1 y^{IV} + B_2 y'' + B_3 y \quad (3.3.4)$$

Where

$$B_1 = -E I_x / (\omega^2 \rho A a_x);$$

$$B_2 = -I_x / (A a_x); \quad B_3 = 1/a_x$$

Eliminating ϕ from eq. 3.3.3b using eq. 3.3.4, there is obtained

$$B_4 y^{VIII} + B_5 y^{VI} + B_6 y^{IV} + B_7 y'' + B_8 y = 0 \quad (3.3.5)$$

Where

$$\left. \begin{aligned} B_4 &= -E^2 I_\omega I_x \\ B_5 &= E GJ I_x - 2w^2 \rho E I_\omega I_x \\ B_6 &= w^2 \rho (E I_\omega A + E I_p I_x + GJ I_x) - w^4 \rho^2 I_\omega I_x \\ B_7 &= -w^2 \rho GJA + w^4 \rho^2 (I_p I_x + I_\omega A) \\ B_8 &= w^4 \rho^2 (-I_p A + A^2 a_x^2) \end{aligned} \right\} (3.3.5a)$$

Assuming solution of the form

$$y = c e^{mz}$$

The characteristic equation is

$$B_4 m^8 + B_5 m^6 + B_6 m^4 + B_7 m^2 + B_8 = 0 \quad (3.3.6)$$

Let the eight roots of this polynomial are

$$m_1, m_2, m_3, \dots, m_8.$$

The solution can then be expressed as

$$v(z) = e^{i\omega t} \sum_{i=1}^8 K_i e^{m_i z}$$

$$v'(z) = e^{i\omega t} \sum_{i=1}^8 K_i m_i e^{m_i z}$$

$$M(z) = E I_x v'' = e^{i\omega t} \left(\sum_{i=1}^8 K_i m_i^2 e^{m_i z} \right) E I_x$$

$$V(z) = -E I_x v''' = -e^{i\omega t} E I_x \sum_{i=1}^8 K_i m_i^3 e^{m_i z}$$

$$\theta(z) = e^{i\omega t} \sum_{i=1}^8 (B_1 m_i^4 + B_2 m_i^2 + B_3) K_i e^{m_i z}$$

$$\theta'(z) = e^{i\omega t} \sum_{i=1}^8 (B_1 m_i^5 + B_2 m_i^3 + B_3 m_i) K_i e^{m_i z}$$

$$B(z) = -E I_\omega \theta'' = -e^{i\omega t} E I_\omega \sum_{i=1}^8 (B_1 m_i^6 + B_2 m_i^4 + B_3 m_i^2) K_i e^{m_i z}$$

$$H(z) = -E I_{\omega} \theta'''' + GJ\theta'$$

$$= e^{i\omega t} \sum_{i=1}^8 \{ -E I_{\omega} B_i m_i^7 + (-E I_{\omega} B_2 + GJB_1) m_i^5 + (-E I_{\omega} B_3 + GJB_2) m_i^3 + GJB_3 m_i \} e^{m_i z}$$

Where K_1, K_2, \dots, K_8 are the constants to be determined from boundary conditions.

The above equations can be expressed in matrix form as

$$\{Z(z)\} = [C][D(z)][K] e^{i\omega t} \quad (3.3.8)$$

Where $Z(z)$ is the state vector at a distance z from the origin and defined as

$$Z(z) = \begin{pmatrix} V(z) \\ V'(z) \\ M(z) \\ V(z) \\ \theta(z) \\ \theta'(z) \\ B(z) \\ H(z) \end{pmatrix} \quad (3.3.8a)$$

$[C]$ is a 8×8 square matrix with elements as follows

$$c(1,i) = 1$$

$$c(2,i) = m_i$$

$$c(3,i) = E I_x m_i^2$$

$$c(4,i) = -E I_x m_i^3$$

$$c(5,i) = B_1 m_i^4 + B_2 m_i^2 + B_3$$

$$c(6,i) = B_1 m_i^5 + B_2 m_i^3 + B_3 m_i$$

$$c(7,i) = -E I_{\omega} (B_1 m_i^6 + B_2 m_i^4 + B_3 m_i^2)$$

$$c(8,i) = -E I_{\omega} B_1 m_i^7 + (-E I_{\omega} B_2 + GJB_1) m_i^5 \\ + (-E I_{\omega} B_3 + GJB_2) m_i^3 + GJB_3 m_i$$

where $i = 1, 2, \dots, 8$

$[D(z)]$ is a 8×8 diagonal matrix with diagonal elements as

$$d(i,i) = e^{m_i z}, \text{ where } i=1, 2, \dots, 8$$

$\{K\}$ is a column matrix consisting of constants K_i $i=1, 2, \dots, 8$

The boundary conditions at the base $z = 0$ are

$$v = v(0), v' = v'(0), M = M(0), V = V(0)$$

$$\theta = \theta(0), \theta' = \theta'(0), B = B(0), H = H(0)$$

Substituting these conditions in eq. 3.3.8, there is obtained

$$\{Z(0)\} = [C][I]\{K\} e^{i\omega t} \quad (3.3.9)$$

The constants can be determined by matrix inversion in eq. 3.3.9

$$\{K\} = [C]^{-1}\{Z(0)\} e^{-i\omega t} \quad (3.3.9a)$$

The boundary conditions at $z = l$ are

$$v = v(l), v' = v'(l), M = M(l), V = V(l)$$

$$\theta = \theta(l), \theta' = \theta'(l), B = B(l), H = H(l)$$

Substituting these conditions in eq. 3.3.8 and using eq. 3.3.9 to eliminate $\{K\}$

$$\{Z(l)\} = [C][D(l)][C]^{-1}\{Z(0)\} \quad (3.3.10)$$

The field transfer matrix is a 8×8 square matrix obtained from matrix multiplication

$$[F] = [C][D(l)][C]^{-1} \quad (3.3.11)$$

3.3.2 Point Transfer Matrix

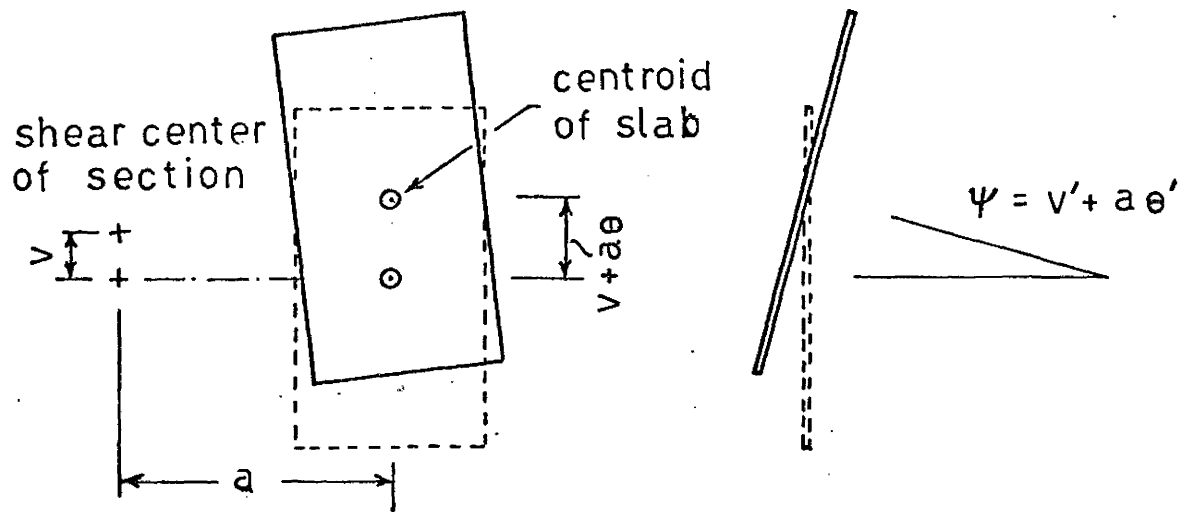
The displacement of the center of gravity of the slab is shown in Fig. 3.3.1a. The inertia forces due to motion of the slab is shown in Figs. 3.3.1b to 3.3.1d. Stiffening action of the slab to contribute bimoment is shown in Fig. 3.3.1e.

Consideration of equilibrium of the slab element yields the following equations.

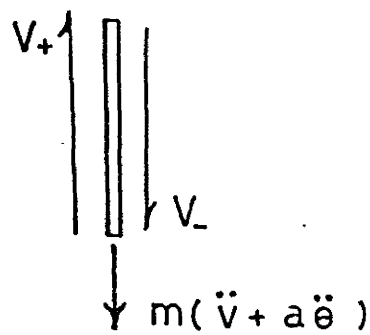
$$\left. \begin{aligned} V_+ &= V_- + mv'' + am\theta'' \\ H_+ &= H_- + J_m\theta'' + amv'' + a^2m\theta'' \\ M_+ &= M_- + J_x\psi'' = M_- + J_xv'' + aJ_x\theta'' \\ B_+ &= B_- - B_s = B_- - D\theta' \end{aligned} \right\} \quad (3.3.12)$$

Notations used in the above eqs. are

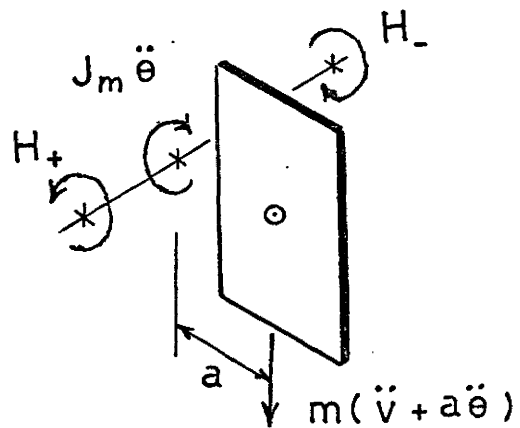
- a Distance between shear center of the section and the center of gravity of the slab
- m Mass of the slab
- J_m Polar mass moment of inertia of the slab about an axis through the center of gravity.
- J_x Mass moment of inertia of the slab about an axis parallel to x and passing through the center of gravity
- D Bimoment contribution factor defined in eq. 2.3.12



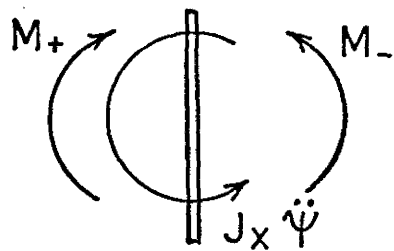
(a) Displacement of Centroid of Slab



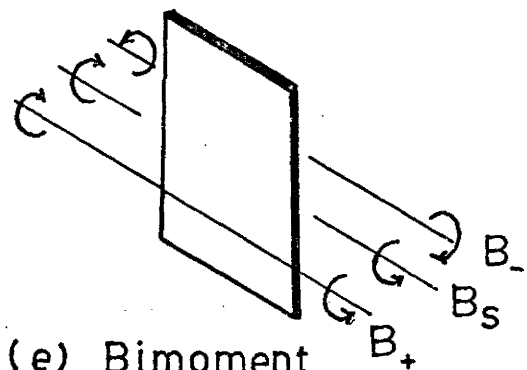
(b) Shear Force from Slab



(c) Torque from Slab



(d) Moment from Slab



(e) Bimoment from Slab

FIG. 3.3.1

For a periodic vibration of frequency w

$$\begin{pmatrix} v(z,t) \\ v'(z,t) \\ \theta(z,t) \\ \theta'(z,t) \end{pmatrix} = \begin{pmatrix} y(z) \\ y'(z) \\ \phi(z) \\ \phi'(z) \end{pmatrix} e^{iwt}$$

Differentiating twice with respect to time

$$\begin{pmatrix} v''(z,t) \\ v'''(z,t) \\ \theta''(z,t) \\ \theta'''(z,t) \end{pmatrix} = -w^2 \begin{pmatrix} v(z,t) \\ v'(z,t) \\ \theta(z,t) \\ \theta'(z,t) \end{pmatrix} \quad (3.3.13)$$

Substituting in eq. 3.3.12, there is obtained

$$\left. \begin{aligned} V_+ &= V_- - w^2 m v - w^2 a m \theta \\ H_+ &= H_- - w^2 a m v - w^2 (J_m + a^2 m) \theta \\ M_+ &= M_- - w^2 J_x v' - w^2 a J_x \theta' \\ B_+ &= B_- - D \theta' \end{aligned} \right\} \quad (3.3.14)$$

Compatibility conditions give

$$\left. \begin{aligned} v_+ &= v_- \\ v'_+ &= v'_- \\ \theta_+ &= \theta_- \\ \theta'_+ &= \theta'_- \end{aligned} \right\} \quad (3.3.15)$$

Point transfer matrix $[P]$ can there be formed from eq. 3.3.14 and eq. 3.3.15 as

$$[P] = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & p_{32} & 1 & 0 & 0 & p_{36} & 0 & 0 \\ p_{41} & 0 & 0 & 1 & p_{45} & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & p_{76} & 1 & 0 \\ p_{81} & 0 & 0 & 0 & p_{85} & 0 & 0 & 1 \end{bmatrix} \quad (3.3.16)$$

The elements of matrix are

$$p_{32} = -w^2 J_x, \quad p_{36} = -w^2 a J_x$$

$$p_{41} = -w^2 m, \quad p_{45} = -w^2 a m$$

$$p_{76} = -D, \quad p_{81} = -w^2 a m$$

$$p_{85} = -w^2 (J_m + a^2 m)$$

3.3.3 Boundary Conditions

The boundary conditions for the shear wall fixed at base and free at top can be written as

$$\{z_o^+\} = \begin{bmatrix} 0 \\ 0 \\ M_o^+ \\ 0 \\ V_o^+ \\ 0 \\ 0 \\ B_o^+ \\ H_o^+ \\ 0 \end{bmatrix} \quad \text{and} \quad \{z_8^+\} = \begin{bmatrix} v_8^+ \\ v_8'^+ \\ 0 \\ 0 \\ \theta_8^+ \\ \theta_8'^+ \\ 0 \\ 0 \end{bmatrix}$$

Substituting these conditions in eq. 3.2.3, there is obtained

$$\begin{bmatrix} v_8^+ \\ v_8'^+ \\ 0 \\ 0 \\ \theta_8^+ \\ \theta_8'^+ \\ 0 \\ 0 \end{bmatrix} = [A] \begin{bmatrix} 0 \\ 0 \\ M_O^+ \\ V_O^+ \\ 0 \\ 0 \\ B_O^+ \\ H_O^+ \end{bmatrix} \quad (3.3.17)$$

Re-arrangement of terms of the above eq. yields

$$[R]\{X\} = 0 \quad (3.3.18)$$

Where

$$[R] = \begin{bmatrix} a_{33} & a_{34} & a_{37} & a_{38} \\ a_{43} & a_{44} & a_{47} & a_{48} \\ a_{73} & a_{74} & a_{77} & a_{78} \\ a_{83} & a_{84} & a_{87} & a_{88} \end{bmatrix}$$

Here 'a' denotes elements of matrix [A].

$$\{X\} = \begin{bmatrix} M_O^+ \\ V_O^+ \\ B_O^+ \\ H_O^+ \end{bmatrix}$$

The condition to determine the natural frequencies is

$$|R| = 0 \quad (3.3.19)$$

After determining the natural frequencies, the relative values of the elements in vector $\{X\}$ can be determined. The mode shape of the structure follows by back substitution of the vector $\{X\}$.

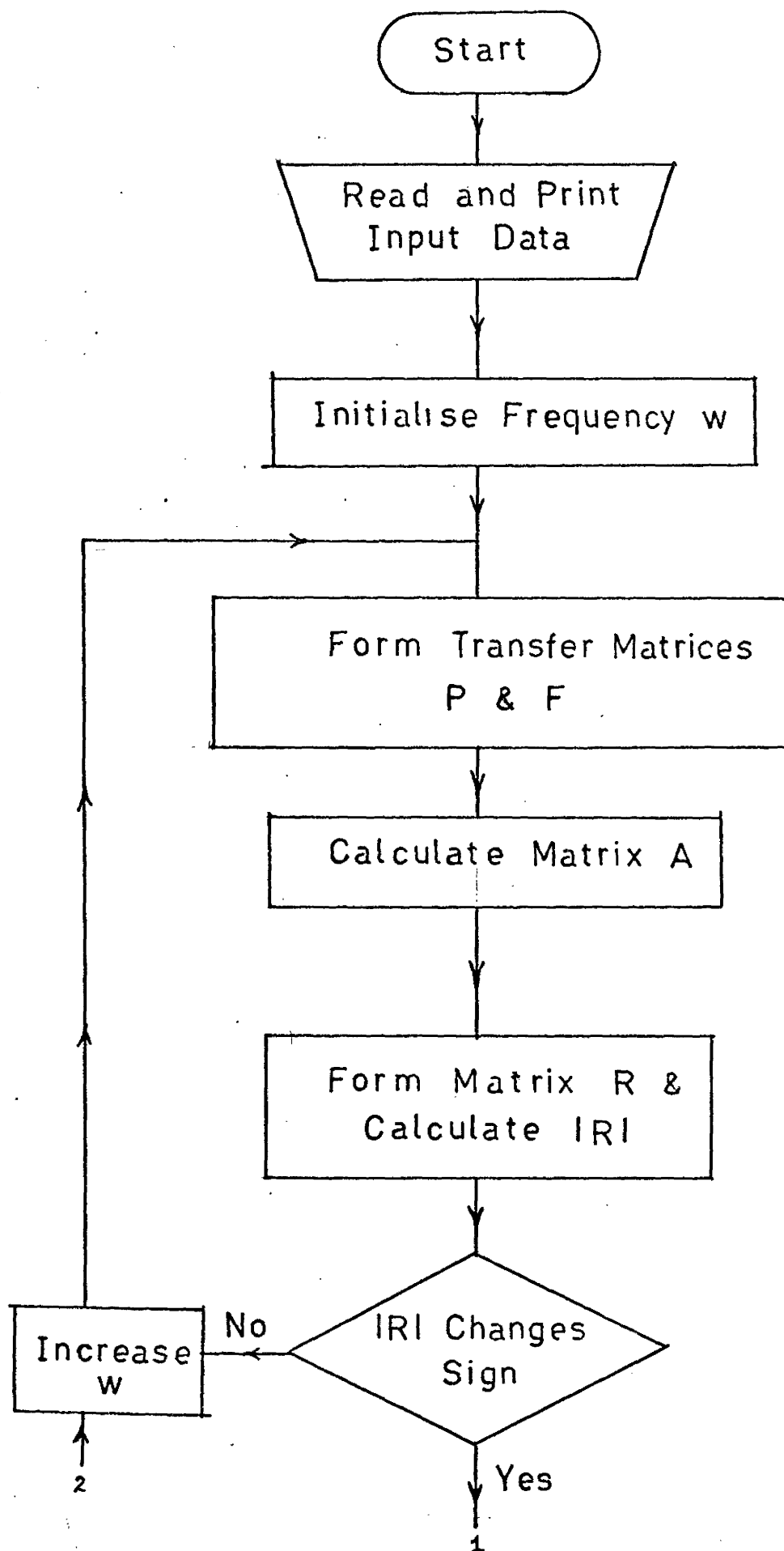
3.4 Computer Program

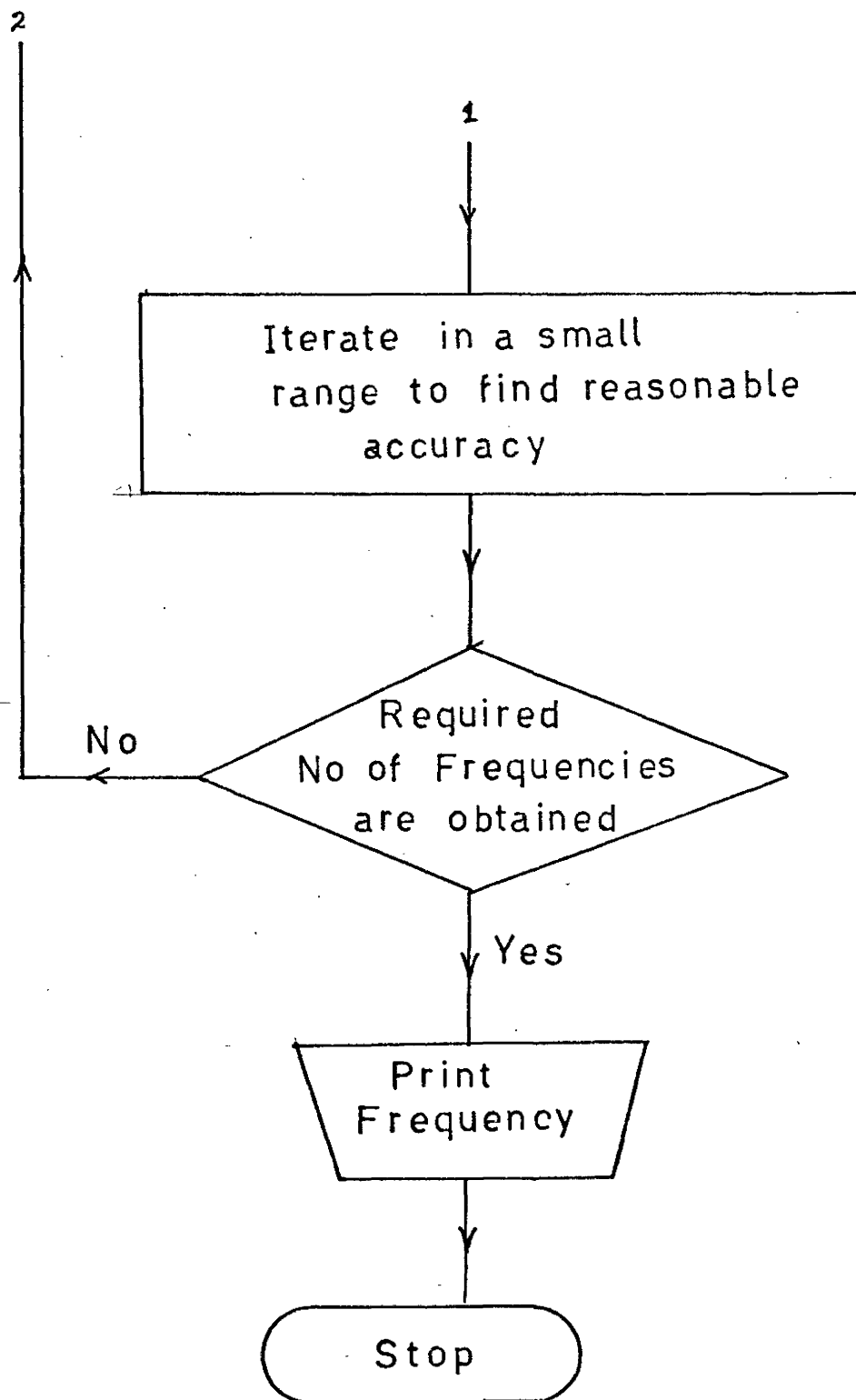
A computer program based on the above analysis has been written. The input data are the geometric and elastic properties of the shear wall structure. The trial frequency w is increased from an initial value and $|R|$ (eq. 3.3.19) is calculated. If $|R| \neq 0$ another value of w is tried. The same procedure is repeated until $|R|$ is reasonably small to be considered as zero. The next higher frequency is then determined following the same scheme. These natural frequencies obtained are then used as inputs in a second program to determine associated mode shapes.

The flow chart for the first program is shown in fig. 3.4.1 and the computer program is included in Appendix B.

3.5 Experiment

A dynamic experiment was carried out on the model (Fig. 2.1.1). The model was fixed on the shaking table and subjected to lateral vibration of known amplitude while the frequency is gradually swept from an initial value upwards. In all the dynamic tests, the shaking table was subjected





FLOW CHART

FIG. 3.4.1

to a constant displacement of 0.005 inch from 10 cps to 44 cps. After cross over frequency of 44 cps it was subjected to constant acceleration of 0.5 g. The response of the accelerometers fixed at different points of the model was studied. The experimental set up is shown in fig. 3.5.1 to 3.5.3. The output from the accelerometers was viewed in an oscilloscope. The RMS response of the accelerometers were plotted in a XY recorder. D.C. voltage proportional to RMS acceleration was fed in X ordinate and Y ordinate was adjusted in a suitable time scale. The arrangement of instruments is shown in fig. 3.5.4. Marked increase in response is noticed in the frequency response plots (Fig. 3.6.1 to 3.6.4) at the resonant frequencies. For locating the resonant peak more accurately, the frequency was manually changed around each resonance zone and the response was monitored on a RMS voltmeter. The strains at resonance were determined from plotting outputs from the strain gauges in the Visicorder. The arrangement of instrument used is shown in fig. 3.5.5. Relative strain distribution is drawn from these plots and shown in fig. 3.6.7 to fig. 3.6.10.

The following is a list of different instruments used in the experiment.

The dynamic testing set-up consists of

A. Sweep Oscillator SD104A-5D

Make: Spectral Dynamic Corporation



Fig. 3.5.1
EXPERIMENTAL SET UP FOR DYNAMIC TEST



Fig. 3.5.2
MODEL FIXED ON THE SHAKING TABLE

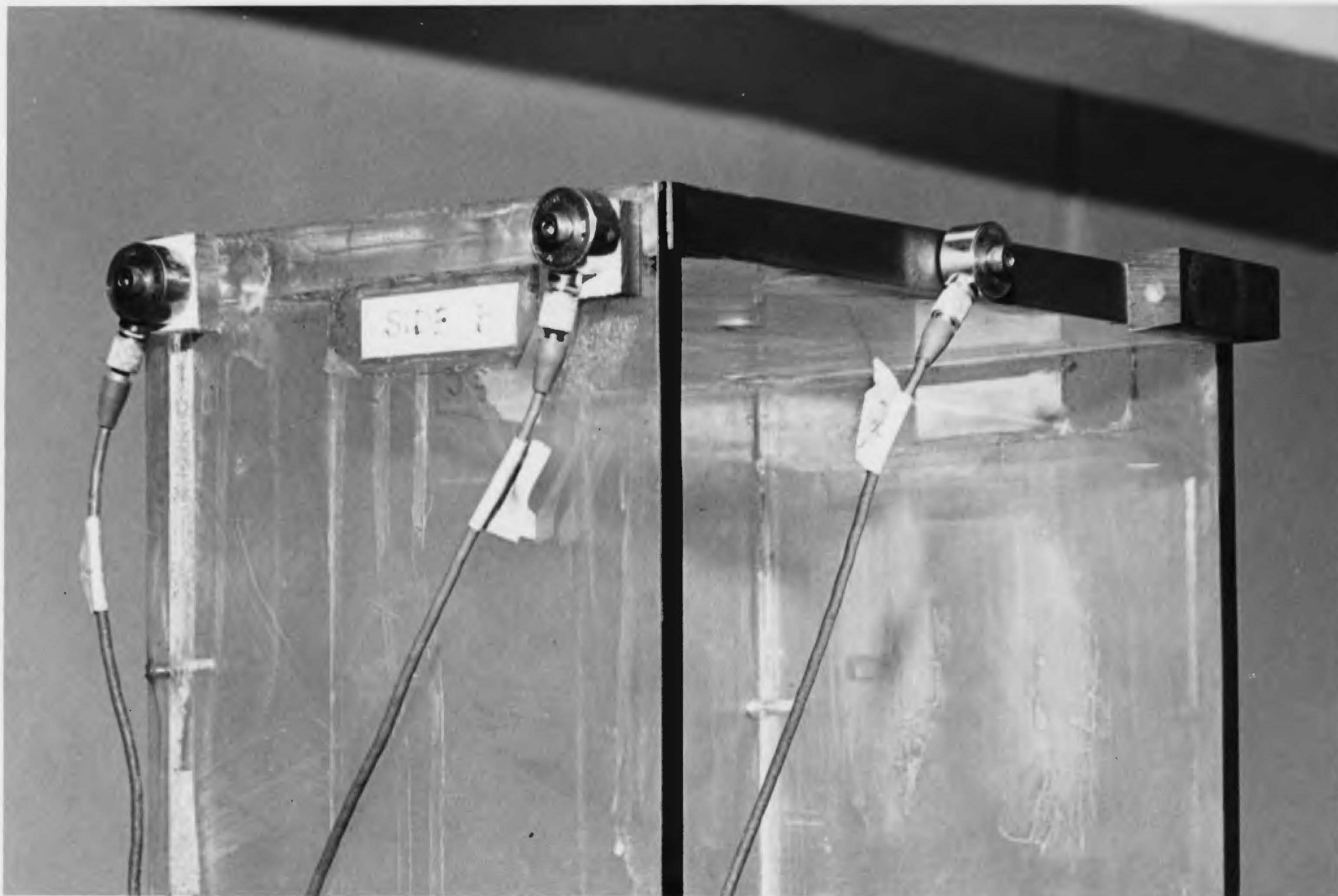
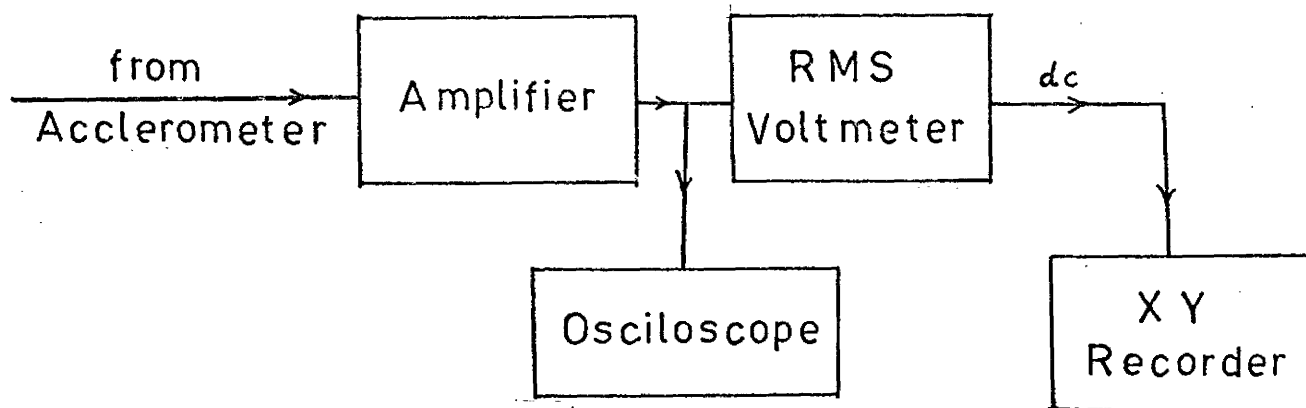


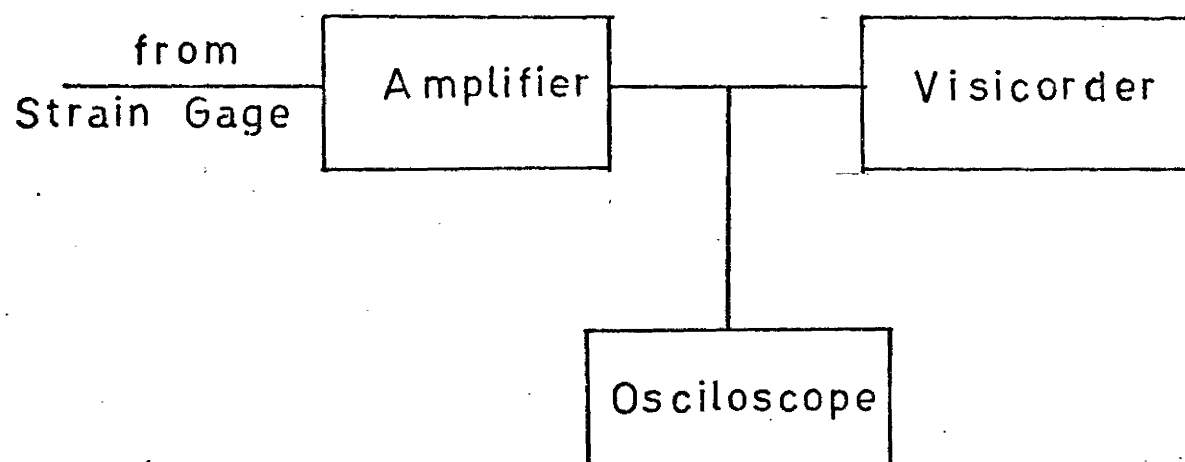
Fig. 3.5.3

ACCELEROMETERS ATTACHED TO THE MODEL



ARRANGEMENT FOR RESONANCE SEARCH

FIG. 3.5.4



ARRANGEMENT FOR MEASURING STRAINS

FIG. 3.5.5

- B. Amplitude Servo/Monitor SD105A
Make: Spectal Dynamic Corporation.
- D. Accelerometer source follower SFA-100
Make: Ling Electronics
- E. Accelerometer Normalizing Amplifier ANA-101
Make: Ling Electronics
- F. Power Amplifier CP-5/6
Make: Ling Electronics
- G. Shaker B 290
Make: Ling Electronics
- H. Vibraglide Sliptable SINGCO 30-30
Make: Marshall Research and Dvelopment Corporation

For monitoring and plotting the response the following instruments were employed.

- I. Dual Beam Oscilloscope
Make: Tectronix Inc.
- J. Accelerometer
- K. Laboratory amplifier 2616B
Make: Endevco Corporation
- L. D.C. Amplifier, High Gain Type 1-165
Make: Endevco Corporation
- M. Bridge Amplifier
Make: Ellis Associates
- N. RMS Voltmeter
Make: Hewlett Packard
- O. Visicorder (2 channels)
Make: Honeywell Controls Ltd.

P. Digital Counter

Make: Hewlet Packard.

3.6 Results and Discussions

The frequency response plots as obtained from the dynamic test for the accelerometers located at different positions are shown in fig. 3.6.1 to 3.6.4. The figures near the peaks are the experimental resonant frequencies. Average of all the experimental frequencies together with the theoretical frequencies for different consideration of floor slabs are shown in table 3.6.1.

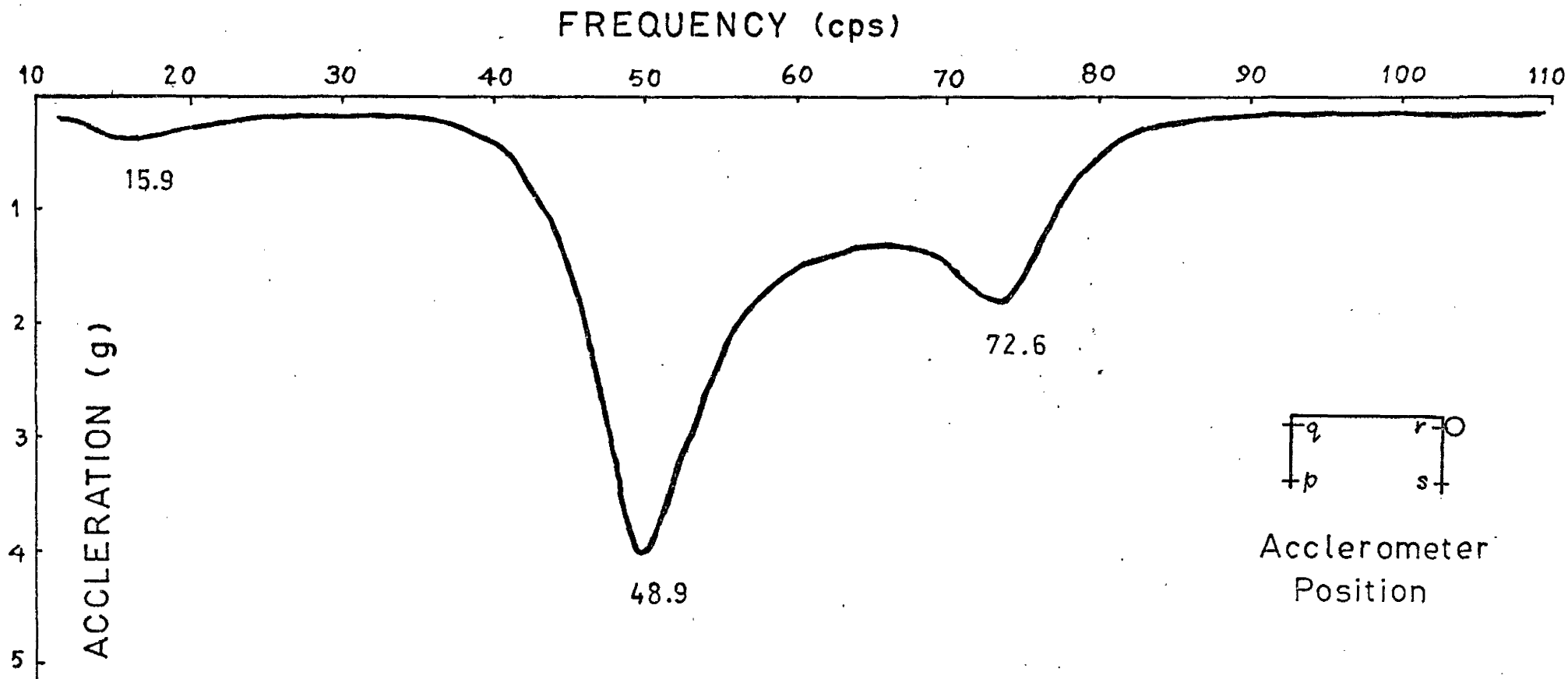
It can be seen that for the first and third modes the experimental freq. lies between the theoretical predicted value when the torsional restraining effect of the slab is considered and the theoretical predicted value when both the torsional and effective bending restraining effect of the slab is considered. The difference between the theoretical and experimental values is about 5%.

In the second mode, the experimental value is 16% lower than the theoretical calculated value. Since the second mode is a bending predominant mode, in this case, the larger difference may be caused by neglecting shear deformation in the mathematical model. The importance of considering shear deformation for bending predominant mode was noted by Tso (29).

NO. FREQ.	THEORETICAL FREQUENCY (cps)				EXPERIMENTAL FREQUENCY (cps)
	Floor -less Structure	Torsion Only of Slab	Torsion + Eff. Bending of Slab	Torsion + Bending of Slab	
1	16.5	14.5	16.8	20.8	15.9
2	65.6	56.7	57.8	60.9	48.7
3	93.9	71.8	76.5	89.0	72.9

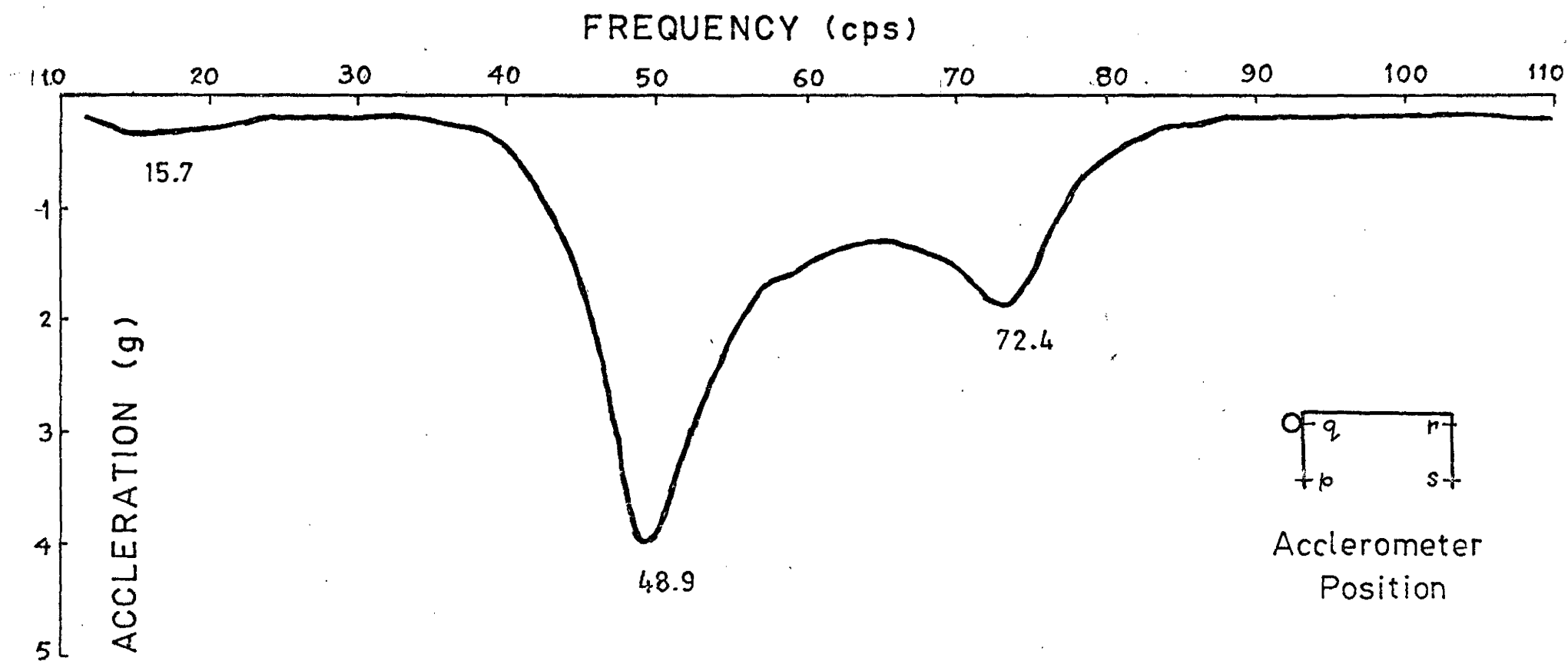
COMPARISON OF FREQUENCIES

TABLE 3.6.1



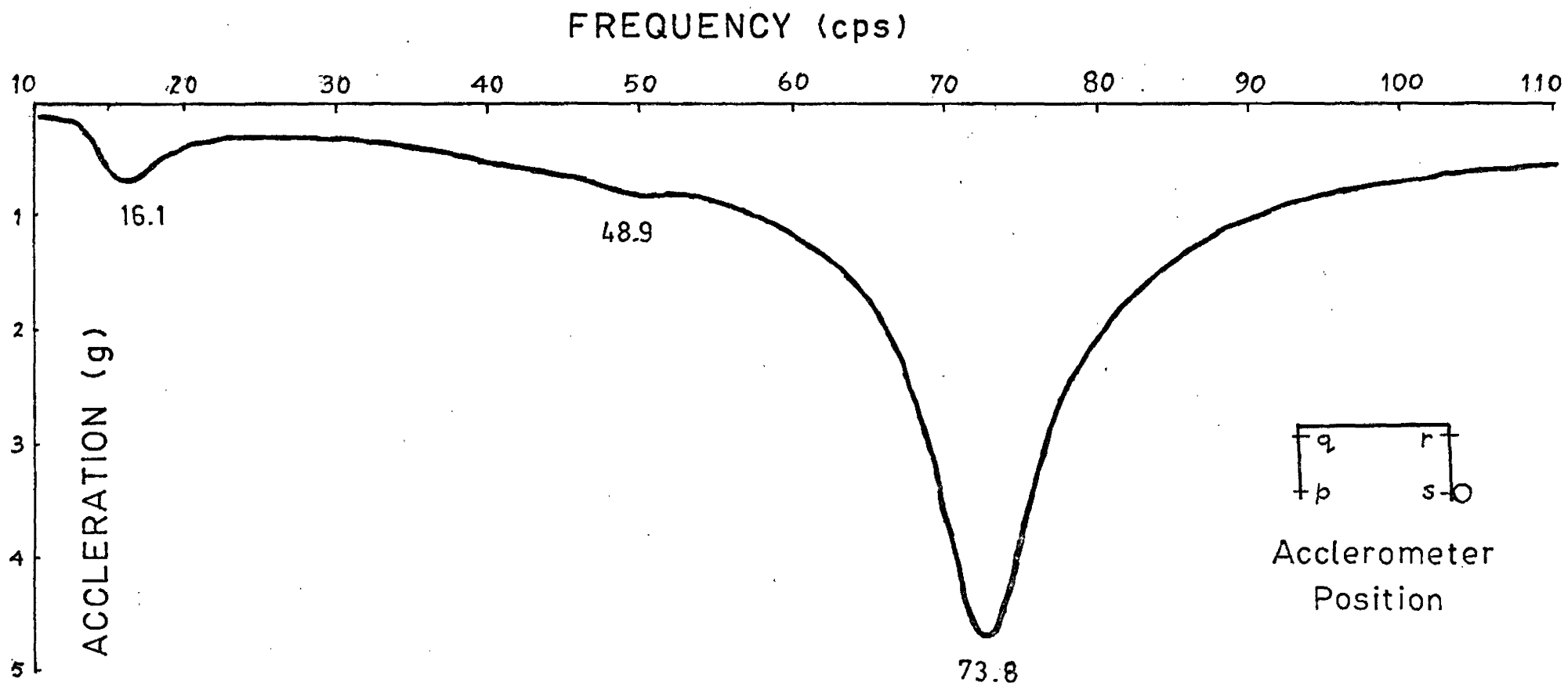
FREQUENCY RESPONSE PLOT

FIG. 3.6.1



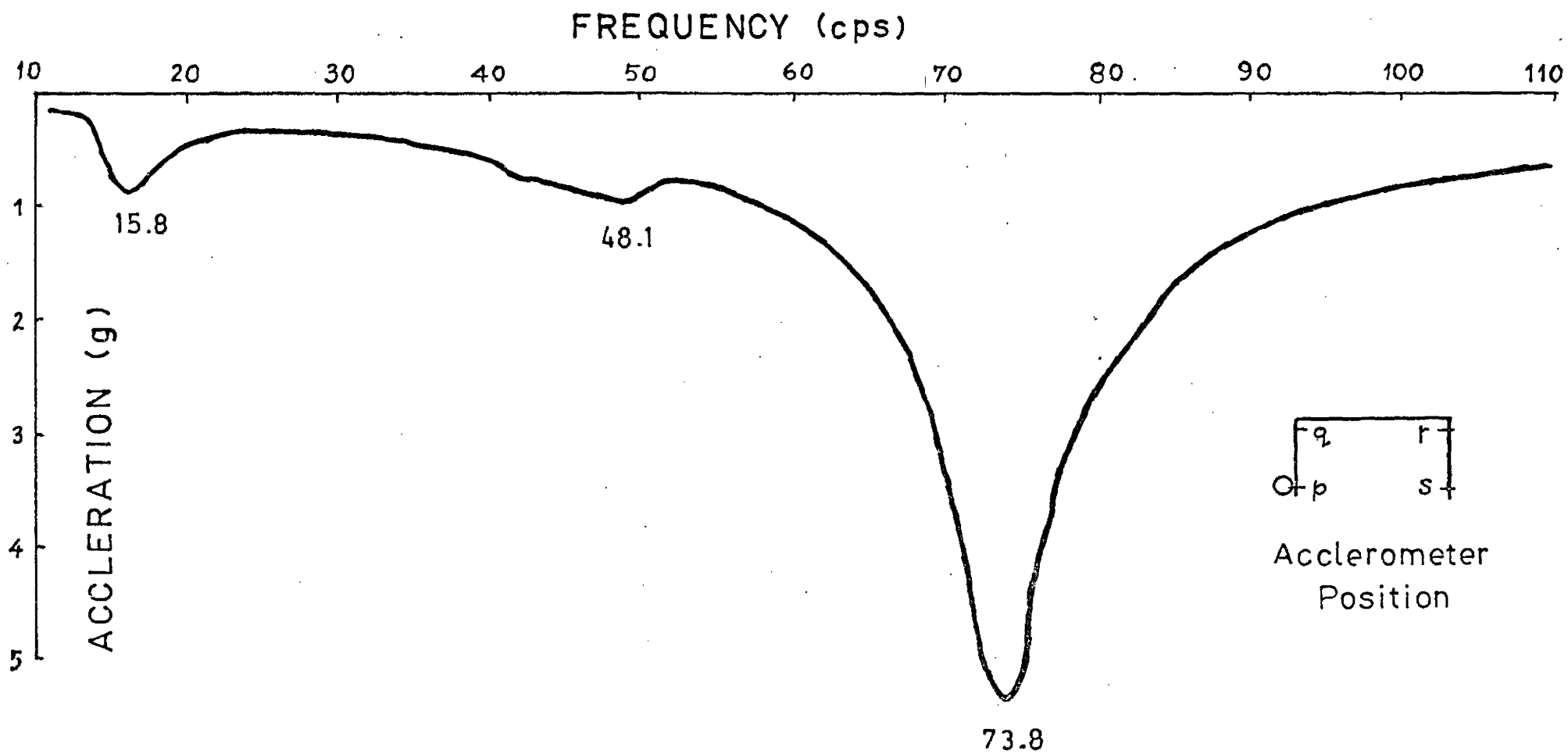
FREQUENCY RESPONSE PLOT

FIG. 3.6.2



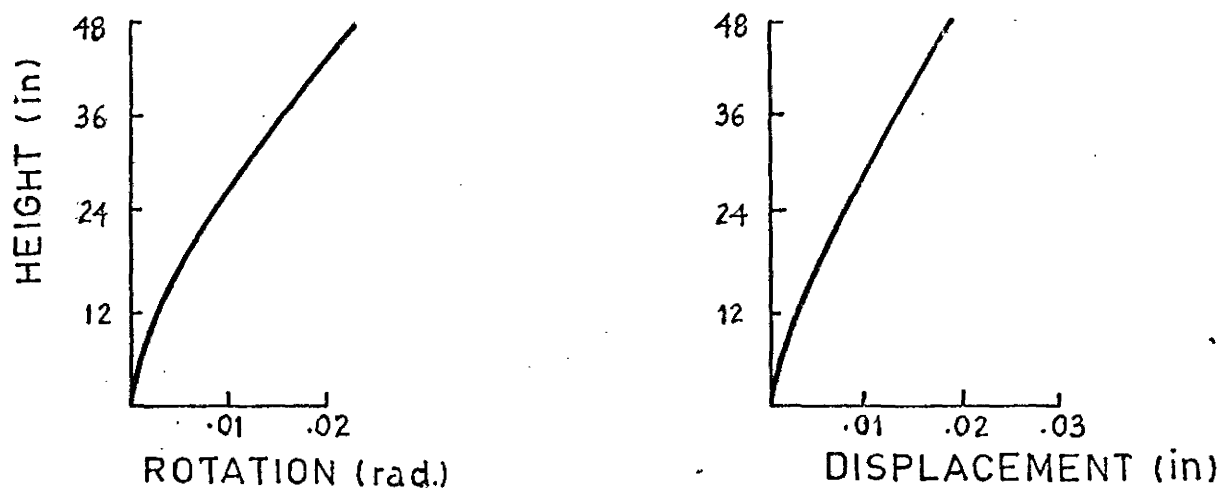
FREQUENCY RESPONSE PLOT

FIG 3.6.3

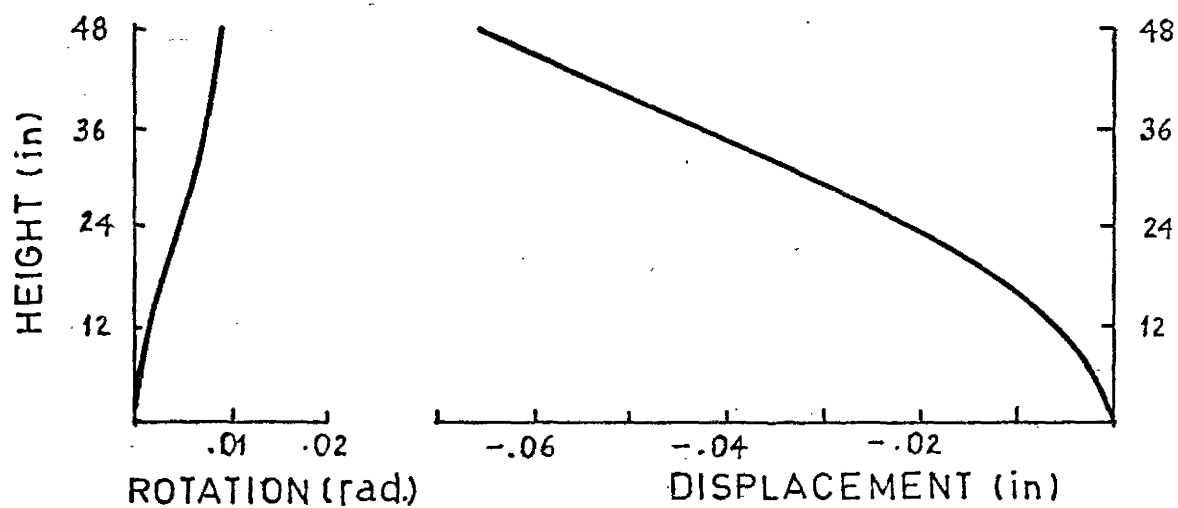


FREQUENCY RESPONSE PLOT

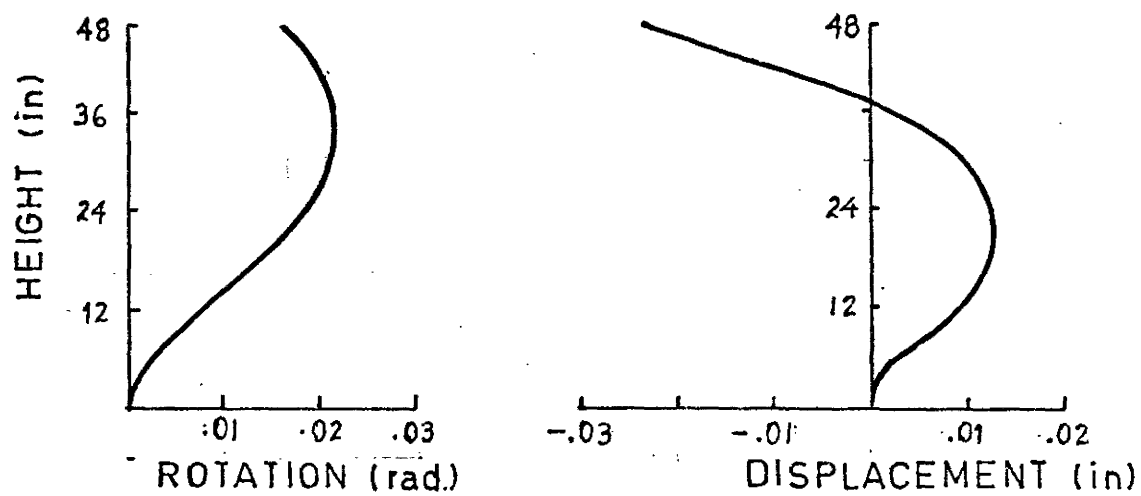
FIG. 3.6.4



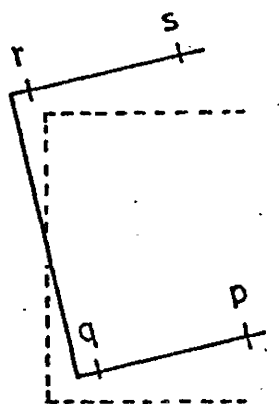
(a) MODE SHAPE FOR FIRST FREQ.



(b) MODE SHAPE FOR SECOND FREQ.



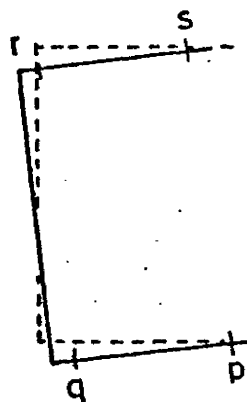
(c) MODE SHAPE FOR THIRD FREQ.



$$v = 0.18 \text{ in}$$

$$\theta = 13.6 \text{ degree}$$

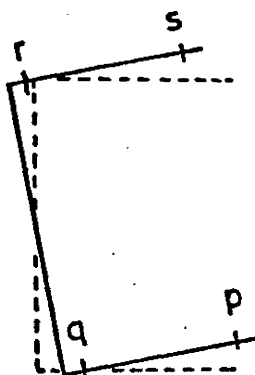
(a) FIRST MODE



$$v = -0.66 \text{ in}$$

$$\theta = 5.5 \text{ degree}$$

(b) SECOND MODE

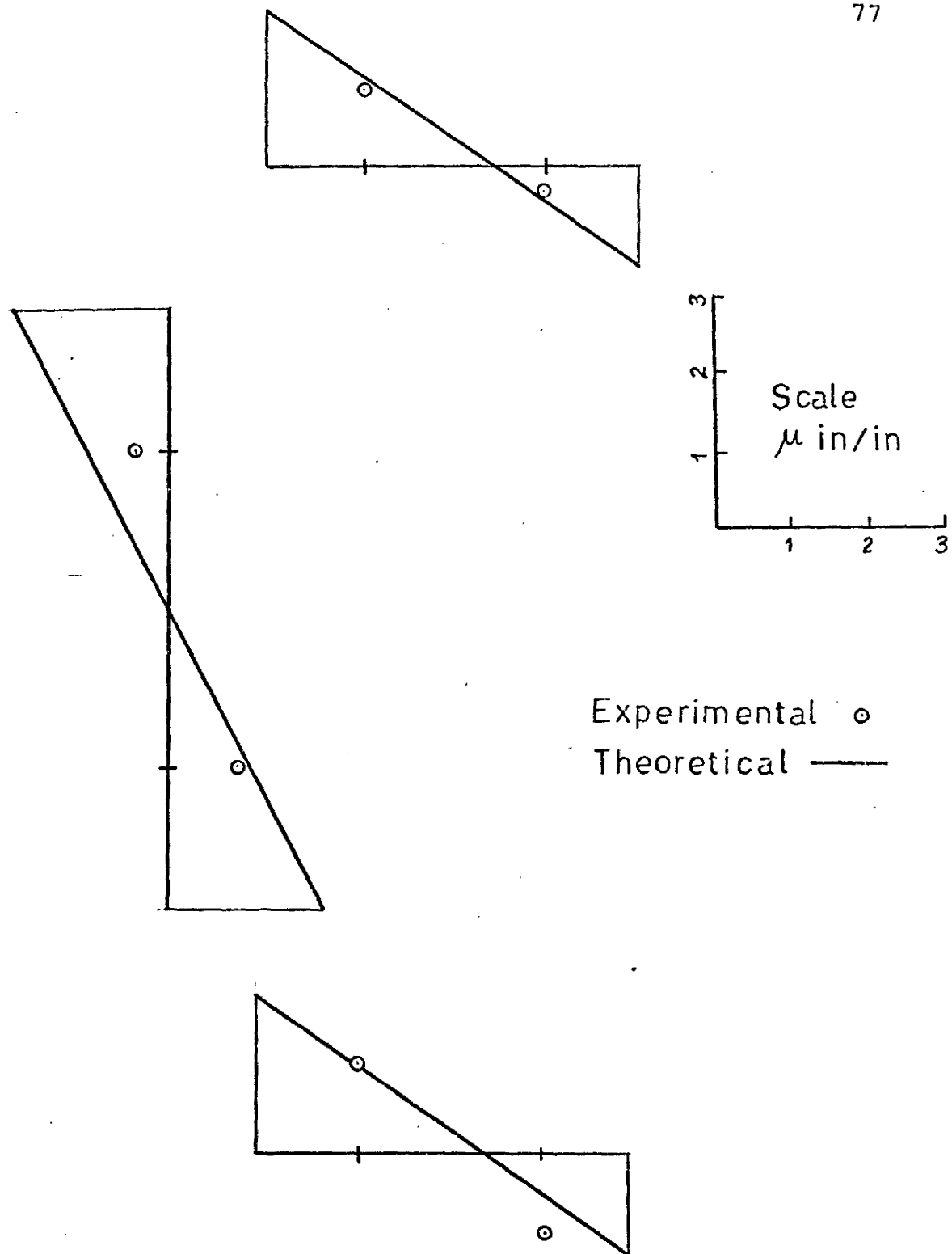


$$v = -0.237 \text{ in}$$

$$\theta = 9.3 \text{ degree}$$

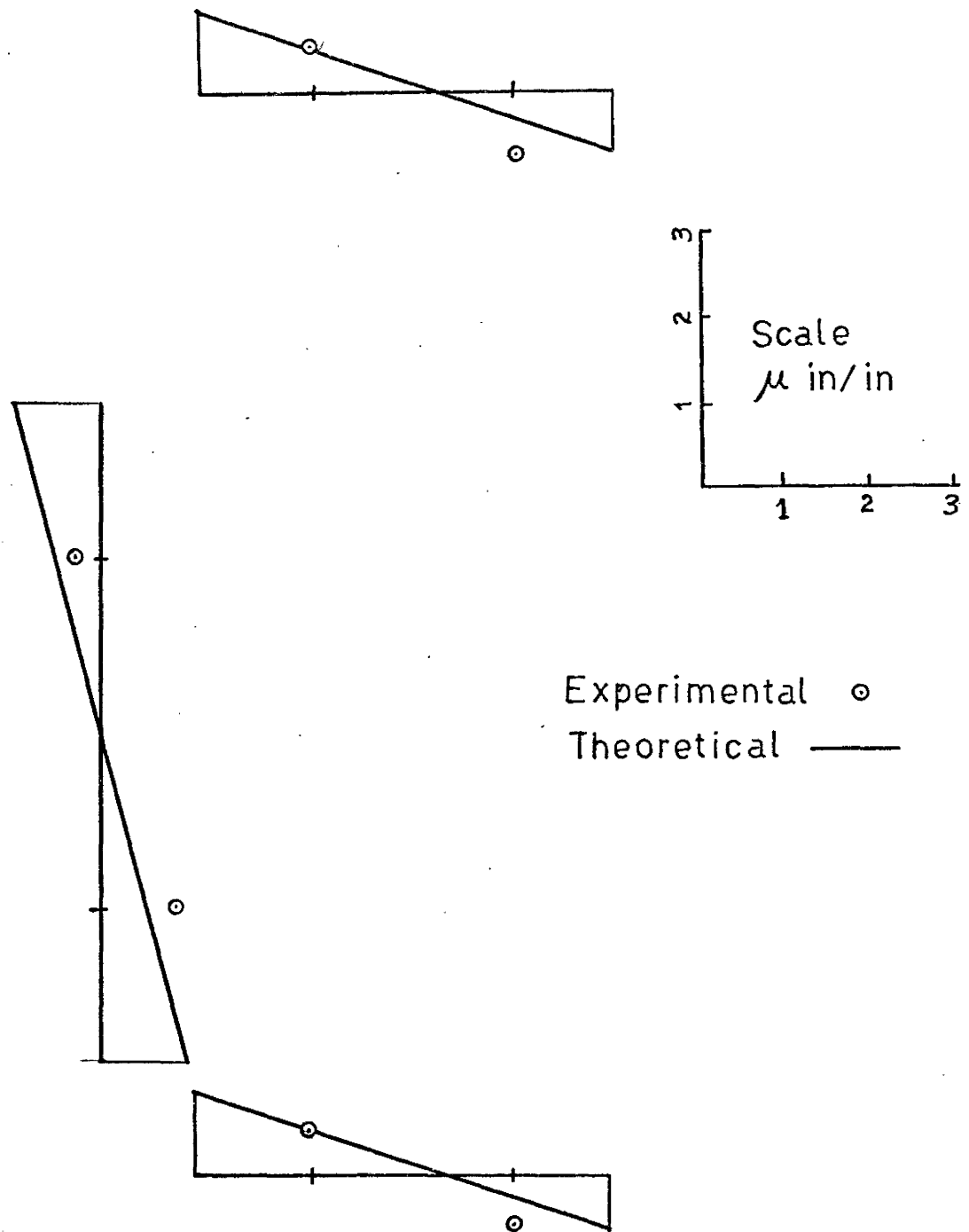
(c) THIRD MODE

RELATIVE ROTATION & DISPLACEMENT
AT 8TH. FLOOR
FIG. 3.6.6



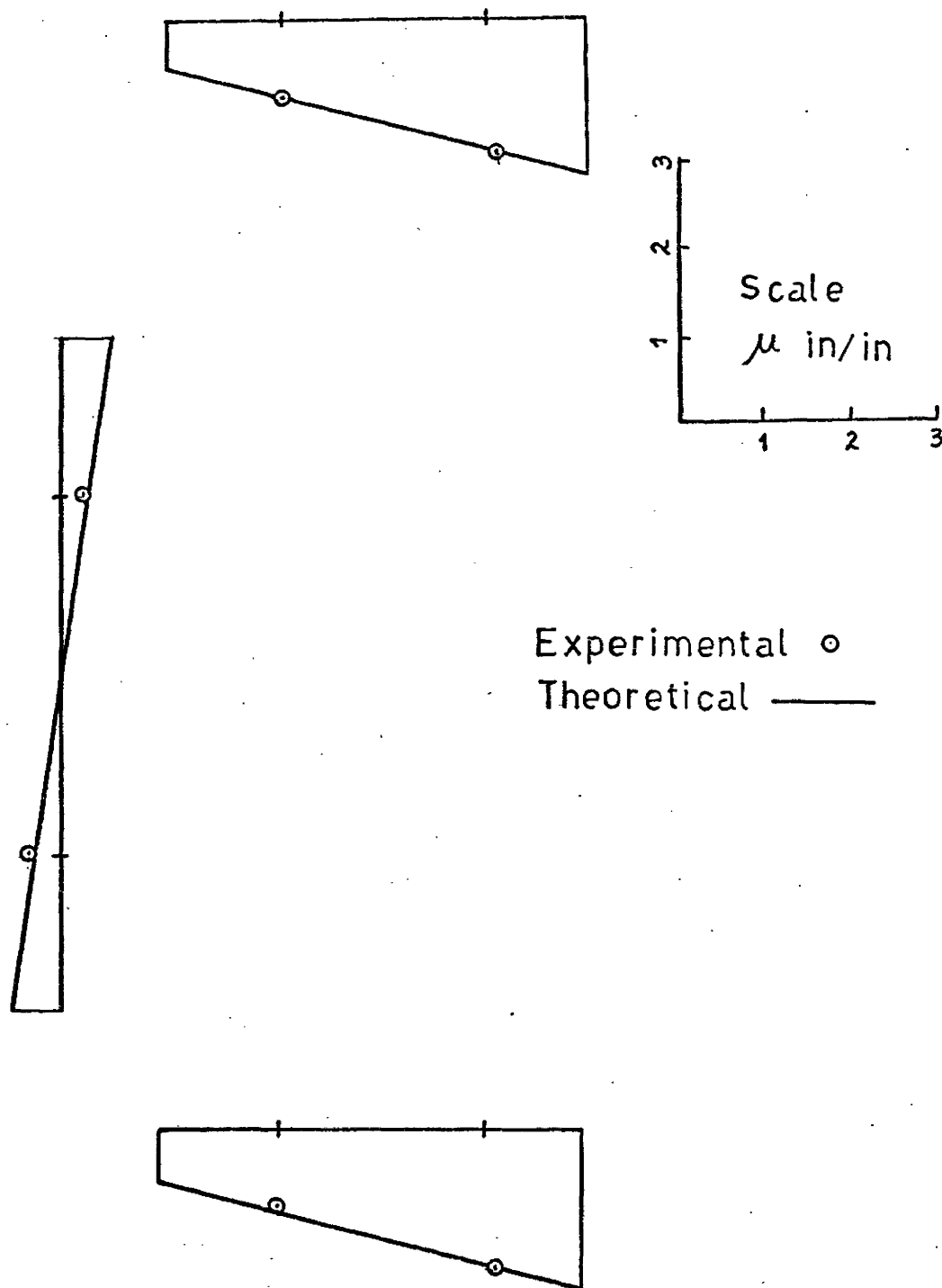
RELATIVE STRAIN DISTRIBUTION AT LEVEL A
AT FIRST FREQUENCY

FIG. 3.6.7



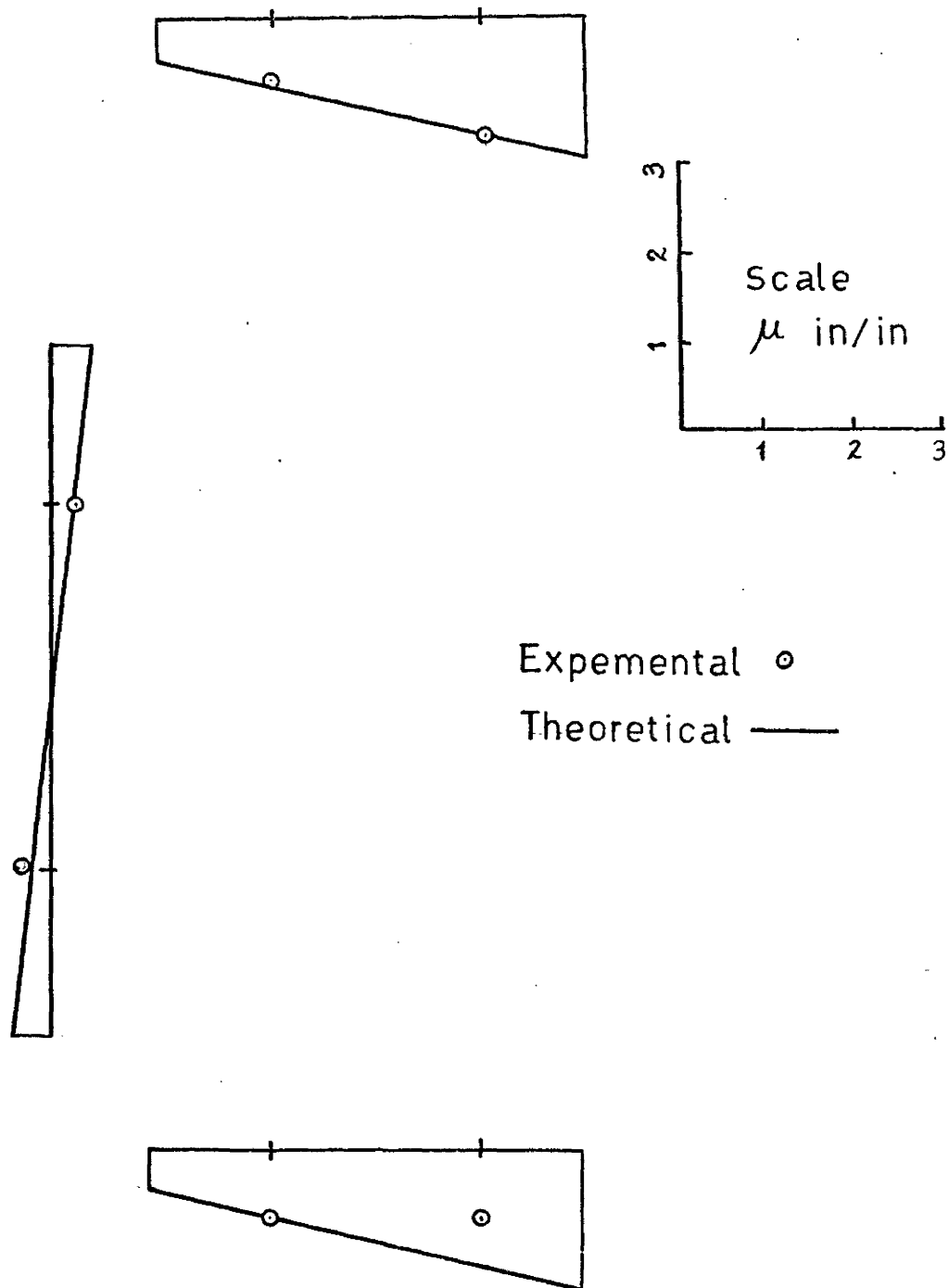
RELATIVE STRAIN DISTRIBUTION AT LEVEL B
AT FIRST FREQUENCY

FIG. 3.6.8



RELATIVE STRAIN DISTRIBUTION AT LEVEL A
AT SECOND FREQUENCY

FIG. 3.6.9



RELATIVE STRAIN DISTRIBUTION AT LEVEL B
AT SECOND FREQUENCY

FIG. 3.6.10

The associated mode shapes of the structure for first three modes obtained theoretically are plotted in fig. 3.6.5. It shows from these plots that the first mode consists of torsion predominant displacement and the second mode consists of bending predominant displacements.

It is also noted that in first mode displacement and rotation are in phase. In second mode they are out of phase. In third mode they are in phase in lower part of the structure but out of phase in top part.

The displacements and rotations at the top of the structure for different modes are plotted in fig. 3.6.6. The values of v and θ written on the figure are the values of mode shapes curve (Fig. 3.6.5) at 48 inch levels are relative values only. In other words the values for the first mode does not have any relation with that for the second or the third mode.

It is seen from the frequency response plots (fig. 3.6.1 to fig. 3.6.4) that at any resonant frequency, the response varies depending on the position of the accelerometer. For example in the third resonance the accelerometers mounted at points q and r show lower response than that mounted at points p and s. In the second resonance, the accelerometers mounted at points q and r shows higher response than that mounted at points p and s. In the first resonance, the response of accelerometers mounted at p and s shows higher response than that mounted at q and r. These

type of behaviour is due to the in phase or out of phase nature of displacements and rotations in different modes and can be explained from fig. 3.6.6. It is seen from the plot that points p and s undergo maximum translational displacement due to combined action of v and θ in the third and the first mode. Whereas points q and r undergo maximum translational displacement in second mode only.

The relative strain distribution at level A and level B (refer fig. 2.5.2 for level A and B) at the first and the second resonance are plotted (Fig. 3.6.7 to Fig. 3.6.10). Theoretical distribution are drawn from the calculated mode shapes. Since the mode shapes are defined by relative values only the theoretical strain diagram is made to pass through one experimental point. A reasonable agreement between the theory and experiment is shown in these plots.

CHAPTER IV

STATIC FORMULATION OF NON PLANAR COUPLED SHEAR WALL

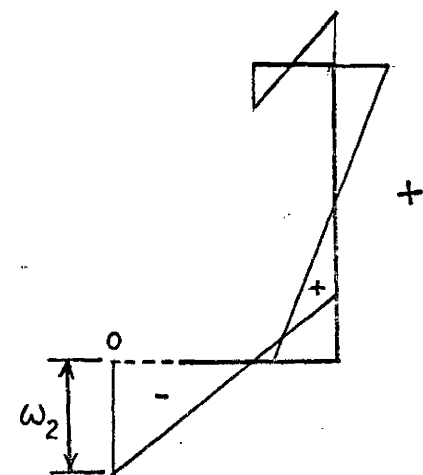
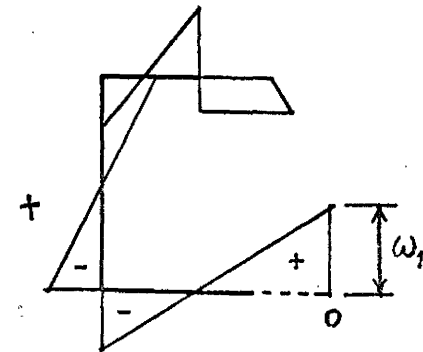
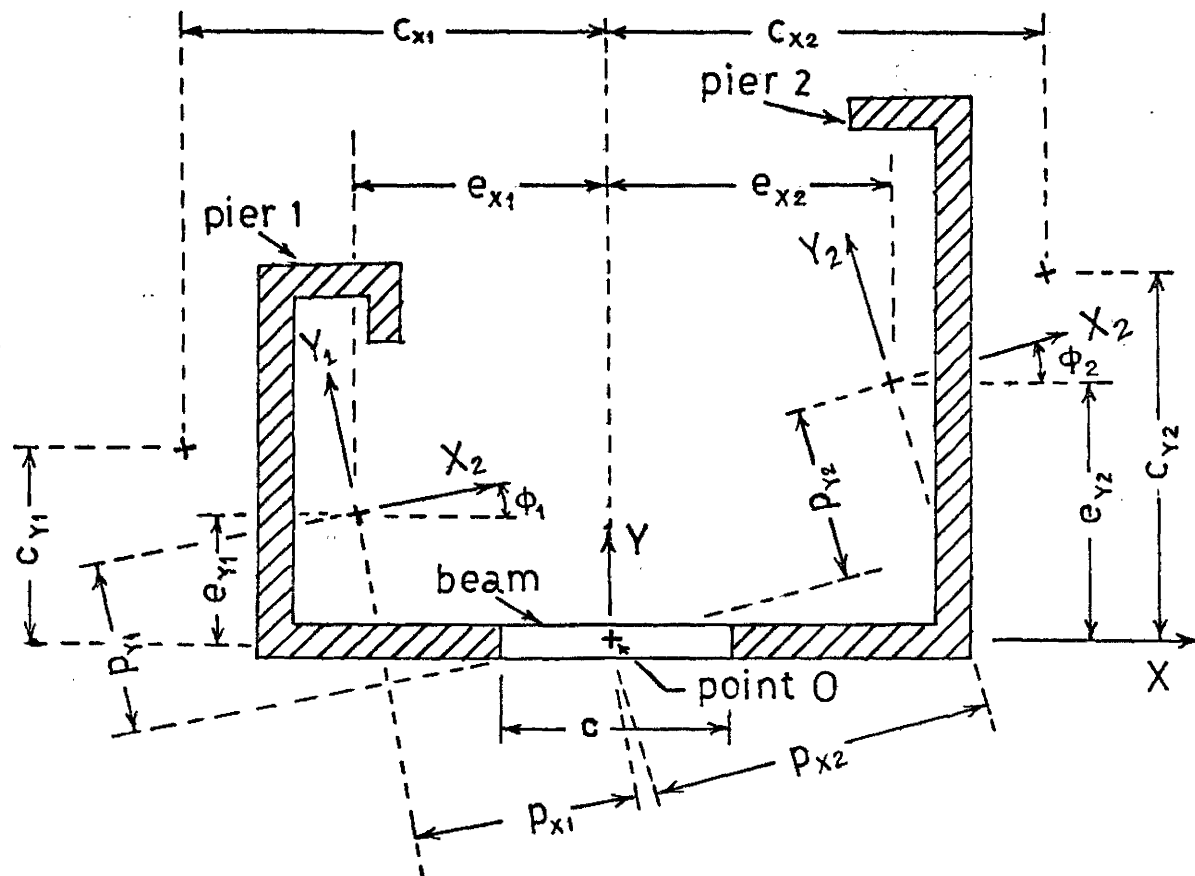
4.1 Summary

In this chapter, the nonplanar coupled shear wall is analysed using the continuum approach. Differential equations are developed for such wall subjected to lateral forces and torques distributed along the height. Special configurations of shear walls are shown and the modification on differential equations and solution is indicated. The experimental study was performed on a small scale plexiglas model consisting of two angles connected by beams at equal spacing (Fig. 4.4.1). The model was subjected to a force and a torque at the top storey and strains and deflections are measured. The experimental results are then compared with the theoretical predictions.

4.2 Theoretical Analysis

Consider two nonplanar piers which are connected by floor beams at equal spacing (Fig. 4.2.1). In the analysis the center of the connecting beam O is taken as the reference point. The differential equations are derived in terms of the displacement variables of point O, namely u , v and θ . The external forces and torques are also referred to the same point O. The theory is based on two assumptions:

(i) The deformation of the connecting beams due to bending in horizontal plane is restricted.



GEOMETRY OF COUPLED SHEAR WALL

FIG. 4.2.1

(ii) Points of contraflexure for the connecting beams due to bending in vertical plane are taken to be at the center.

In addition to the above assumptions, Vlasov's theory is taken to be valid for individual section constituting the coupled wall.

Using the continuum approach, the connecting beams are replaced by independently acting laminae of appropriate stiffness (Fig. 4.2.2a).

4.2.1 Notations Used

The notations used in the present analysis are listed below and illustrated in fig. 4.2.1.

X_1, Y_1	Orthogonal principal axes for pier 1.
u_1, v_1, θ_1	Generalised displacements of the shear center of pier 1.
X, Y, Z	Orthogonal global axes with origin at point O. X is parallel to the length of the beam.
\bar{X}_1, \bar{Y}_1	Orthogonal axes parallel to X and Y and passing through centroid of pier 1 (Fig. 4.2.4).
u, v, θ	Generalised global displacements of point O.
e_{X1}, e_{Y1}	Co-ordinates of the centroid of pier 1 referred to global axes.
c_{X1}, c_{Y1}	Co-ordinate of the shear center of pier 1 referred to global axes.
p_{X1}, p_{X2}	Distances of the centroid of pier 1 from point O measured parallel to X_1, Y_1 axes.
q_{X1}, q_{X2}	Distances of the shear center of pier 1 from point O measured parallel to X_1, Y_1 axes.
ϕ_1	Angle between X_1 and X axes.
ω_1	Sectoral ordinate of pier 1 at the point O referred to shear center of the same pier.

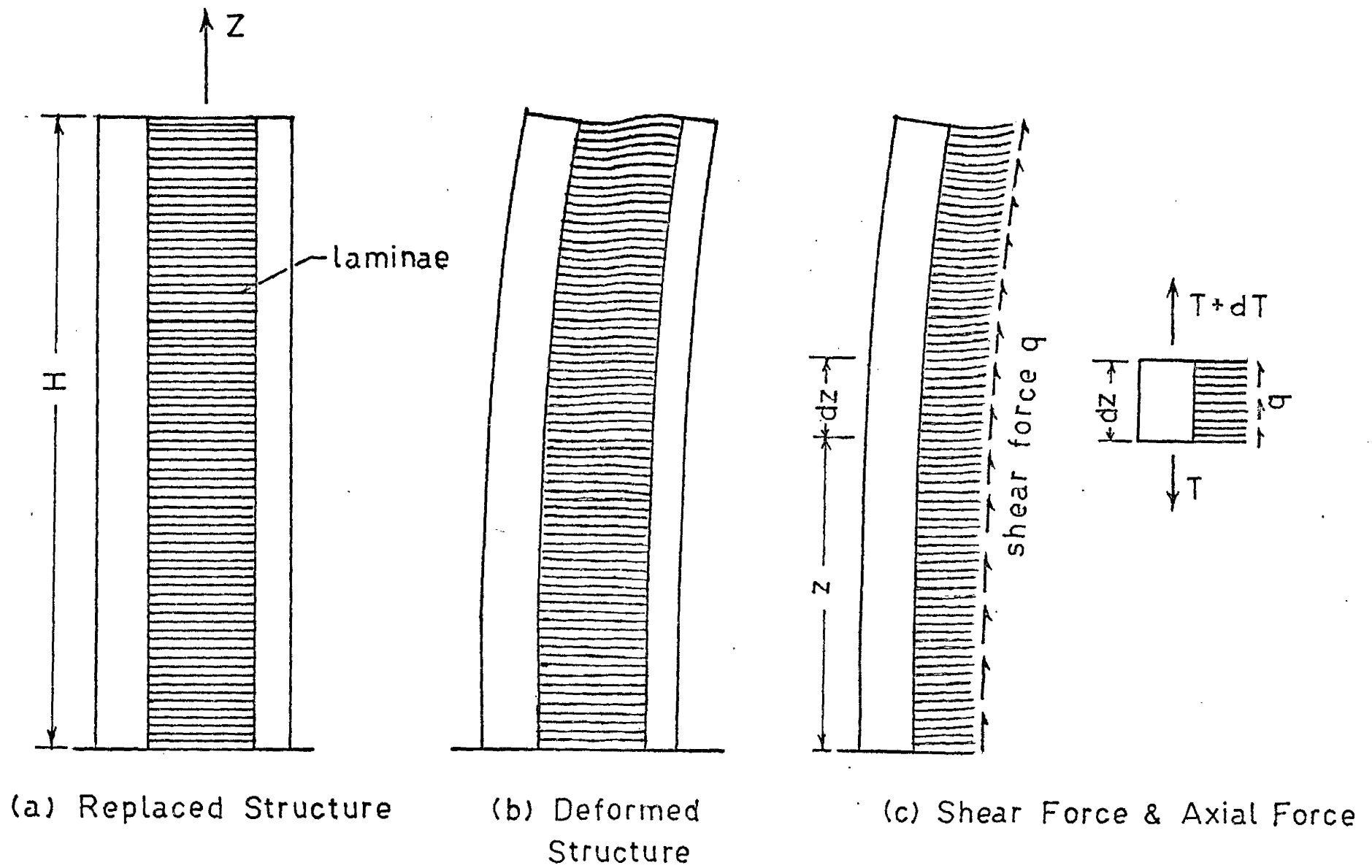


FIG. 4.2.2

I_{X1}, I_{Y1}	Moment of inertia of pier 1 about X_1, Y_1 axes.
S_{X1}, S_{Y1}	Moment of inertia of pier 1 about \bar{X}_1 and \bar{Y}_1 axes (Fig. 4.2.4).
S_{XY1}	Product moment of inertia of pier 1 about axes parallel to X,Y and passing through centroid i.e. \bar{X}_1, \bar{Y}_1 axes respectively.
$I_{\omega 1}$	Principal sectorial moment of inertia.
M_X	External moment about X axis
M_Y	External moment about Y axis
Q_t	External torque about point O.
I_b	Moment of inertia of the connecting beam
A_b	Area of the connecting beam
c	Clear span of connecting beam.
h	Storey height

Note:--The subscript 1 in the above notations are replaced by 2 for pier 2.

4.2.2. Geometric Relations

The global axes of the structure are X, Y and Z are the reference axes about which the displacements and forces are referred. The X axis is parallel to the longitudinal axis of the connecting beam. The principal axis of the piers are inclined at angles ϕ_1 and ϕ_2 to the X axis (Fig. 4.2.1). The Z axis is the vertical axis through point O.

The assumption (i) along with Vlasov's hypothesis of non deformable cross section leads to a rigid section of the coupled wall, for which the following geometric relation are valid for the transformation of displacement of the cross section from one reference point to another.

The Transfer matrices $[R_j]$ and $[T_j]$ can be defined as:

$$[R_j] = \begin{pmatrix} \cos \phi_j & \sin \phi_j & 0 \\ -\sin \phi_j & \cos \phi_j & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad (4.2.1)$$

$$[T_j] = \begin{pmatrix} 1 & 0 & -c_{yi} \\ 0 & 1 & c_{xj} \\ 0 & 0 & 1 \end{pmatrix} \quad (4.2.2)$$

$$(j = 1, 2)$$

The relation between the global displacement variables u , v and θ and the displacement variables of the piers are:

$$\begin{pmatrix} u_j \\ v_j \\ \theta_j \end{pmatrix} = [R_j][T_j] \begin{pmatrix} u \\ v \\ \theta \end{pmatrix} \quad (4.2.3)$$

$$(j = 1, 2)$$

Other geometric relations which relate the distances measured along the principal axes of the piers to the global directions are:

$$\begin{pmatrix} p_{xj} \\ p_{yi} \end{pmatrix} = \begin{pmatrix} \cos \phi_j & \sin \phi_j \\ -\sin \phi_j & \cos \phi_j \end{pmatrix} \begin{pmatrix} e_{xj} \\ e_{yj} \end{pmatrix} \quad (4.2.4)$$

$$\begin{pmatrix} q_{xj} \\ q_{yj} \end{pmatrix} = \begin{pmatrix} \cos \phi_j & \sin \phi_j \\ -\sin \phi_j & \cos \phi_j \end{pmatrix} \begin{pmatrix} c_{xj} \\ c_{yj} \end{pmatrix} \quad (4.2.5)$$

$$\text{in which } j = 1, 2$$

4.2.3 Displacement Consideration

An imaginary cut is made along the center line of the laminae. Due to deflection of the piers, there is a relative displacement δ_1 to the left and the right of the cut as shown in fig. 4.2.3(a).

$$\begin{aligned} \delta_1 = & u_2' p_{x2} - u_1' p_{x1} + v_2' p_{y2} - v_1' p_{y1} \\ & - \theta_2' I_{\omega 2} + \theta_1' I_{\omega 1} \end{aligned} \quad (4.2.6)$$

Using relations in eq. 4.2.3 and 4.2.4 this equation becomes:

$$\delta_1 = u'a + v'b + \theta'(\omega + d) \quad (4.2.7)$$

Where

$$a = e_{x2} - e_{x1}$$

$$b = e_{y2} - e_{y1}$$

$$\omega = \omega_1 - \omega_2$$

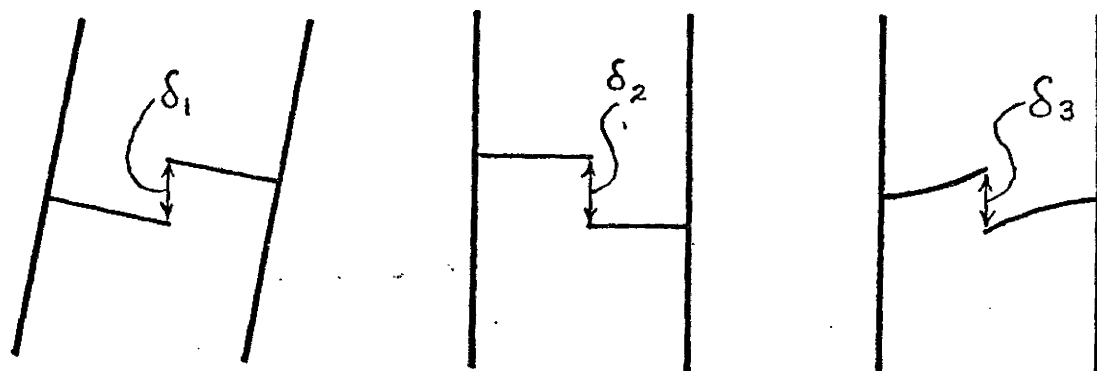
$$d = c_{x2}e_{y2} - c_{y2}e_{x2} + c_{y1}e_{x1} - c_{x1}e_{y1}$$

The shear force distribution q induced at the center of laminae produces compressive force on one pier and tensile force in the other. This axial force T in pier is related to q (Fig. 4.2.2c) as:

$$q = - \frac{dT}{dz} \quad (4.2.8)$$

This axial force T produces axial deformation of the pier which decreases the relative displacement at center of laminae by δ_2 .

$$\delta_2 = \int_0^z \frac{1}{E} \left(\frac{1}{A_1} + \frac{1}{A_2} \right) T(\xi) d\xi \quad (4.2.9)$$

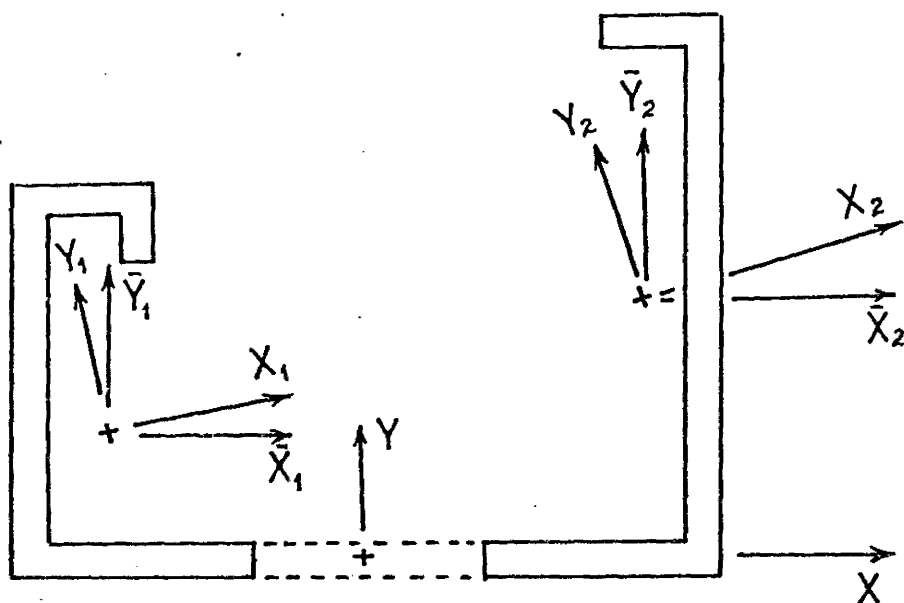


(a) Gap due to
Deflection of
Piers

(b) Gap due to
Axial Deformation
of Piers

(c) Gap due to
Deformation of
Laminas

FIG. 4.2.3



TRANSFORMATION OF MOMENT OF INERTIA

FIG. 4.2.4

Finally the force q will produce in the laminae a deformation δ_3 due to bending and shear.

$$\delta_3 = \frac{qc^3}{12EJ_b} \quad (4.2.10)$$

Where J_b is the equivalent moment of inertia of laminas taking into account shear deformation of beam.

$$J_b = \frac{I_b}{h(1 + \frac{12EI_b}{c^2GA_b})} \quad (4.2.11)$$

To satisfy the condition of compatibility it is required

$$\delta_1 = \delta_2 + \delta_3$$

Substituting above expressions for δ_1 , δ_2 , and δ_3

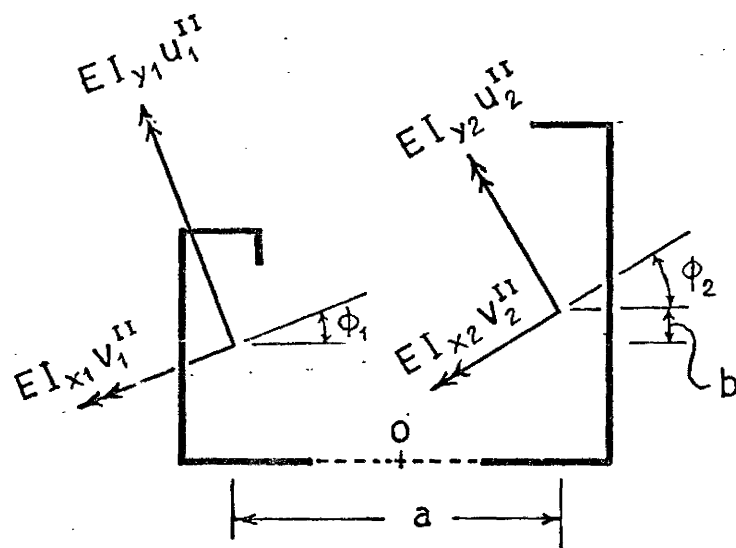
$$\begin{aligned} u'a + v'b + \theta'(\omega+d) \\ = \int_0^z \frac{1}{E} \left(\frac{1}{A_1} + \frac{1}{A_2} \right) T(\xi) d\xi + \frac{qc^3}{12EJ_b} \end{aligned} \quad (4.2.12(a))$$

Differentiating once

$$\begin{aligned} u''a + v''b + \theta''(\omega+d) \\ = \frac{1}{E} \left(\frac{1}{A_1} + \frac{1}{A_2} \right) T + \frac{c^3}{12EJ_b} \frac{dq}{dz} \end{aligned} \quad (4.2.12)$$

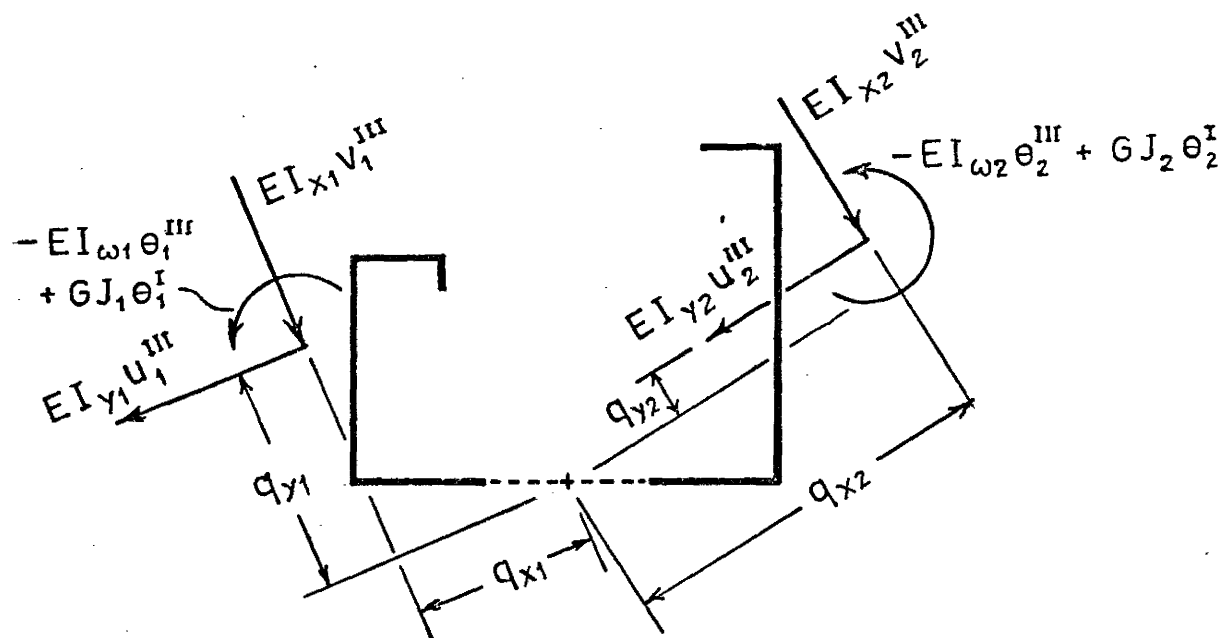
4.2.4 Force Equilibrium Conditions

The internal moments acting on the different components of coupled shear wall are shown in fig. 4.2.5. These internal moments along with couple produced by axial force T balances the external moment. For equilibrium of moments about Y and X axis, the following expressions can be derived.



INTERNAL MOMENTS

FIG. 4.2.5



INTERNAL TORQUES & SHEAR FORCES

FIG. 4.2.6

$$\begin{aligned}
EI_{y1} u_1'' \cos \phi_1 - EI_{x1} v_1'' \sin \phi_1 + EI_{y2} u_2'' \cos \phi_2 \\
- EI_{x2} v_2'' \sin \phi_2 + T_a = M_y
\end{aligned} \quad (4.2.13)$$

$$\begin{aligned}
EI_{y1} u_1'' \sin \phi_1 + EI_{x1} v_1'' \cos \phi_1 + EI_{y2} u_2'' \sin \phi_2 \\
+ EI_{x2} v_2'' \cos \phi_2 + T_b = M_x
\end{aligned} \quad (4.2.14)$$

The relations between the moment of inertia with reference to different axes (Fig. 4.2.4) are:

$$\left. \begin{aligned}
S_{xj} &= I_{yj} \sin^2 \phi_j + I_{xj} \cos^2 \phi_j \\
S_{yj} &= I_{yj} \cos^2 \phi_j + I_{xj} \sin^2 \phi_j \\
S_{xyj} &= (I_{yj} - I_{xj}) \sin \phi_j \cos \phi_j
\end{aligned} \right\} \quad (4.2.15)$$

in which $j = 1, 2$

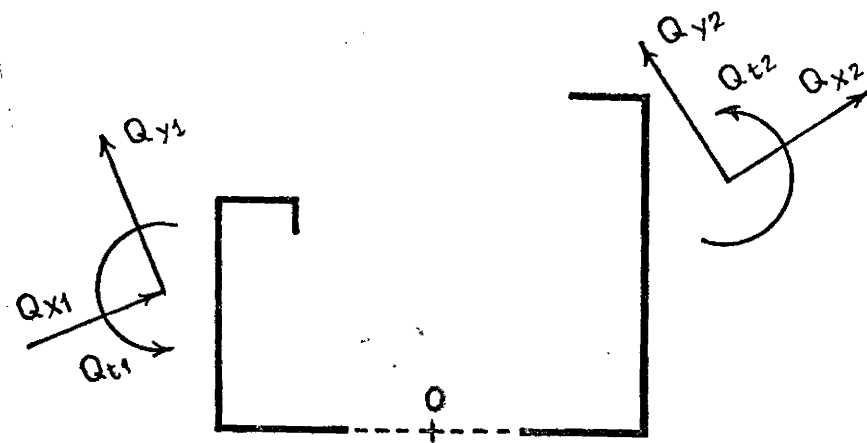
Using relations in eq. 4.2.3 and 4.2.15, the above equations are simplified as follows.

$$E S_y u'' + E_{xy} v'' - E S_{yc} \theta'' + T_a = M_y \quad (4.2.16)$$

$$E S_{xy} u'' + E S_x v'' + E S_{xc} \theta'' + T_b = M_x \quad (4.2.17)$$

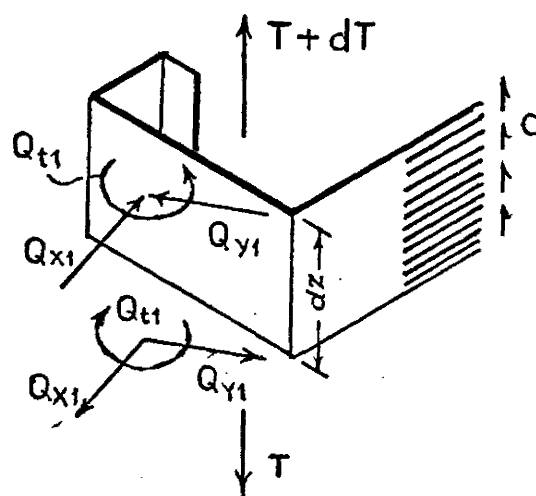
Where

$$\left. \begin{aligned}
S_y &= S_{y1} + S_{y2} \\
S_x &= S_{x1} + S_{x2} \\
S_{xy} &= S_{xy1} + S_{xy2} \\
S_{yc} &= c_{y1} S_{y1} + c_{y2} S_{y2} - c_{x1} S_{xy1} - c_{x2} S_{xy2} \\
S_{xc} &= c_{x1} S_{x1} + c_{x2} S_{x2} - c_{y1} S_{xy1} - c_{y2} S_{xy2}
\end{aligned} \right\} \quad (4.2.18)$$



TORQUES & SHEAR FORCES DUE TO
BEAM SHEAR q

FIG. 4.2.7'



FORCES ACTING ON AN ELEMENT

FIG. 4.2.8

The internal torques and shear forces acting on different components of coupled shear wall are shown in fig. 4.2.6.

Let \bar{Q}_t be the resultant torque of all these forces about point O. Therefore:

$$\begin{aligned} Q_t = & -EI_{x1}v_1''''q_{x1} + EI_{y1}u_1''''q_{y1} - EI_{x2}v_2''''q_{x2} \\ & + EI_{y2}u_2''''q_{y2} - E(I_{\omega 1} + I_{\omega 2})\theta'''' + G(J_1 + J_2)\theta' \end{aligned} \quad (4.2.19)$$

Using relations in eq. 4.2.3 and 4.2.15 the above equation is simplified as follows:

$$\bar{Q}_t = E S_{yc} u'''' - E S_{xc} v'''' - EI_{\omega} \theta'''' + GJ\theta' \quad (4.2.20)$$

Where

$$\begin{aligned} I_{\omega} = & I_{\omega 1} + I_{\omega 2} + c_{x1}^2 I_{x1} + c_{x2}^2 I_{x2} + c_{y1}^2 I_{y1} \\ & + c_{y2}^2 I_{y2} - 2c_{x1}c_{y1}I_{xy1} - 2c_{x2}c_{y2}I_{xy2} \end{aligned} \quad (4.2.21)$$

$$J = J_1 + J_2$$

Additional shear forces Q_{x1} , Q_{y1} , Q_{x2} and Q_{y2} and torques Q_{t1} and Q_{t2} develop in the section due to the shear force q in the laminae (Fig. 4.2.7). These forces can be expressed in terms of q from consideration of equilibrium of an element as shown in fig. 4.2.8. Thus,

$$\left. \begin{aligned} Q_{x1} &= -qp_{x1} \quad ; \quad Q_{y2} = -qp_{y1} \\ Q_{x2} &= qp_{x2} \quad ; \quad Q_{y2} = qp_{y2} \\ Q_{t1} &= q\omega_1 \quad ; \quad Q_{t2} = -q\omega_2 \end{aligned} \right\} \quad (4.2.22)$$

The resultant torque \bar{Q}_t about the point O then becomes

$$\begin{aligned}\bar{Q}_t = Q_{t1} + Q_{t2} - Q_{x1} q_{y1} + Q_{y1} q_{x1} - Q_{x2} q_{y2} \\ + Q_{y2} q_{x2}\end{aligned}\quad (4.2.23)$$

Using relations in eq. 4.2.4, eq. 4.2.5 and eq. 4.2.22, the above equation is simplified as follows:

$$\bar{Q}_t = q(\omega+d) = -(\omega+d) \frac{dT}{dz} \quad (4.2.24)$$

Equilibrium of torque about O gives the internal torque \bar{Q}_t together with torque due to shear force \bar{Q}_t must balance the external torque Q_t . Thus

$$\bar{Q}_t + \bar{Q}_t = Q_t$$

$$\text{or} \quad -E S_{xc} v'''' + E S_{yc} u'''' - EI_\omega \theta'''' + GJ\theta'$$

$$-(\omega+d) \frac{dT}{dz} = Q_t \quad (4.2.25)$$

4.2.5 Differential Equations

The compatibility condition (eq. 4.2.12) and the three force equilibrium condition in eq. 4.2.16, eq. 4.2.17 and eq. 4.2.25 are four equations relating the unknown of the problem u , v , θ and T . In the following paragraphs, simplification is made to reduce the four coupled equations to a single equation in θ . From eq. 4.2.16 and 4.2.17 the following expressions are obtained by algebraic elimination.

$$u'' = C_1 \theta'' + C_2 T + C_3 M_y + C_4 M_x \quad (4.2.26)$$

$$v'' = C_5 \theta'' + C_6 T + C_7 M_y + C_8 M_x \quad (4.2.27)$$

Where

$$\begin{aligned}
 SS_{xy} &= S_x S_y - S_{xy}^2; \\
 C_1 &= (S_{xy} S_{xc} + S_x S_{yc})/SS_{xy}; \\
 C_2 &= -(a S_x - b S_{xy})/E SS_{xy}; \\
 C_3 &= S_x/E SS_{xy}; \quad C_4 = -S_{xy}/E SS_{xy}; \\
 C_5 &= -(S_y S_{xc} + S_{xy} S_{yc})/SS_{xy}; \\
 C_6 &= -(b S_y - a S_{xy})/E SS_{xy}; \\
 C_7 &= -S_{xy}/E SS_{xy}; \quad C_8 = S_y/E SS_{xy}
 \end{aligned} \quad (4.2.28)$$

Differentiating eq. 4.2.26 and eq. 4.2.27 and substituting in eq. 4.2.25,

$$-q = \frac{dT}{dz} = -\frac{E\bar{I}_\omega}{r_1} \theta''' + \frac{GJ}{F_1} \theta' + C_9 M'_y + C_{10} M'_x + C_{11} Q_t \quad (4.2.29)$$

Where

$$\begin{aligned}
 \bar{I}_\omega &= I_\omega - (S_y S_{xc}^2 + S_x S_{yc}^2)/SS_{xy} \\
 r_1 &= \omega + d - S_{xc}(b S_y - a S_{xy})/SS_{xy} \\
 &\quad + S_{yc}(a S_x - b S_{xy})/SS_{xy} \\
 C_9 &= (S_{xy} S_{xc} + S_x S_{yc})/SS_{xy} r_1 \\
 C_{10} &= -(S_y S_{xc} + S_{xy} S_{yc})/SS_{xy} r_1 \\
 C_{11} &= -1/r_1.
 \end{aligned}$$

Substituting eq. 4.2.26 and eq. 4.2.27 in eq. 4.2.12,

$$Er_2 \theta''' - \frac{1}{A} T + C_{12} M_y + C_{13} M_x + \gamma \frac{d^2 T}{dz^2} = 0 \quad (4.2.30)$$

Where

$$\begin{aligned}
 r_2 &= \omega + d + a(S_{xy} S_{xc} + S_x S_{yc})/SS_{xy} \\
 &\quad - b(S_y S_{xc} + S_{xy} S_{yc})/SS_{xy}
 \end{aligned}$$

$$1/A = 1/A_1 + 1/A_2 + a(a S_x - b S_{xy})/SS_{xy} + \\ + b(b S_y - a S_{xy})/SS_{xy}$$

$$C_{12} = (a S_x - b S_{xy})/SS_{xy} \\ C_{13} = (b S_y - a S_{xy})/SS_{xy} \quad (4.2.31) \\ \gamma = c^3/12J_b$$

Differentiating eq. 4.2.31 once

$$E r_2 \theta'''' - \frac{1}{A} \frac{dT}{dZ} + C_{12} M'_y + C_{13} M'_x + \gamma \frac{d^3 T}{dZ^3} = 0 \quad (4.2.32)$$

Eliminating $\frac{dT}{dZ}$ and $\frac{d^3 T}{dZ^3}$ using eq. 4.2.29 the following fifth order differential equation in θ is obtained.

$$\beta_1 \theta^V - \beta_2 \theta'''' + \beta_3 \theta' = C_{14} M'_y + C_{15} M'_x + C_9 M_y'''' \\ + C_{10} M_x'''' + C_{16} Q_t - Q_t' \quad (4.2.33)$$

Where

$$\left. \begin{aligned} \beta_1 &= E \bar{I}_\omega \\ \beta_2 &= \frac{E \bar{I}_\omega}{A \gamma} + GJ + \frac{E r_1 r_2}{\gamma} \\ \beta_3 &= GJ/A \gamma ; C_{16} = 1/A \gamma \\ C_{14} &= \frac{C_{12} r_1}{\gamma} - \frac{C_9}{A \gamma} \\ C_{15} &= \frac{C_{13} r_1}{\gamma} - \frac{C_{10}}{\gamma} \end{aligned} \right\} \quad (4.2.34)$$

The Eq. 4.2.33 along with eq. 4.2.26, eq. 4.2.27 and eq. 4.2.29 are the final equations used in the analysis.

4.2.5 Boundary Conditions

For no rotation and displacement at base

$$\theta(0) = 0 \quad (4.2.41(a))$$

$$u(0) = 0 \quad (4.2.41(b))$$

$$v(0) = 0 \quad (4.2.41(c))$$

For no slope and warping at base

$$\theta'(0) = 0 \quad (4.2.42(a))$$

$$u'(0) = 0 \quad (4.2.42(b))$$

$$v'(0) = 0 \quad (4.2.42(c))$$

For no moment and bi moment at top

$$\theta''(H) = 0 \quad (4.2.43(a))$$

$$u''(H) = 0 \quad (4.2.43(b))$$

$$v''(H) = 0 \quad (4.2.43(c))$$

Substituting eq. 4.2.42 in eq. 4.2.12a the following condition is obtained at $Z = 0$.

$$q(0) = - \left. \frac{dT}{dZ} \right|_{Z=0} = 0$$

Using eq. 4.2.29

$$\left(\frac{EI_\omega}{r_1} \theta'''' - \frac{GJ}{r_1} \theta' - C_9 M'_y - C_{10} M'_x - C_{11} Q_t \right) \bigg|_{Z=0} = 0 \quad (4.2.44)$$

Axial force T can be expressed in terms of q by the following relation

$$T = \int_z^H q(\xi) d\xi \quad (4.2.45)$$

$$\text{From which at } Z = H, T = 0 \quad (4.2.46)$$

Substituting eq. 4.2.43 and 4.2.46 in eq. 4.2.12

$$\left. \frac{dq}{dZ} \right|_{Z=H} = - \left. \frac{d^2 T}{dZ^2} \right|_{Z=H} = 0 \quad (4.2.47)$$

Differentiating expression in eq. 4.2.29 and using in eq. 4.2.47

$$\left(\frac{EI_{\omega}}{r_1} \theta^{IV} - \frac{GJ}{r_1} \theta''' - C_9 M_Y'' - C_{10} M_X'' - C_{11} Q_t' \right) \Big|_{x=H} = 0 \quad (4.2.48)$$

The solution of the differential equation subjected to the above boundary conditions are the complete solution to the problem.

4.2.7 Solution

The solution of the differential equation (eq. 4.2.33) will consist the complimentary solution θ_C and a particular integral θ_P . Thus

$$\theta = \theta_C + \theta_P$$

The complimentary part θ_C will satisfy the following equation

$$\beta_1 \theta^V - \beta_2 \theta''' + \beta_3 \theta' = 0 \quad (4.2.50)$$

Assuming solution of the form $\theta_C = Ke^{mz}$ the following characteristic equation is obtained.

$$\beta_1 m^5 - \beta_2 m^3 + \beta_3 m = 0 \quad (4.2.51)$$

The roots are: $m_1 = 0$

$$m_2, m_3 = \pm \sqrt{\frac{\beta_2 + \sqrt{\beta_2^2 - 4\beta_1\beta_3}}{2\beta_1}}$$

$$m_4, m_5 = \pm \sqrt{\frac{\beta_2 - \sqrt{\beta_2^2 - 4\beta_1\beta_3}}{2\beta_1}}$$

In the above expressions β_1 , β_2 and β_3 are positive. It can be shown that

$$\beta_2^2 > 4\beta_1\beta_3$$

Therefore all roots are real.

Defining

$$\alpha_1 = |m_2| = |m_3|$$

$$\alpha_2 = |m_4| = |m_5|$$

The solution is of the following form.

$$\begin{aligned} \theta_c = & \kappa_1 + \kappa_2 \cosh \alpha_1 Z + \kappa_3 \sinh \alpha_1 Z + \kappa_4 \cosh \alpha_2 Z \\ & + \kappa_5 \sinh \alpha_2 Z \end{aligned} \quad (4.2.52)$$

The particular integral θ_p depends on loading.

Case-I, Concentrated load W_x , W_y and torque W_t at top. The loads are acting in the +ve directions and referred to the center line of beams. For such loading

$$M_y = W_x (H-Z)$$

$$M_x = W_y (H-Z)$$

$$Q_t = W_t$$

Particular integral for this case is:

$$\theta_p = K_p Z$$

Where

$$K_p = - (C_{14} w_x + C_{15} w_y - C_{16} w_t) \quad (4.2.53)$$

Case-II, Uniformly distributed load w_x , w_y and torque w_t acting through out the height. The loads are acting in the positive directions and referred to the center line of beams.

For such loading:

$$M_y = w_x (H-Z)^2 / 2$$

$$M_x = w_y (H-Z)^2 / 2$$

$$Q_t = w_t Z$$

Particular integral for this case is:

$$\theta_p = K_{p1} Z^2 + K_{p2} Z$$

Where

$$K_{p1} = (C_{14} w_x + C_{15} w_y - C_{16} w_t) / 2\beta_3$$

$$K_{p2} = - H(C_{14} w_x + C_{15} w_y - C_{16} w_t) / \beta_3 \quad (4.2.54)$$

are

The constants $\kappa_1, \kappa_2, \kappa_3, \kappa_4, \kappa_5$ (eq. 4.2.52) are determined from the boundary condition in eq. 4.2.41(a), eq. 4.2.42(a), eq. 4.2.43(a), eq. 4.2.44 and eq. 4.2.48.

After obtaining complete solution for θ , the expression for shear force q is determined from eq. 4.2.29. The expression for axial force T is determined from direct integration of the expression for q subjected to boundary condition as in eq. 4.2.46.

The expression of displacements u and v are determined from direct integration of the expressions in eq. 4.2.26 and eq. 4.2.27 subjected to boundary condition as in eq. 4.2.41(b), eq. 4.2.41(c), eq. 4.2.42(b) and eq. 4.2.42(c).

The displacement of the individual piers u_j , v_j and θ_j (where $j = 1, 2$) are determined from the relation in eq. 4.2.3. The average moment and bimoment of the individual piers are determined from the following relations.

$$\begin{aligned} M_{yj} &= EI_{yj} u_j'', \quad M_{xj} = EI_{xj} v_j'' \\ B_j &= -EI_{\omega j} \theta_j'' \end{aligned} \quad (4.2.55)$$

where $j = 1, 2$

4.2.8 Special Configurations

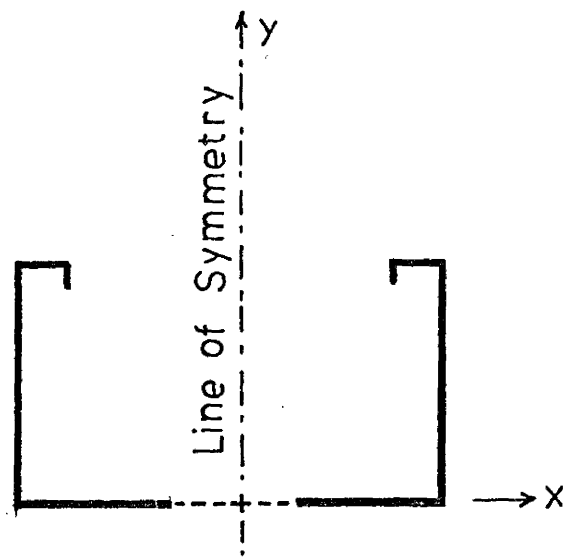
For a mono symmetric configuration of piers as shown in Case-A (fig. 4.2.11) $b = 0$; $S_{xy} = 0$; $S_{xc} = 0$. Substituting these conditions in the governing equations (eq. 4.2.12, eq. 4.2.16, eq. 4.2.17 and eq. 4.2.25) they are reduced to

$$u''a + \theta''(\omega + d) = \frac{1}{E} \left(\frac{1}{A_1} + \frac{1}{A_2} \right) T + \frac{c^3}{12J_b E} \frac{dq}{dz} \quad (4.2.56)$$

$$E S_y u'' + E S_{yc} \theta'' + Ta = M_y \quad (4.2.57)$$

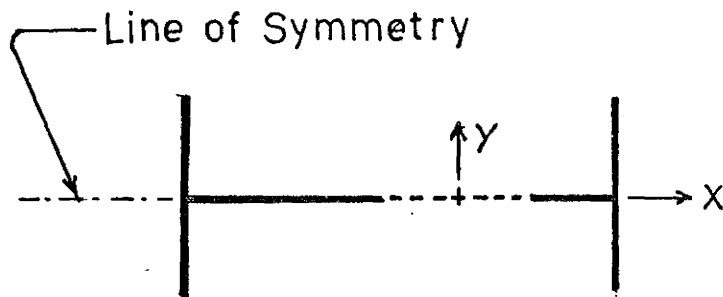
$$E S_x v'' = M_x \quad (4.2.58)$$

$$E S_{yc} u''' - EI_{\omega} \theta''' + GJ\theta' - (\omega + d) \frac{dT}{dz} = Q_t \quad (4.2.59)$$



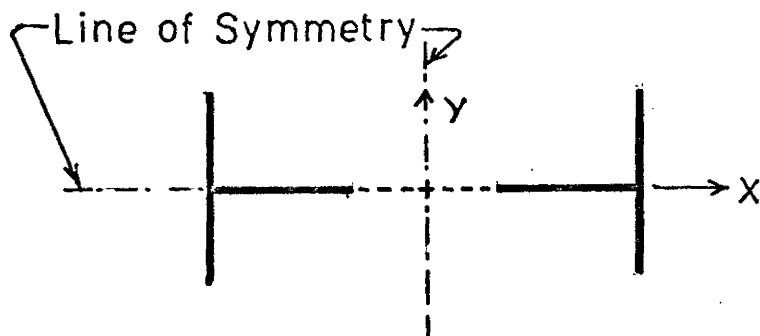
$$\begin{aligned} b &= 0 \\ S_{xy} &= 0 \\ S_{xc} &= 0 \end{aligned}$$

CASE - A



$$\begin{aligned} b &= 0 \\ S_{xy} &= 0 \\ S_{yc} &= 0 \\ \omega + d &= 0 \\ I_{\omega 1} &= I_{\omega 2} = 0 \end{aligned}$$

CASE - B



$$\begin{aligned} b &= 0 \\ S_{xy} &= 0 \\ S_{xc} &= 0 \\ S_{yc} &= 0 \\ \omega + d &= 0 \\ I_{\omega 1} &= I_{\omega 2} = 0 \end{aligned}$$

CASE - C

FIG. 4. 2.11

The eq. 4.2.58 is uncoupled and represents independent bending of the structure about the X axis. The other equations can be combined in a single differential equation in θ and solved following the same scheme as before.

For a mono symmetric configuration of piers as in Case-B (fig. 4.2.11) $b = 0$, $S_{xy} = 0$, $S_{xc} = 0$; $(\omega+d) = 0$. Substituting these conditions the governing equations (eq. 4.2.12, eq. 4.2.16, eq. 4.2.17 and eq. 4.2.25) they are reduced to.

$$u''a = \frac{1}{E} \left(\frac{1}{A_1} + \frac{1}{A_2} \right) T + \frac{c^3}{12EJ_b} \frac{dq}{dz} \quad (4.2.60)$$

$$E S_y u'' + T a = M_y \quad (4.2.61)$$

$$E S_x v''' + E S_{xc} \theta'' = M_x \quad (4.2.62)$$

$$-E S_{xc} v''' - EI_\omega \theta''' + GJ \theta' = Q_t \quad (4.2.63)$$

In the above equations u and T are coupled in the first two equations. In the last two equations v and θ are coupled but independent of u and T . The first two equations (eq. 4.2.60 and eq. 4.2.61) represent plane coupled case for bending about y axis but the bending about x axis and rotation are coupled in the other two equations.

For a symmetric configuration of piers as shown in Case-C (fig. 4.2.11) $b = 0$, $I_{xy} = 0$, $I_{xc} = 0$; $I_{yc} = 0$; $(\omega+d) = 0$. Substituting these conditions in the governing equations (eq. 4.2.12, eq. 4.2.16, eq. 4.2.17 and eq. 4.2.25),

they are reduced to:

$$u'' a = \frac{1}{E} \left(\frac{1}{A_1} + \frac{1}{A_2} \right) T + \frac{c^3}{12EJ_b} \frac{dq}{dz} \quad (4.2.64)$$

$$E S_{yc} u'' + T a = M_y \quad (4.2.65)$$

$$E S_x v'' = M_x \quad (4.2.66)$$

$$-EI_\omega \theta'''' + GJ\theta' = Q_t \quad (4.2.67)$$

The first two equations (eq. 4.2.64 and eq. 4.2.65) represent the plane coupled case for bending about y axis. The other two equations (eq. 4.2.66 and eq. 4.2.67) are uncoupled representing independent rotation and bending about x axis. It should be noted that though the individual piers do not have any sectorial moment of inertia, the group has an equivalent I_ω as defined in eq. 4.2.21. In this case

$$I_\omega = C_{x1}^2 I_{x1} + C_{x2}^2 I_{x2} \quad (4.2.68)$$

Therefore, the torsional resistance of the combined structure is substantially larger than the sum of individual resistance.

4.2.9 Effect of Neglecting Axial Deformation of Piers

Gluck (22) assumed the axial deformation of piers to be very small and neglected it in his analysis. It is of interest to indicate the present analysis is reducible to the equation given by Gluck.

If axial deformation of piers is neglected, $\delta_2 = 0$ (eq. 4.2.9). For which the compatibility equation (eq. 4.2.12a) is reduced to,

$$u' a + v' b + \theta' (\omega + d) = \frac{c^3}{12EJ_p} q \quad (4.2.69)$$

Differentiating eq. 4.2.16 and eq. 4.2.17 and using eq. 4.2.8

$$E S_y u'''' + S_{xy} v'''' - E S_{yc} \theta'''' - aq = M'_y \quad (4.2.70)$$

$$E S_{xy} u'''' + E S_x v'''' + E S_{xc} \theta'''' - bq = M'_x \quad (4.2.71)$$

Using eq. 4.2.8, eq. 4.2.25 becomes

$$\begin{aligned} -E S_{xc} v'''' + E S_{yc} u'''' - EI_\omega \theta'''' + GJ\theta' \\ + (\omega + d)q = Q_t \end{aligned} \quad (4.2.72)$$

Eliminating eq. q from eq. 4.2.70, eq. 4.2.71 and eq. 4.2.72 by the help of eq. 4.2.69 and expressing in matrix form:

$$\begin{aligned} - \begin{pmatrix} E S_y & E S_{xy} & -E S_{yc} \\ E S_{xy} & E S_x & E S_{xc} \\ -E S_{yc} & E S_{xc} & E I_\omega \end{pmatrix} \begin{pmatrix} u'''' \\ v'''' \\ \theta'''' \end{pmatrix} \\ + \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & GJ \end{pmatrix} \begin{pmatrix} u' \\ v' \\ \theta' \end{pmatrix} \end{aligned}$$

$$+ \frac{12EJ_b}{c^3} \begin{bmatrix} a^2 & ab & a(\omega+d) \\ ab & b^2 & b(\omega+d) \\ a(\omega+d) & b(\omega+d) & (\omega+d)^2 \end{bmatrix} \begin{Bmatrix} u' \\ v' \\ \theta' \end{Bmatrix} = \begin{Bmatrix} -M'_y \\ -M'_x \\ Q_t \end{Bmatrix} \quad (4.2.73)$$

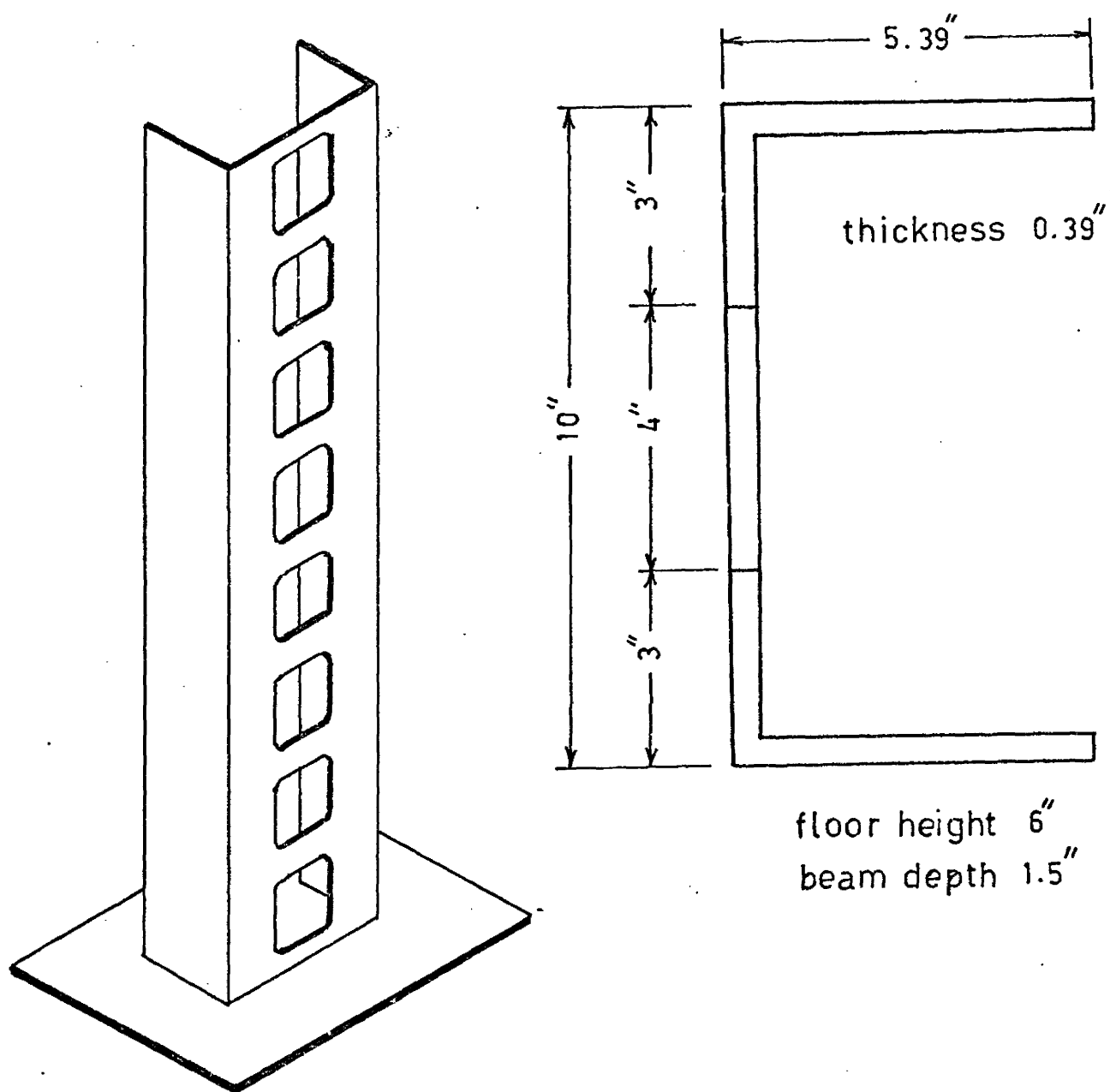
This equation is same as Gluck (22) except the factor d which is omitted. The error has been noted and suggested by Biswas and Tso (24).

4.3 Computer Program

A computer program based on the analysis has been written. This program covers the general configuration and the special configuration in Case-A (Fig. 4.2.11). subjected to concentrated force and/or torque at top. The input data are the geometric and elastic properties of the shear wall. The program determines the value of constants $\kappa_1, \kappa_2, \dots, \kappa_5$ (eq. 4.2.52) by solving a set of linear simultaneous equation obtained from the boundary conditions. The output consists of the generalised displacement of reference point O, shear force in the connecting beams and bending moment of the individual piers at chosen levels. The computer program is included in Appendix B.

4.4 Experiment

An experiment was done on a small scale model (fig. 4.4.1). It consisted of two equal angles connected by floor beams at equal spacing. The model was made from plexiglas sheet. It was loaded by a concentrated force at the top by hanging weights over a pulley system (Fig. 4.4.2). A



TEST MODEL

FIG. 4.4.1



Fig. 4.4.2
EXPERIMENTAL SET UP

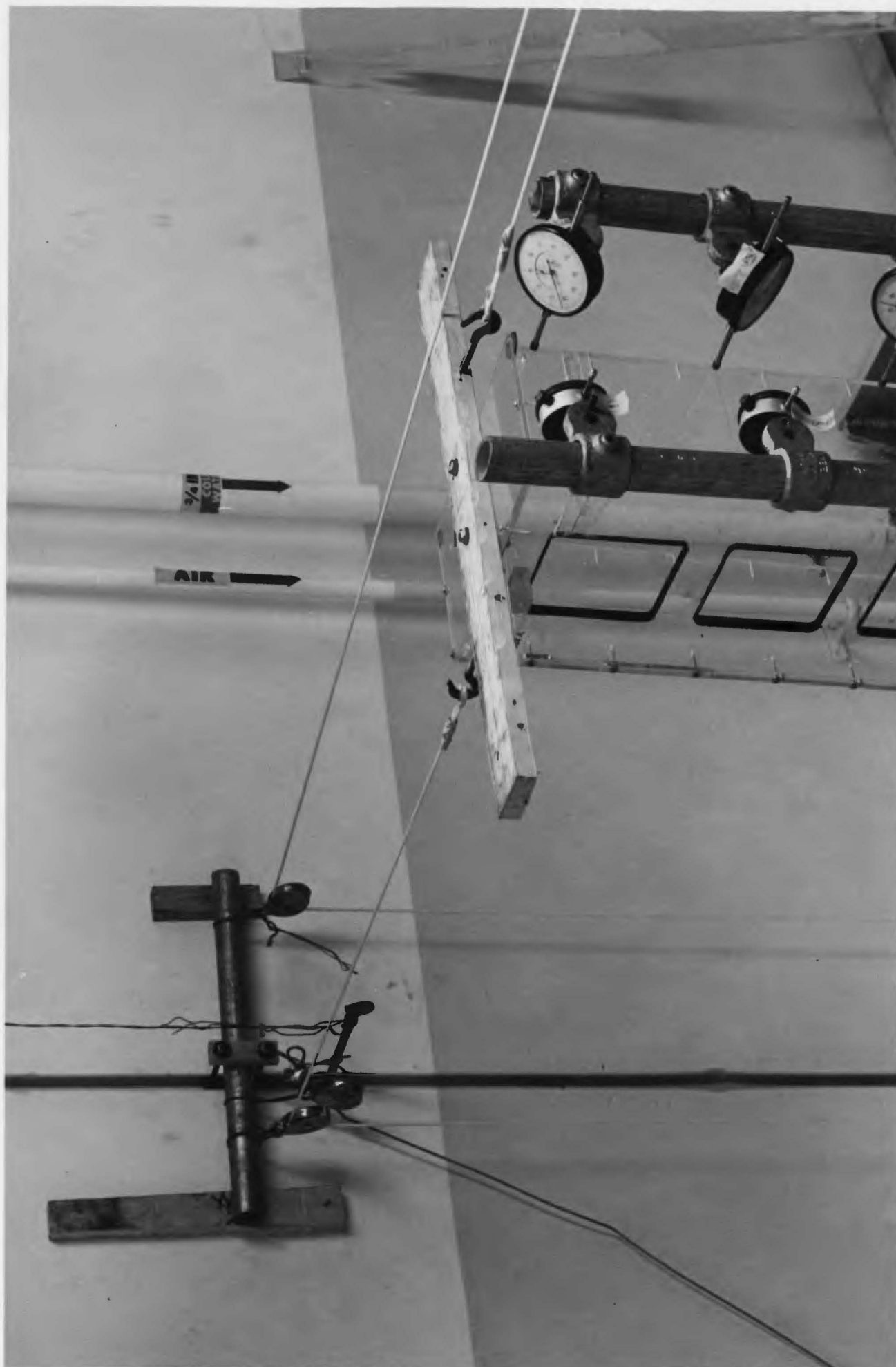
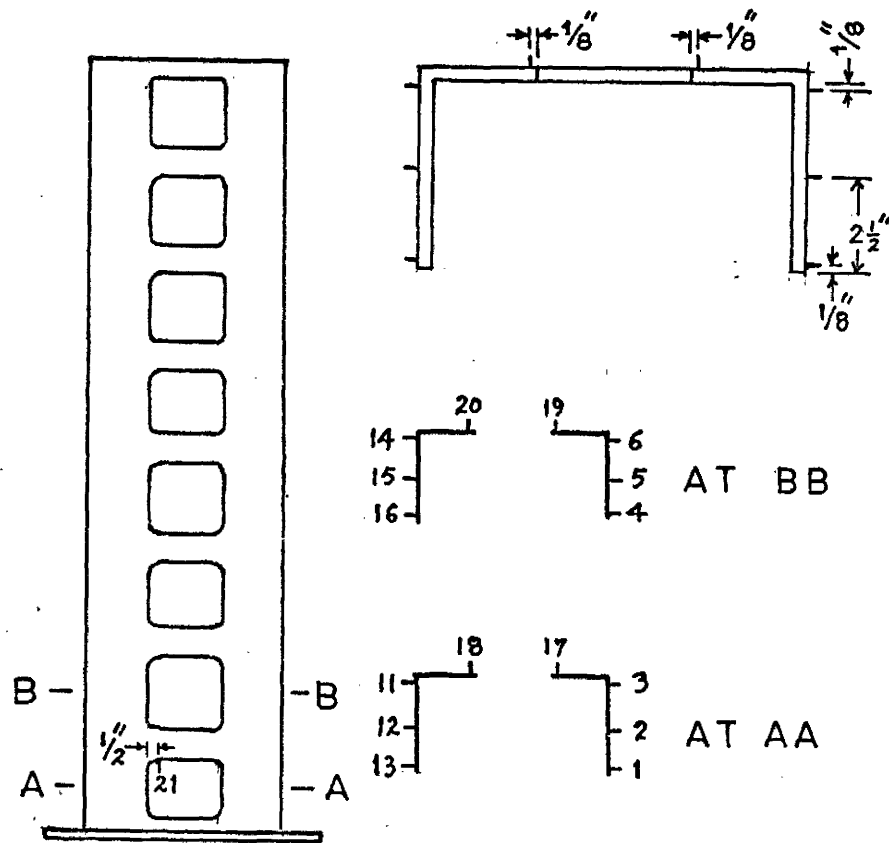
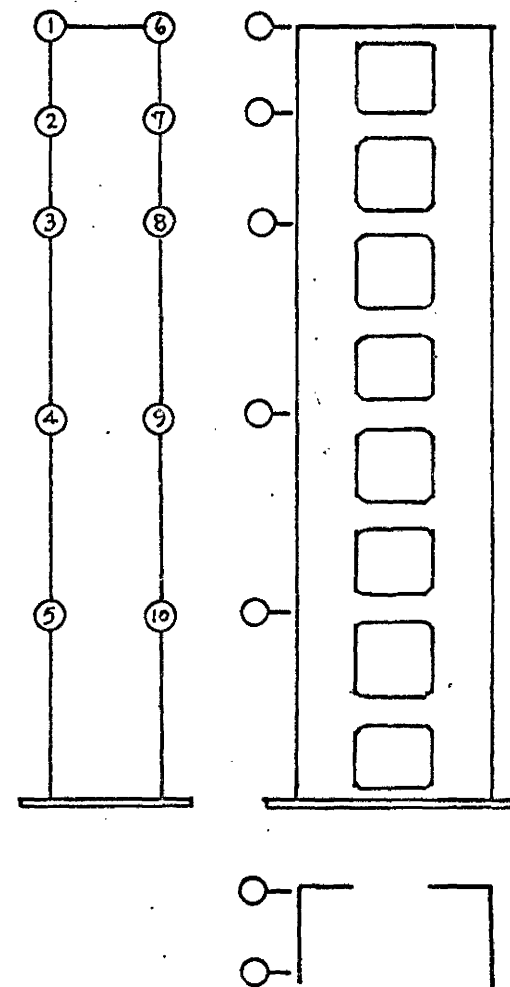


Fig. 4.4.3
TORQUE APPLYING DEVICE



(a) Strain Gauge Positions



(b) Dial Gauge Positions

FIG. 4.4.4

second loading configuration consists of a torque applied at top of the same structure. This was done by applying two equal but opposite forces as shown in fig. 4.4.3. Strain gauges and dial gauges were fixed at different points of the model (Fig. 4.4.4) and readings were taken at every increment of loading. Strain gauge and dial gauge readings are tabulated in Appendix-C. The same set of instruments used for static test of model with floor (Chapter 2) was used.

4.5 Results and Discussion

The linearity of the test structure is checked in fig.

4.8.10.. The rotation and displacement of the model subjected to concentrated load and torque at top are plotted in fig. 4.5.1 to fig. 4.5.4. The moment in pier 1 in the principal directions and distributed shear force q in the connecting beam are plotted in fig. 4.5.5. From the moment diagram, the theoretical strains at different points in level AA and BB (Fig. 4.4.4) are determined. The strain distribution thus obtained together with the experimental strains are plotted in Fig. 4.5.6 to 4.5.9.

The experimental results of rotation and displacement compared reasonably well with the theory except for the case of displacement measured due to an applied torque as shown in fig. 4.5.3. The probable reason for difference in fig. 4.5.3 may be due to the imperfection of the torque applying device. It is conceivable that some lateral load may develop in addition to the applied torque during the test.

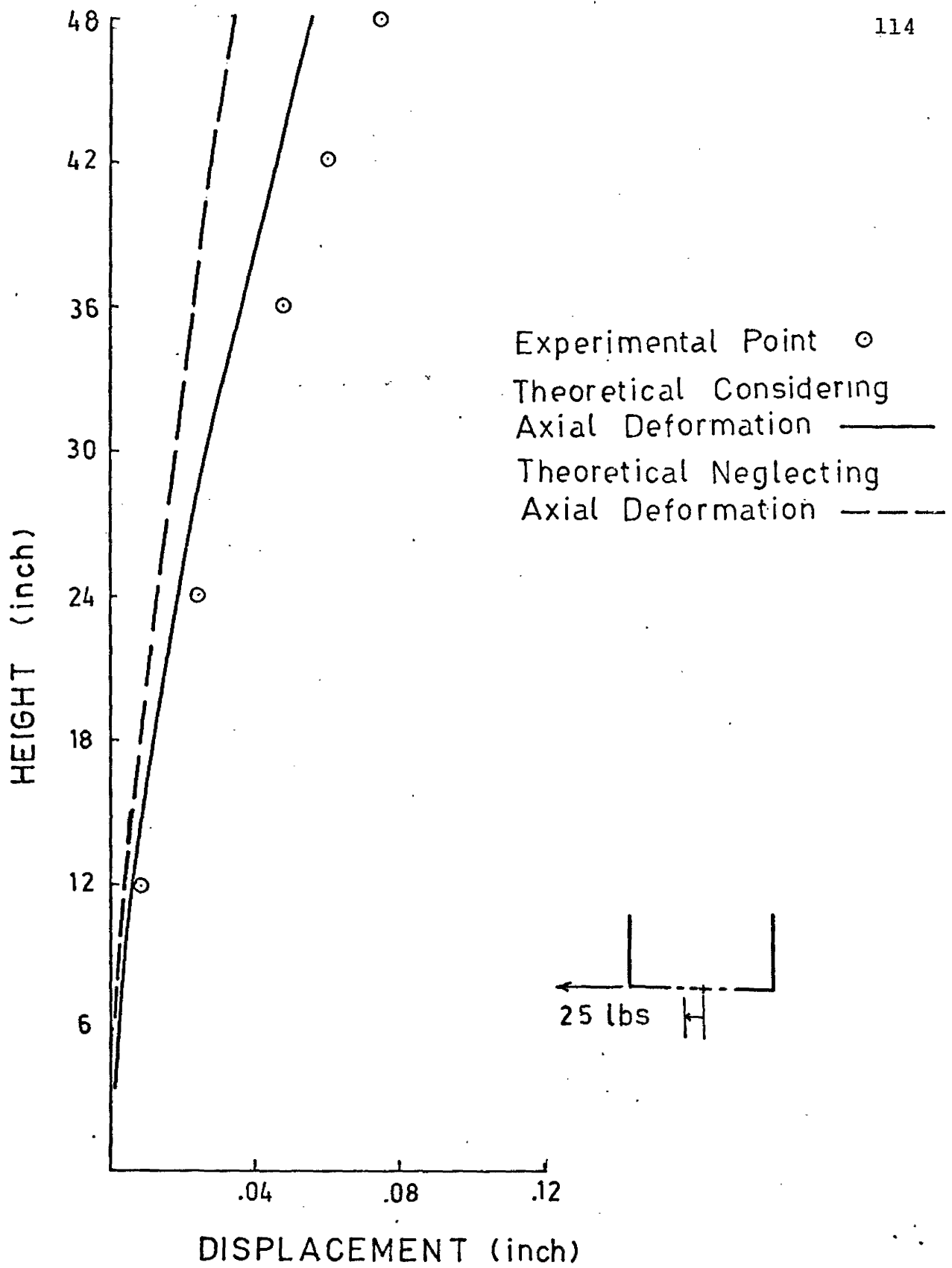


FIG. 4.5.1

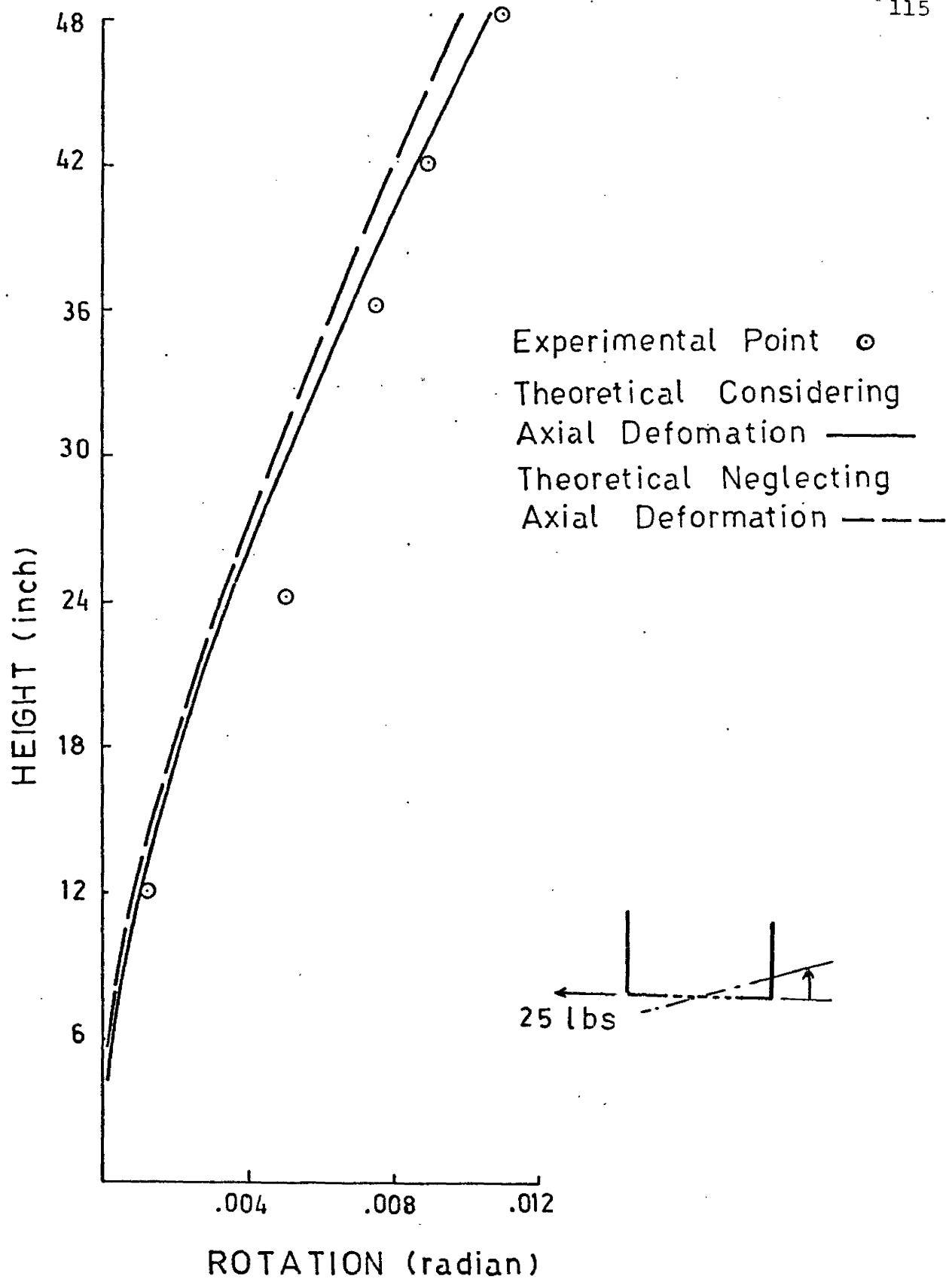


FIG. 4.5.2

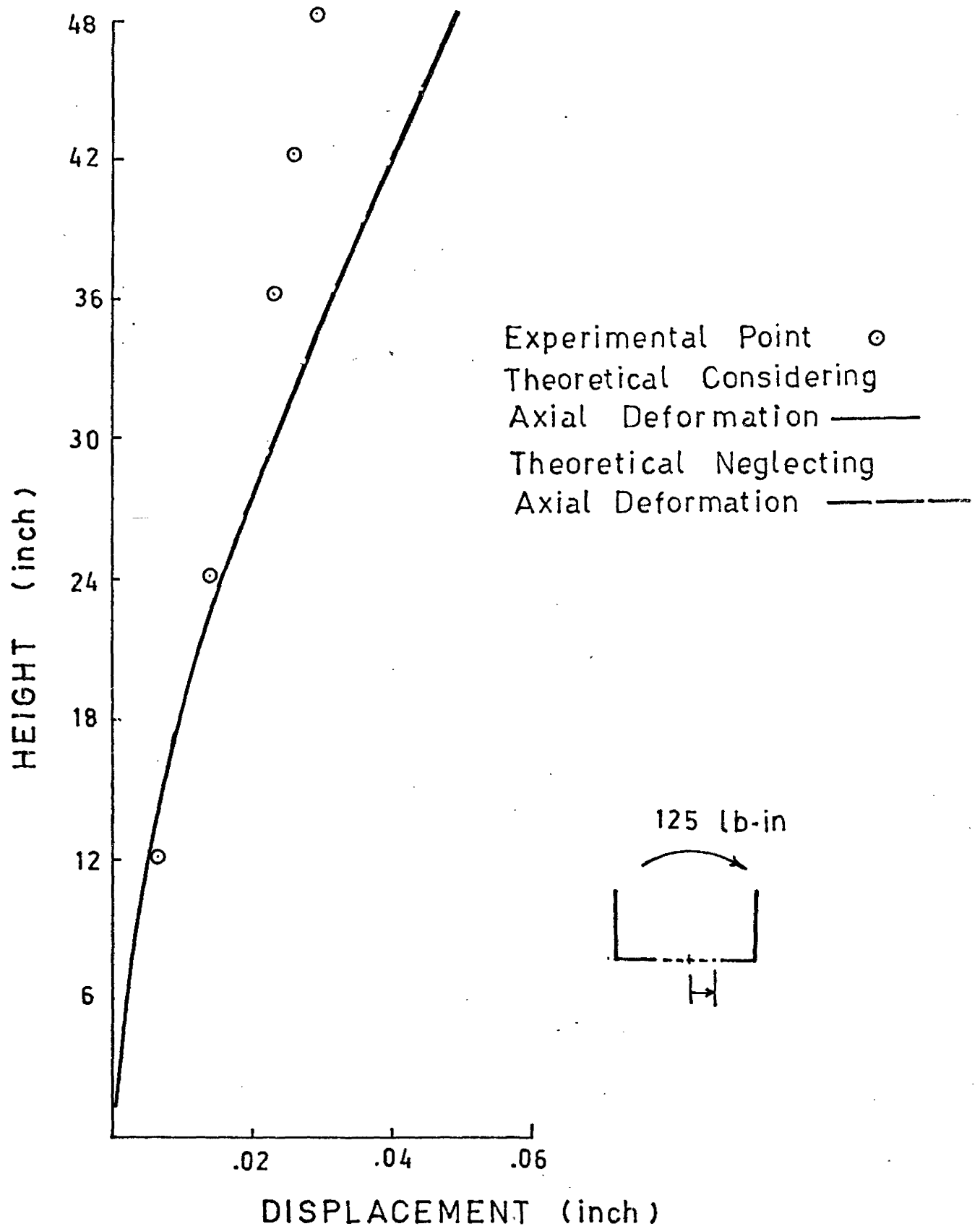


FIG. 4.5.3

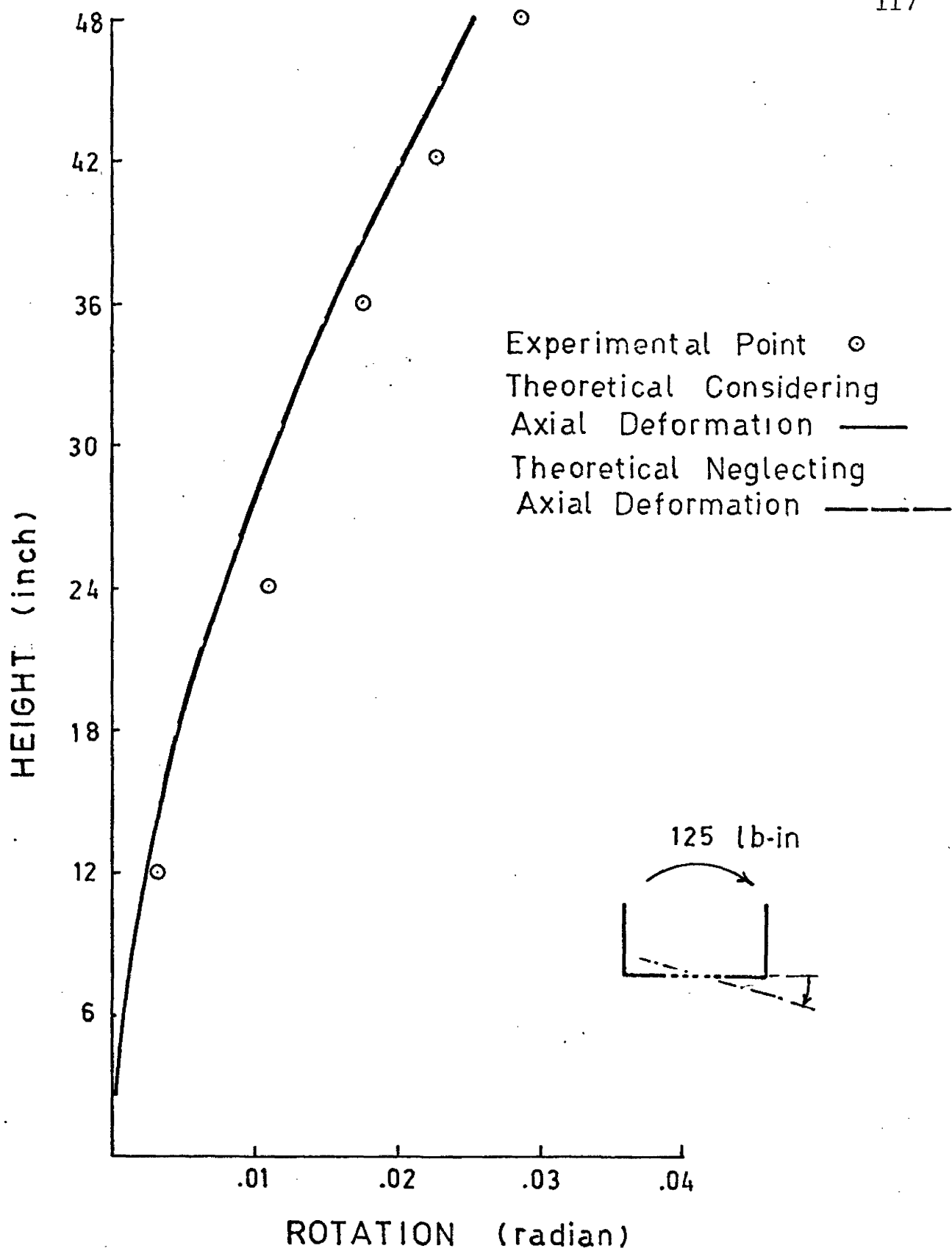
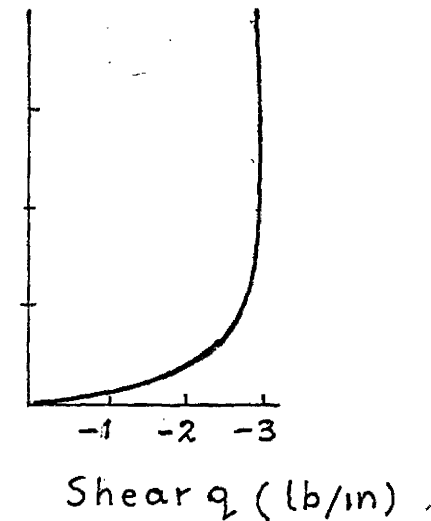
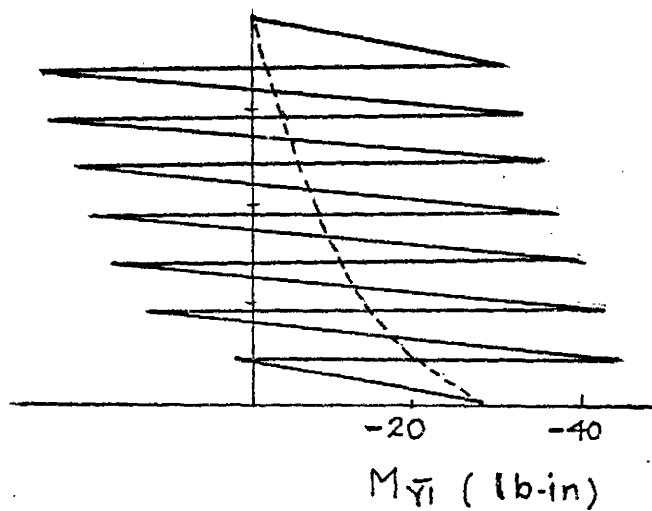
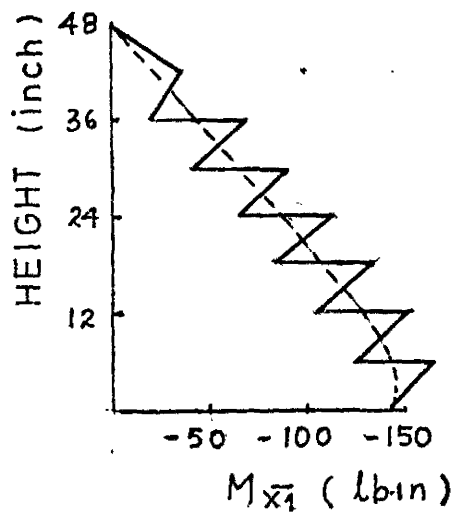
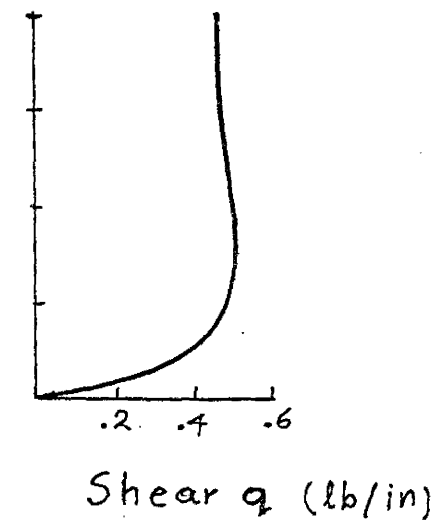
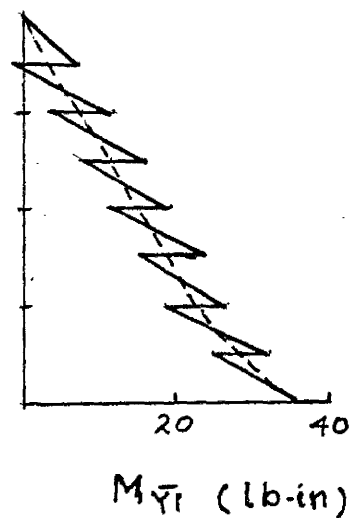
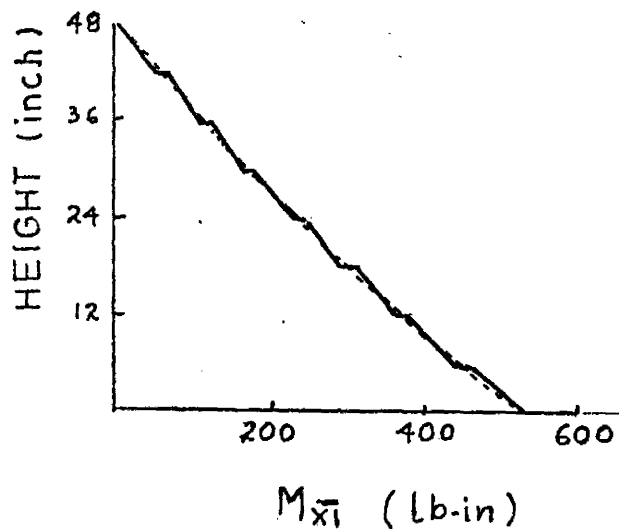


FIG . 4.5.4

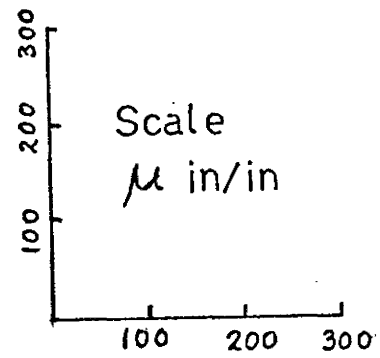
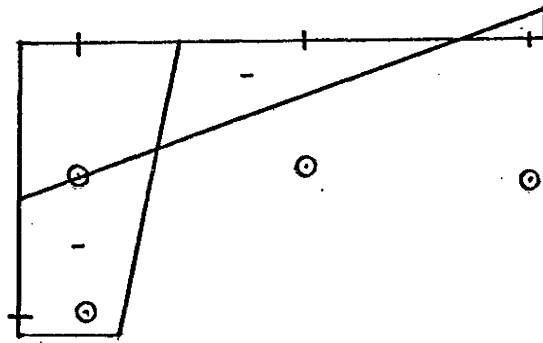


DUE TO CONCENTRATED FORCE 25 lbs AT TOP

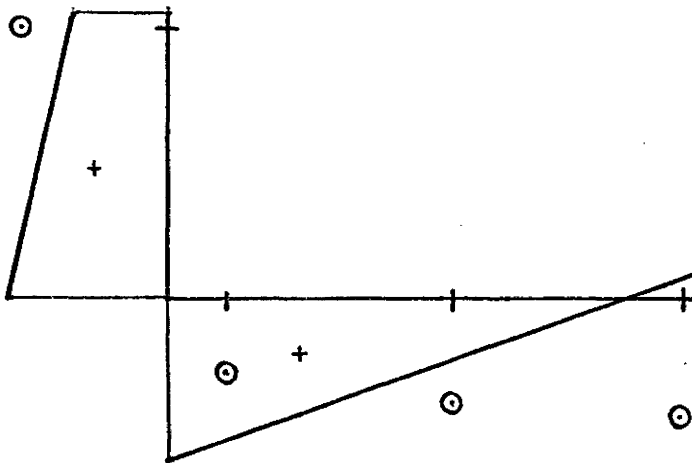


DUE TO CONCENTRATED TORQUE 125 lb-in AT TOP

FIG. 4.5.5

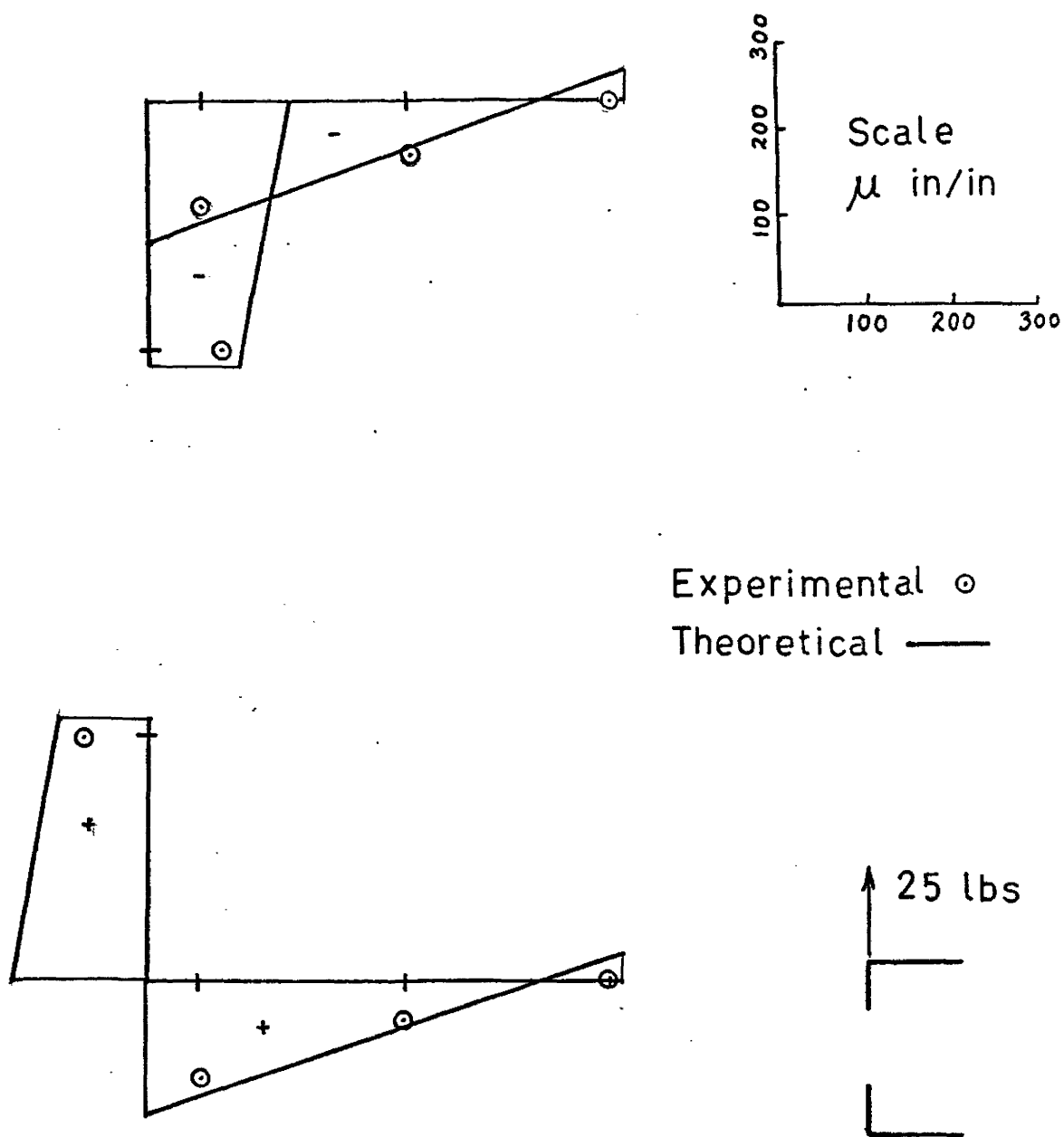


Experimental ⊙
 Theoretical —



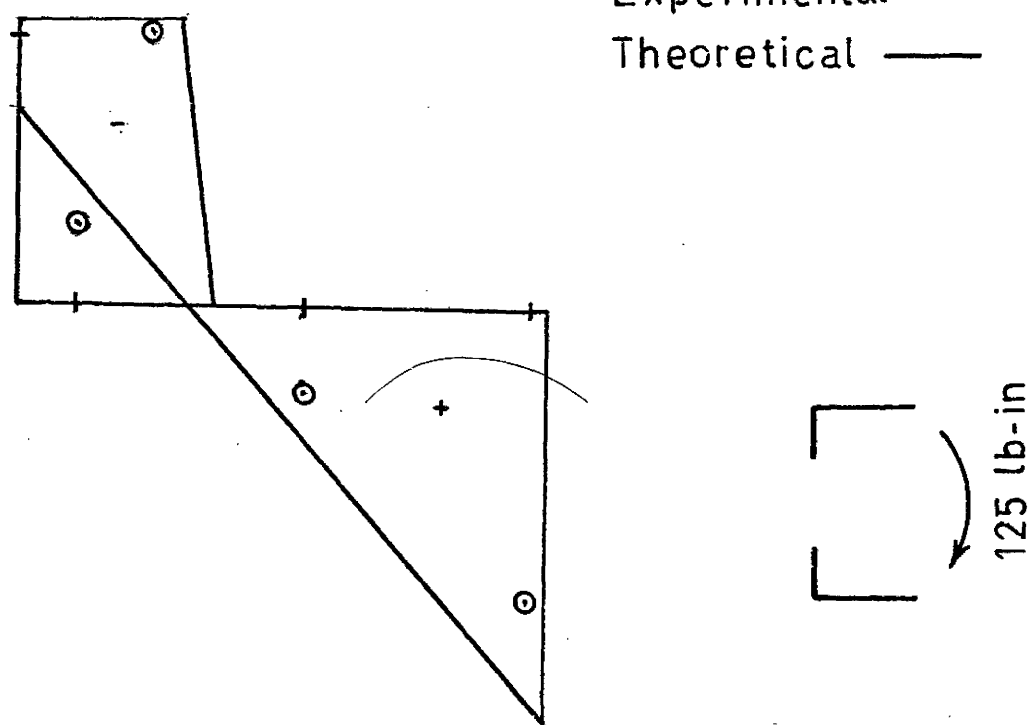
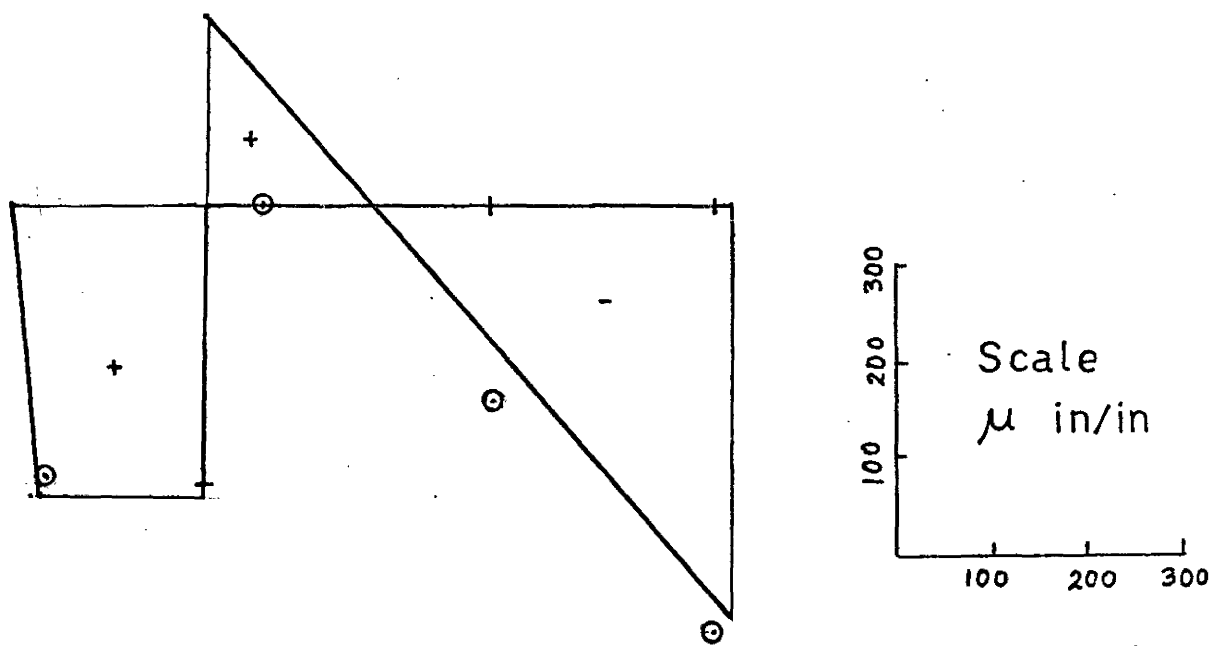
STRAIN DISTRIBUTION AT LEVEL AA

FIG. 4.5.6



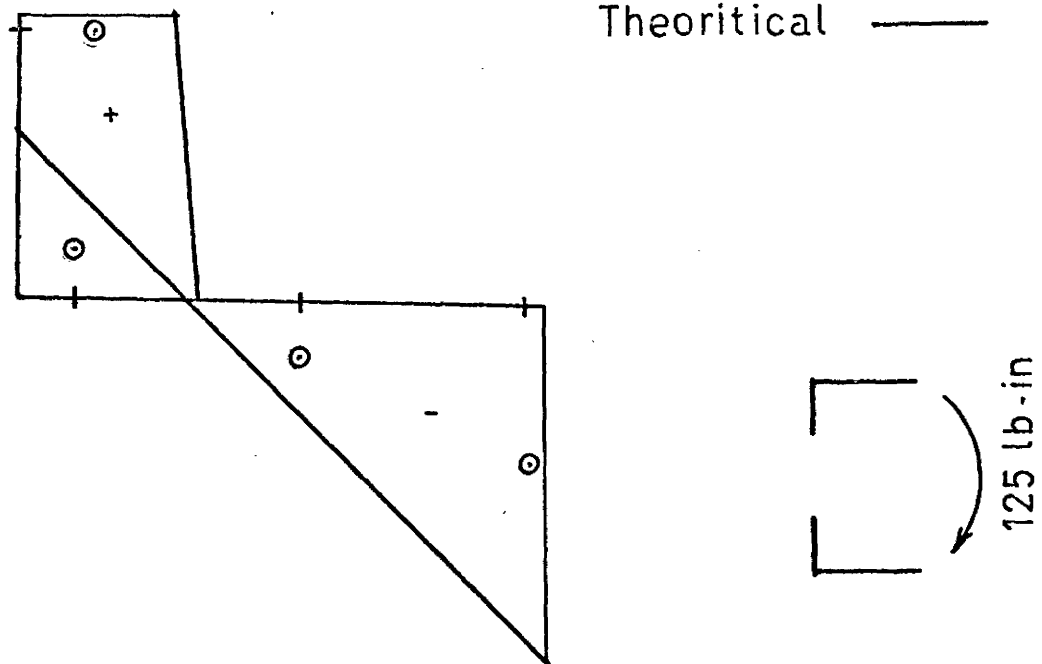
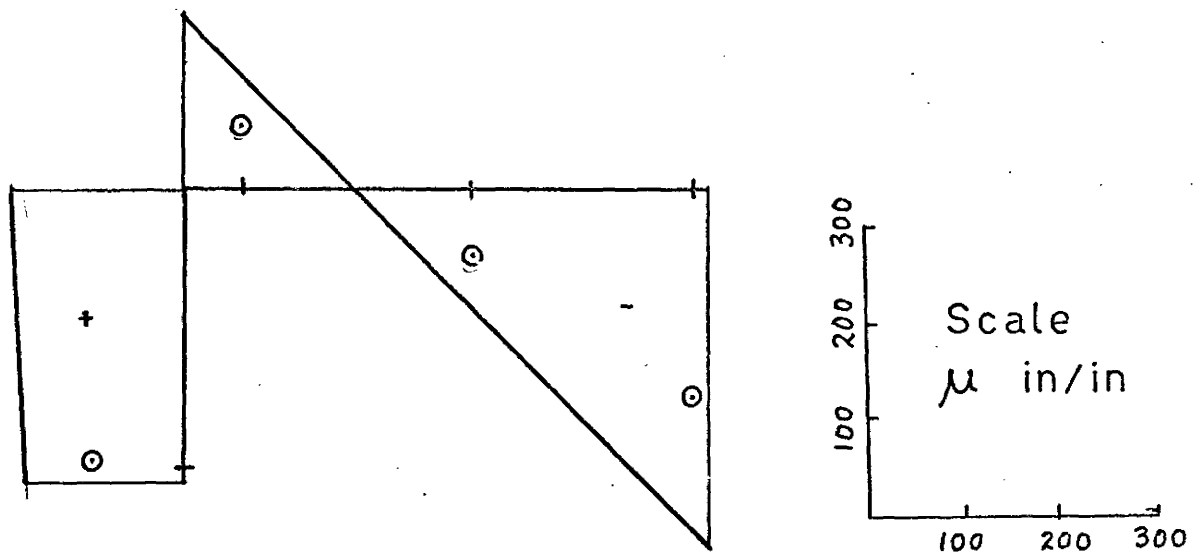
STRAIN DISTRIBUTION AT LEVEL BB

FIG. 4.5.7



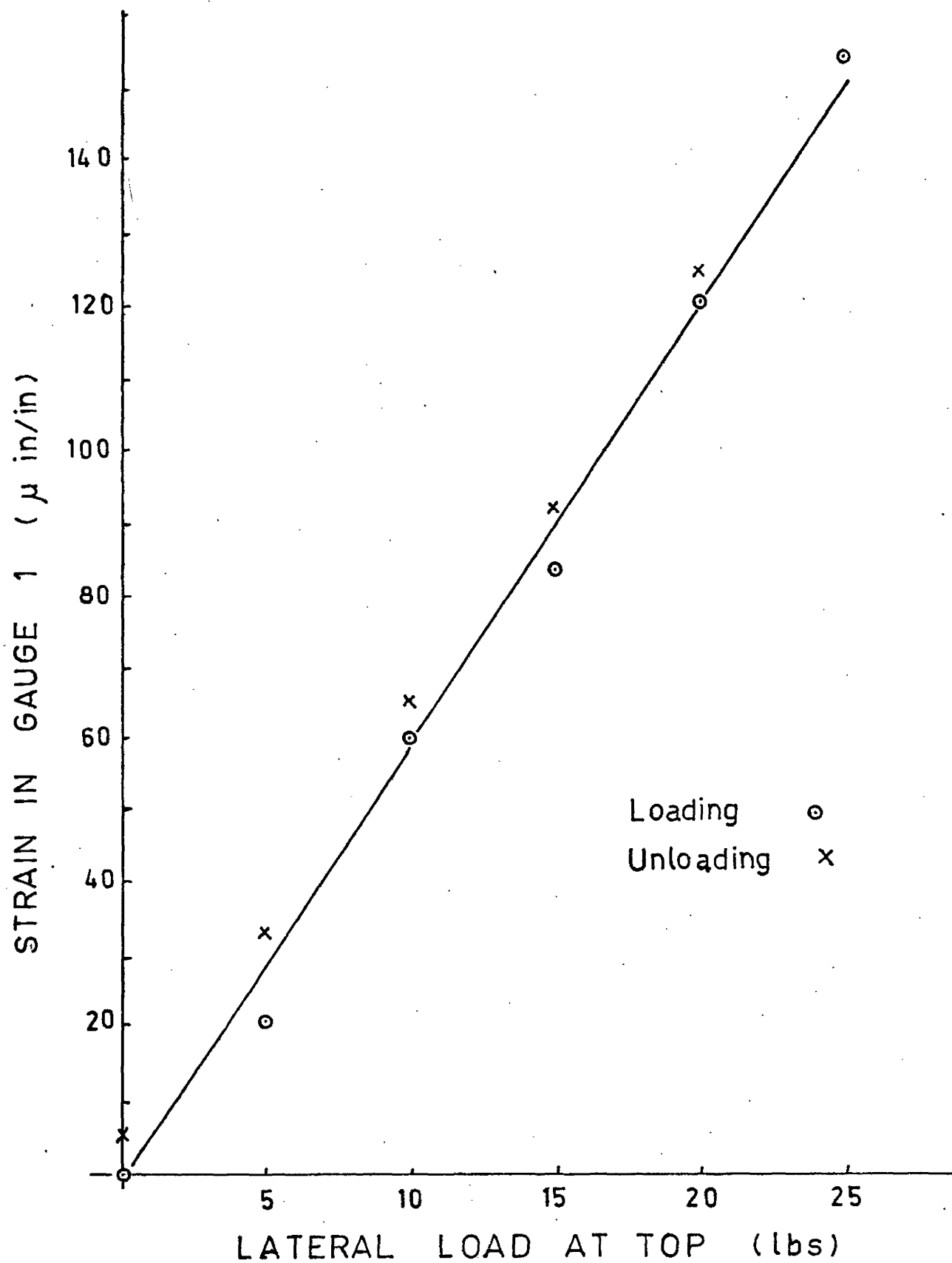
STRAIN DISTRIBUTION AT LEVEL AA

FIG. 4.5.8



STRAIN DISTRIBUTION AT LEVEL BB

FIG. 4.8.9



LOAD vs. STRAIN DIAGRAM FOR MODEL WITH BEAMS

FIG. 4.8.10

The effect of axial deformation of piers is found to have no effect on results for loading case with applied torque. However, the effect is considerable for in the case of lateral loading. The displacement at top decreases by 37% and rotation at top decreases by 7% as a result of neglecting axial deformation. The error introduced by neglecting axial deformation becomes significant if the axial force in the piers is large as in the case of the lateral loading. A comparison from the shear (q) plot in fig. 4.5.5 shows that the shear is about five times in the case of applied loading as compared to the case of applied torque. Since the axial force in the piers is the sum of the distributed shear q , neglecting the deformation due to axial force in the case of applied loading affect the results considerably.

In the strain diagrams (Fig. 4.5.6 to 4.5.9) the comparison between the experiment and the theory is less favourable. Since strain is a local measure, the experimental results are affected by the local imperfection of the test model. The continuum approach used in the analysis is expected to give results to the overall behaviour of the structure but with less accurate results to the local behaviour. Nevertheless, the trend of the strain distribution as predicted by the theory is varified by the experimental points.

In general it is expected that the continuum approach of analysis is best suited for structures with a large number of stories. In the present study, the model used in the experiment consists of eight floor beams only. Yet, the results obtained from the analysis compare well with the

experimental points.

CHAPTER V

CONCLUSIONS AND RECOMMENDATIONS

5.1 Conclusions

The results from the static analysis and the testing of shear wall model with floors show that floor slab can provide considerable torsional stiffness by providing restraint against warping. Neglect of the floor slab stiffness will underestimate the stiffness of the structure. Moreover the actual structure is not as stiff as predicted by considering both torsional and flexural stiffness of the slabs. The bending stiffness of the slab is less effective due to the local rotation of the joints and a modified bending stiffness is to be considered. This is taken into account by a factor K in eq. 2.6.2. Theoretical values thus obtained compare reasonably well with experiments for displacement and strain measurements.

In the dynamic study of shear wall with floor slabs, the mass and mass moment of inertia of the slabs have to be considered in addition to the stiffness provided. The vibration is in general coupled. The first mode is a torsion predominant one and the second mode is bending predominant. The theoretical frequency compared reasonably with the experimental values except for the second frequency which is 16% higher. This is due to the neglect of shear deformation in the theory, which have considerable effect on the frequency for bending predominant modes.

In the static and dynamic analysis of shear walls with floors, the 'Matrix Transfer' method has been used. The method is ideal for continuous systems with discrete points. The main advantage is that the size of the matrix handled in the analysis is independent of the number of floors. An increase of the number of floors will only increase the number of matrix multiplication which can be done easily in a digital computer.

The non-planar coupled shear wall was analysed using the continuum method. It is shown that the effect of neglecting axial deformation of piers will lead to gross overestimation of stiffness of the structure for certain cases. Simplicity can be achieved by this assumption but should always be done with caution. The theoretical analysis shows reasonable agreement with the experimental results for displacements but less favourable for strains. So it can be concluded that continuum method gives good results for overall behaviour but is less accurate for local behaviour of a structure.

5.2 Recommendations

The present study of shear wall structures with special interest on floor slabs and floor beams are carried out on simple structures. The arrangement of the shear walls in an actual multi-storey building is very complex. The analysis is also complex due to various interacting elements. Simplicity can be achieved by assumption but a complete understanding of the behaviour of different interacting elements

is necessary for making reasonable assumptions.

More experimental and theoretical study on the behaviour of floor slab with different geometrical arrangement is necessary. In such cases, it may not be justified to treat the floor slab as a series of beams but as an elastically supported plate.

Extension of the formulation on non-planar coupled shear wall is necessary for more generalised cases. The present analysis is applicable for two shear walls connected by a single row of beams. Extension for cases with more piers and more rows of connecting beam is required. A study on the dynamic analysis for non-planar coupled shear walls is also recommended.

In general, the analytical tools presently available to design engineers are very limited and most of the time very restricted in its applications. Therefore, more theoretical work supported by experimental data is necessary in the shear wall field to assist the practicing engineers to analyse and design structures which will be safe at the same time economic.

APPENDIX A

VLASOV'S THEORY OF THIN WALLED BEAM

The method of analysis used in the present work is based on the theory presented by Vlasov (8). Vlasov's theory is based on two geometric hypothesis:

(a) a thin walled beam of open section can be considered as a shell of rigid (undeformable) cross section.

(b) the shear deformation of the middle surface (characterising change in the angle between the co-ordinate line) can be neglected.

In shear wall structure, the concrete clear wall can be treated as thin walled beams connected by floor slabs which are normally located at regular intervals. The action of the floor slab is to prevent any deformation of the section which supports the hypothesis (a). Hypothesis (b) requires shear deformation to be negligible compared with the torsional and flexural deformations. Vlasov states that this is satisfied if for the structure shown in Fig. A.1

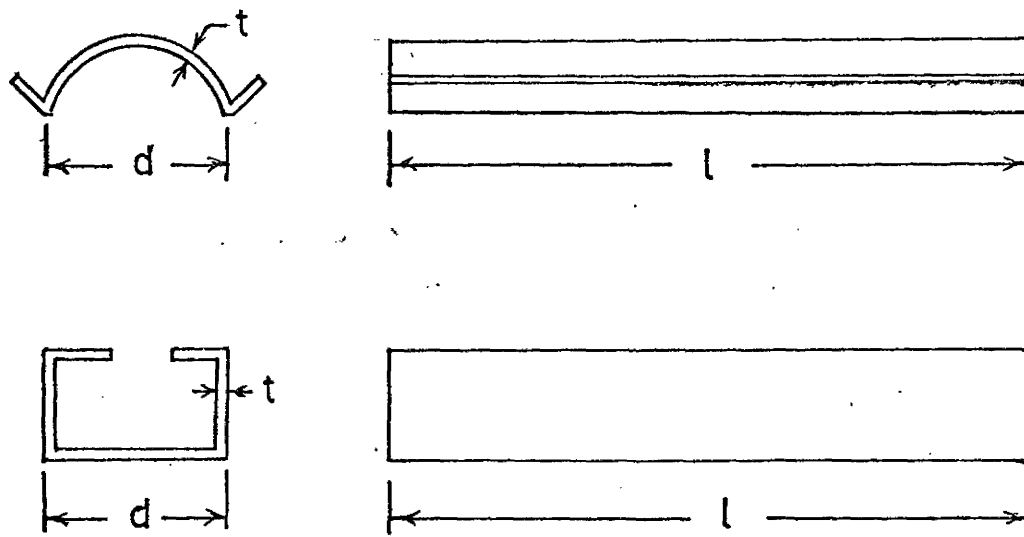
$$t/d \leq 0.1 \text{ and } d/l \leq 0.1$$

For components of tall building this conditions are satisfied.

The expression of longitudinal stress in Vlasov's theory is

$$\sigma = \frac{N}{A} - \frac{M_y \cdot x}{I_y} - \frac{M_x \cdot y}{I_x} + \frac{B \cdot \omega}{I_\omega} \quad (\text{A.1})$$

The first three terms coincide with the equation known from elementary theory. The last term of the expression is



Dimensional Limitations
FIG. A 1

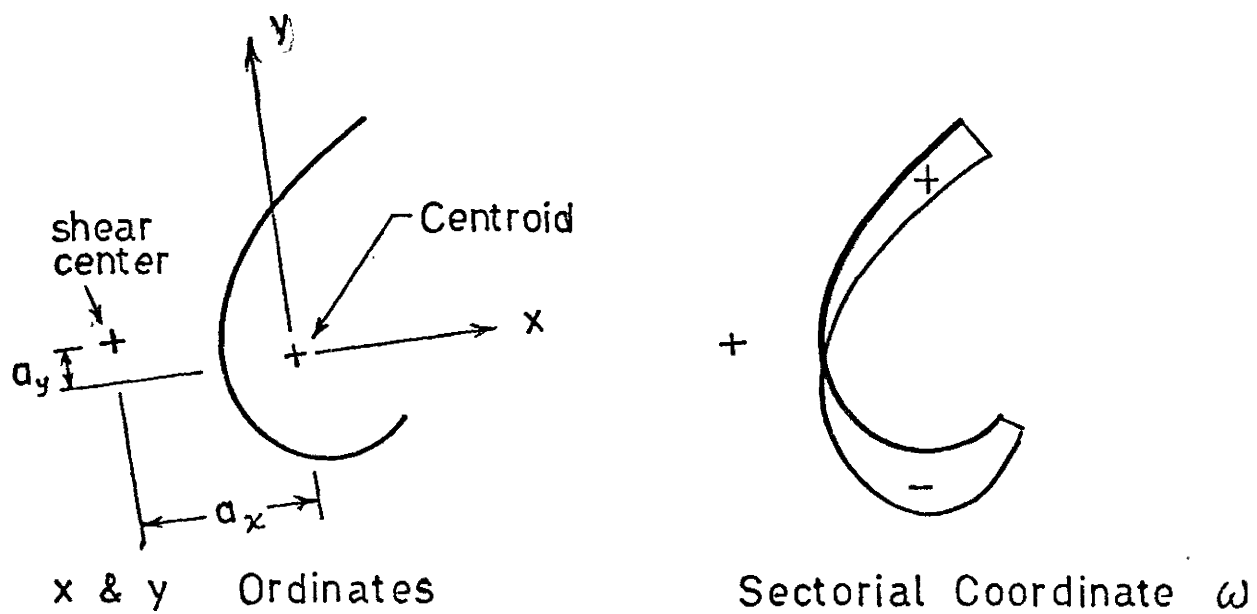
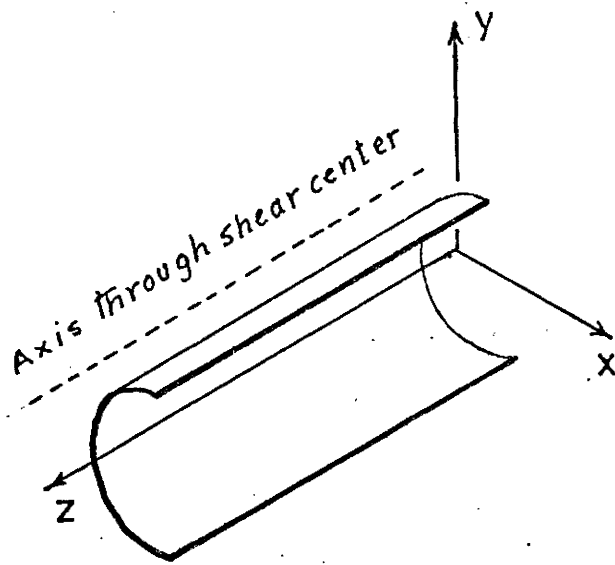
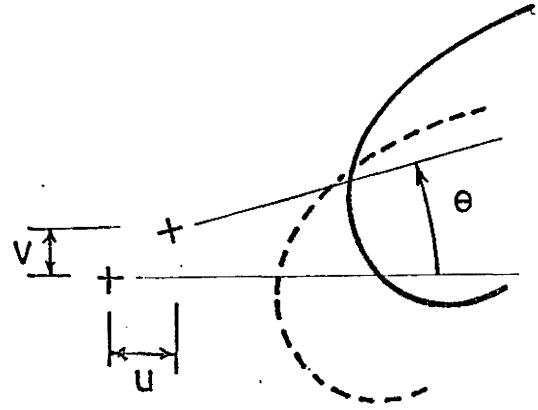


FIG. A 2



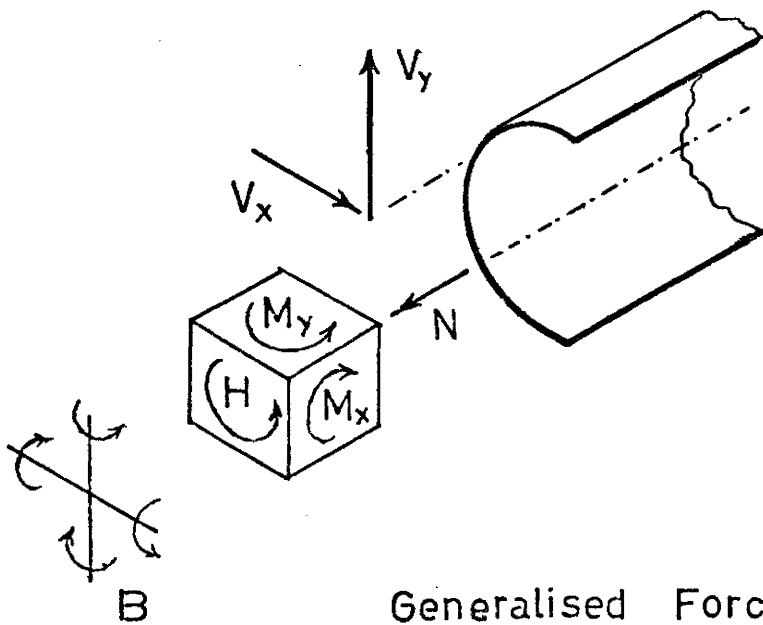
Co-ordinate System

FIG. A 3



Generalised Displacements

FIG. A 4



Generalised Forces

FIG. A 5

the longitudinal stress due to warping. The notations of the last terms are explained as:

Bimoment B is a generalised balanced system of forces statically equivalent to zero. Units are force \times (length)² e.g. lbs. in².

ω is the sectorial area of the point on the section where the stress is being measured. Units are (length)² e.g. in².

I_ω is sectorial moment of Inertia of the section and is defined as $\int_A \omega^2 dA$ units are (length)⁶ e.g. in⁶.

The distribution of sectorial co-ordinate for some open section is shown in Fig. A.2. Right hand co-ordinate system used in the present work (Fig. A.3). The generalised displacement variables are shear center displacements u, v and θ in xy plane and centroidal displacement s in z direction (Fig. A.4). Sign convention for generalised forces are shown in Fig. A.5.

The relation between the generalised forces and displacement variables are:

$$\left. \begin{aligned} N &= EAs', \quad M_x = EI_x v'', \quad M_y = EI_y u'', \\ B &= -EI_\omega \theta'', \quad H = -EI_\omega \theta'''' + GJ\theta', \\ V_y &= -EI_x v''', \quad V_x = -EI_y u'''. \end{aligned} \right\} \quad (A.2)$$

Of these quantities, axial force and bending moments are referred to the centroid and shear forces and torque are referred to the shear center of the section.

The longitudinal stress at any point can be expressed

in terms of displacement variables as:

$$\sigma = E[s' - u''x - v''y - \theta''\omega] \quad (A.3)$$

Displacement of any point in z direction is obtained as:

$$\delta = s - u'x - v'y - \theta'\omega \quad (A.3a)$$

Tensile stress and displacements in the direction of positive x, y & z are positive in eq. A.3 and eq. A.3a. When referred to principal generalised co-ordinates of an open section, the differential equation of a thin walled beam statically loaded at its ends are:

$$\left. \begin{aligned} EAs'' &= 0 \\ EI_x v^{IV} &= 0 \\ EI_y u^{IV} &= 0 \\ EI_\omega \theta^{IV} - GJ\theta'' &= 0 \end{aligned} \right\} \quad (A.4)$$

Differential equations for free vibration of a thin walled beam are:

$$\left. \begin{aligned} EAs'' - \rho As'' &= 0 \\ EI_x v^{IV} + \rho A v'' - \rho I_x v'''' - \rho A a_x \theta'' &= 0 \\ EI_y u^{IV} + \rho A u'' - \rho I_y u'''' + \rho A a_y \theta'' &= 0 \\ EI_\omega \theta^{IV} - GJ\theta'' + \rho I_p \theta'' - \rho I_\omega \theta'''' + \rho A a_y u'' - \rho A a_x v'' &= 0 \end{aligned} \right\} \quad (A.5)$$

Eq. A.4 and eq. A.5 are referred in Chapter 2 and Chapter 3 and solved accordingly as they appeared in the analysis.

APPENDIX - B

COMPUTER PROGRAMS

PROGRAM TST (INPUT,OUTPUT,TAPE5=INPUT,TAPE6=OUTPUT)
 STATIC ANALYSIS OF SHEAR WALL WITH FLOORS
 M. ENG. THESIS BISWAS

COMPUTER NOTATIONS USED

XI MOMENT OF INERTIA OF SECTION ABOUT X AXIS
 YI MOMENT OF INERTIA OF SECTION ABOUT Y AXIS
 WI SECTORIAL MOMENT OF INERTIA OF SECTION
 DI POLAR MOMENT OF INERTIA ABOUT SHEAR CENTER
 E MODULUS OF ELASTICITY
 G MODULUS OF RIGIDITY
 H FLOOR HEIGHT
 NF NUMBER OF STOREY
 DM BIMOMENT CONTRIBUTION FACTOR
 PLY(I) LATERAL LOAD AT I TH FLOOR IN Y DIRECTION
 PLR(I) APPLIED TORQUE AT I TH FLOOR

DIMENSION P(8,8),F(8,8),PFA(8,8),PFB(8,8),PF(8,8),A(8,8),T(8,8),
 1AT(8,1),BT(8,1),N1(32),WORK(8),PLY(50),PLR(50),S(8,1),PS(8,1)

INPUT DATA

XI= 42.7089
 YI= 11.0322
 WI= 108.6931
 DI= .09767
 E= .4E+06
 PSN= .35
 G= E/(2.*(1.+PSN))
 NF= 8
 H= 6.
 DM= 20.109E+04
 DM= 124.909E+04
 DM= 0.
 DO 75 I=1,NF
 PLR(I)= 0.
 75 PLY(I)= 0.
 PLY(8)= 12.5
 PLR(8)= 6.2425*12.5

PRINTING OUT INPUT DATA

WRITE (6,1)

1 FORMAT (1X,* INPUT DATA *)

WRITE (6,2) NF,H,XI,YI,WI,DI,PSN,E,G,E ,DM

2 FORMAT (5X,15,10E12.4)

WRITE (6,3)

3 FORMAT (1X,* FLOOR LOADS *)

7 FORMAT (1X,110,5X,3E18.6)

DO 70 I=1,NF

FIELD TRANSFER MATRIX F

70 WRITE (6,7) I,PLY(I),PLR(I)

GD= G*DI

AK= H*SQRT(GD/(E*WI))

AL= AK/H

DO 4 I= 1,8

DO 4 J= 1,8

IF (I.NE.J) F(I,J)= 0.

4 F(J,J)= 1.

```

000143 F(1,2) = H
000145 F(1,3) = -(H*H)/(2.*E*XI)
000150 F(1,4) = -(H*H*H)/(6.*E*XI)
000153 F(2,3) = -H/(E*XI)
000155 F(2,4) = -(H*H)/(2.*E*XI)
000160 F(3,4) = H
000161 F(5,6) = SINH(AK)/AL
000164 F(5,7) = (1. - COSH(AK))/GD
000170 F(5,8) = (H. - SINH(AK)/AL)/GD
000175 F(6,6) = COSH(AK)
000200 F(6,7) = -(AL*SINH(AK))/GD
000204 F(6,8) = (1. - COSH(AK))/GD
000210 F(7,8) = -GD*SINH(AK)/AL
000214 F(7,7) = COSH(AK)
000216 F(7,8) = SINH(AK)/AL
C POINT TRANSFER MATRIX
DO 6 I= 1,8
DO 6 J= 1,8
IF (I.NE.J) P(I,J)= 0.
6 P(J,J)= 1.
P(7,6) = -DM
32 CALL FLOOR(P,F,NF,A ,PFA,PFB,PF)
DO 10 I= 1,8
10 AT(I,1)= 0
AT(4,1)= -PLY(NF)
AT(8,1)= -PLR(NF)
WRITE (6,9)
9 FORMAT (1X, * FIELD TRANSFER MATRIX *)
WRITE (6,5) (( F(I,J),J=1,8),I=1,8)
WRITE (6,13)
13 FORMAT (1X, * POINT TRANSFER MATRIX *)
WRITE (6,5) (( P(I,J),J=1,8),I=1,8)
5 FORMAT (1X,8E15.6)
NC= 1
11 CALL FLOOR(P, F,NC,T,PFA,PFB,PF)
NCC= NF- NC
DO 12 I= 1,8
12 AT(I,1)= AT(I,1)- T(I,4)*PLY(NCC)- T(I,8)*PLR(NCC)
NC= NC+1
IF (NCC.EQ.1) GO TO 14
GO TO 11
14 DO 16 I= 1,8
DO 16 J= 1,8
16 A(I,J)= -A(I,J)
DO 18 I= 1,4
DO 18 J= 1,2
K= I+4
L= J+4
A(I,J)= 0.
18 A(K,L)= 0.
A(1,1)= 1.
A(2,2)= 1.
A(5,5)= 1.
A(6,6)= 1.
CALL MINVSE(A,8,8,1.E-08,IERR,N1,WORK)
IF (IERR.EQ.0) GO TO 30
57 FORMAT (5X, * UNSUCCESSFUL INVERSION *)
WRITE (6,57)

```

```

000425 GO TO 100
000426 30 CALL GMPRD(A,AT,BT,8,8,1)
000432 BT(1,1)= 0.
000433 BT(2,1)= 0.
000434 BT(5,1)= 0.
000435 BT(6,1)= 0.
C BT IS THE STATE VECTOR AT BASE
000436 WRITE (6,31)
000441 31 FORMAT (1X)
000441 WRITE (6,37)
000445 37 FORMAT (4X, *FLOOR *,10X, *DISPLACEMENT*,7X, *SLOPE*,6X,
1 *MOMENT*,6X, *SHEAR*, 6X, *ROTATION*, 6X, *TWIST*, 6X,* BIMOMENT
1 *,6X,*TORQUE* )
000445 WRITE (6,33) ((BT(I,J),J= 1,1),I= 1,8)
000463 33 FORMAT (4X, *0*,4X,8E14.4)
000463 NA= 1
000464 17 DO 27 I= 1,8
000466 27 AT(I,1)= 0.
000472 AT(4,1)= -PLY(NA)
000474 AT(8,1)= -PLR(NA)
000475 CALL GMPRD(F,BT,S,8,8,1)
000501 WRITE (6,21) NA,((S(I,J),J= 1,1),I= 1,8)
000521 20 FORMAT (4X,13,* BELOW*, 8E14.4)
000521 CALL GMPRD(P,S,PS,8,8,1)
000525 DO 22 I= 1,8
000527 22 PS(I,1)= PS(I,1)+ AT(I,1)
000534 WRITE (6,21) NA,((PS(I,J),J= 1,1),I= 1,8)
000554 21 FORMAT (4X,13,* ABOVE *, 8E14.4)
000554 IF (NA.EQ.NF) GO TO 100
000556 DO 24 I= 1,8
000560 24 BT(I,1)= PS(I,1)
000565 NA = NA+ 1
000566 GO TO 17
000567 100 STOP
000571 END

```

UNUSED COMPILER SPACE

0004000


```

C      SUBROUTINE FLOOR(P,F,N,T,PFA,PFB,PF )
C      MULTIPLICATION OF MATRICES ACCORDING TO NO OF FLOOR
C      P, F ARE THE INPUT MATRICES
C      N IS THE NO OF FLOORS
C      T IS THE OUTPUT MATRIX
000012 PFA, PFB, PF ARE THE WORKING MATRICES OF SIZE 8 * 8
000012 DIMENSION P(8,8),F(8,8),T(8,8),PFA(8,8),PFB(8,8),PF(8,8)
000014 CALL MULT(P,F,PF,8)
000015 KOUNT=1
000024 IF (N.GT.1) GO TO 4
000025 DO 10 I=1,8
000026 DO 10 J=1,8
000041 10 T(I,J)= PF(I,J)
000042 GO TO 30
000044 4 DO 5 I=1,8
000045 DO 5 J=1,8
000060 5 PFA(I,J)= PF(I,J)
000063 6 CALL MULT(PFA,PF,PFB,8)
000065 KOUNT = KOUNT+1
000072 DO 7 I=1,8
000073 DO 7 J=1,8
000106 7 PFA(I,J)= PFB(I,J)
000110 IF (KOUNT.EQ.N) GO TO 9
000110 GO TO 6
000112 9 DO 12 I=1,8
000113 DO 12 J=1,8
000126 12 T(I,J)= PFB(I,J)
000127 30 RETURN
      END

```

```

UNUSED COMPILER SPACE
006200

```

```

      C      SUBROUTINE MULT(A,B,C,N)
      MULTIPLICATION OF SQUARE MARRICES
      DIMENSION A(N,N),B(N,N),C(N,N)
000007      DO 1 I=1,N
000007      DO 1 J=1,N
000010      C(I,J)=0.
000011      DO 1 K=1,N
000015      C(I,J)= C(I,J)+ A(I,K)*B(K,J)
000016
000032      1  CONTINUE
000040      RETURN
000040      END

```

```

UNUSED COMPILER SPACE
000500

```

00000000000000000000000000000000

C

33

1

139

000116

```

2 FORMAT (1X,* SECTION PROPERTIES*/9X,* AREA =*,F9.4/9X,* MOMENT OF
1INERTIA ABOUT X=*,F9.4/9X,* SECTORIAL MOMENT OF INERTIA =*,F9.4/
29X,* TORSION FACTOR =*,F9.4/9X,* POLAR MOMENT OF INERTIA =*,F9.4/
39X,* ORDINATE OF SHEAR CENTRE =*,F9.4/1X,* DIAPHRAGM PROPERTIES*/9
4X,* MASS OF PLATE =*,E15.6/9X,* MASS MOMENT OF INERTIA ABOUT X=*,
5E15.6/9X,* POLAR MASS MOMENT OF INERTIA =*,E15.6/9X,* BIMOMENT FAC
6TOR =*,E15.6/9X,* SHEAR CENTRE TO C.G. OF THE PLATE =*,F9.4)
1 FORMAT (6F9.4/7F9.4/3E15.6)

```

000116

000116

000122

000122

000123

000124

000125

```

42 FORMAT (1X,* RESULTS FOR FIXED FREE CASE * )

```

```

    KKK= 333

```

```

    NCOUNT= 6666

```

```

    DETO= 0.

```

```

    NFREQ= 0

```

```

C    POINT TRANSFER MATRIX P

```

```

40 DO 4 I=1,8

```

```

    DO 4 J=1,8

```

```

        IF (I.NE.J) P(I,J)=0.

```

```

4    P(J,J)=1.

```

```

        P(3,2)=-XJ*W*W

```

```

        P(3,6)=-AP*AM*W*W

```

```

        P(4,1)=-AG*W**2

```

```

        P(4,5)=-AP*AM*W**2

```

```

        P(7,6)=-DM

```

```

        P(8,1)=-AP*AM*W**2

```

```

        P(8,5)=-ZJ*W**2

```

```

C    CUEFFECIENTS OF CHARACTERISTICS EQUATION

```

```

    DO 6 I=1,9

```

```

6    AA(I)=0.

```

```

        AA(1)=-WI*XI*E**2

```

```

        AA(3)=UI*XI*E*G-2.*WI*XI*RO*E*W**2

```

```

        AA(5)=(WI*A*E+UI*XI*G+PI*XI*E)*RO*(W**2) -WI*XI*(W**4)*(RO**2)

```

```

        AA(7)=-DI*A*RO*G*(W**2) +(PI*XI+WI*A)*(W**4)*(RO**2)

```

```

        AA(9)=(-PI*A+(A*AX)**2)*(W**4)*(RO**2)

```

```

C    SOLUTION BY SUBROUTINE BAIRST

```

```

    CALL BAIRST (AA,RR,RI,8)

```

```

    DO 5 I=1,8

```

```

C    5 RC(I)=CMPLX(RR(I),RI(I))

```

```

    FIELD TRANSFER MATRIX F

```

```

    DO 11 J=1,8

```

```

        C(1,J)=CMPLX(1.,0.)

```

```

        C(2,J)=RC(J)

```

```

        C(3,J)=E*XI*(RC(J)**2)

```

```

        C(4,J)=-E*XI*(RC(J)**3)

```

```

        C(5,J)=-((F*XI)/(W**2)*RO*A*AX))*(RC(J)**4)-(XI/(A*AX))*(RC(J)**2

```

```

1) + (1./AX)
        C(6,J)=-((E*XI)/((W**2)*RO*A*AX))*(RC(J)**5)-(XI/(A*AX))*(RC(J)**3

```

```

1) + (1./AX)*RC(J)
        C(7,J)=-((E**2)*WI*XI)/((W**2)*RO*A*AX))*(RC(J)**6)+((E*WI*XI)/

```

```

1(A*AX))*(RC(J)**4)-((E*WI)/AX)*(RC(J)**2)

```

```

        C(8,J)=-((7E**2)*WI*XI)/((W**2)*RO*A*AX))*(RC(J)**7)+((E*WI*XI)/

```

```

1(A*AX)-(E*G*DI*XI)/((W**2)*RO*A*AX))*(RC(J)**5)+((-E*WI)/AX-(G*DI*

```

```

2XI)/(A*AX))*RC(J)**3)+((G*DI)/AX)*RC(J)

```

```

11 CONTINUE

```

```

    DO 13 I=1,8

```

```

        DO 13 J=1,8

```

```

            IF (I.NE.J) D(I,J)=0.

```

```

            D(J,J)=CEXP(H*RC(J))

```

000611

000613

000614

000615

000622

```

000640      13 B(I,J)=C(I,J)
000653      CALL INVCPIX(B,8,8,1.E-8,IERR,NI,TEMP )
000662      CALL MATMOLT(C,D,CD,8)
000665      CALL MATMOLT(CD,B,Q,8)
000670      DO 24 I=1,8
000672      DO 24 J=1,8
000673      24 F(I,J)=REAL(Q(I,J))
000705      CALL FLOOR(P,F,NFLOOR,T,PFA,PFR,PF )
000713      CALL UCFXFR(T,TT)
000715      CALL DETVAL(TT,DET,4)
000720      46 PROD= DETO*DET
000722      IF (NCOUNT.EQ.5555) GO TO 47
000724      IF (PROD.GE.0.) GO TO 8
000726      44 DETN=DET
000730      45 WAP= WO+(W-WO)*DETO/(DETO-DETN)
000737      PC= (WAP-WO)/WAP
000741      IF (PC.LT..01 ) GO TO 55
000743      WT= W
000745      W= WAP
000745      NCOUNT= 5555
000746      GO TO 40
000747      47 IF (PROD.LT.0.) GO TO 44
000751      DETO= DET
000752      KKK= 222
000753      WO= W
000755      W= WT
000756      GO TO 45
000756      55 NFREQ= NFREQ+1
000760      WW= WAP/6.28
000762      WRITE (6,60) NFREQ,WW
000771      60 FORMAT (7X,'FREQ #, 13,5X,F10.4 )
000771      8 DETO= DET
000773      IF (KKK.EQ.222 ) DETO= DETN
000776      KKK= 111
000777      WO= W
001001      NCOUNT= 3333
001002      W= W+ WINC
001003      IF (NFREQ.EQ.NFQ) GO TO 50
001005      GO TO 40
001006      50 STOP
001010      END

```

UNUSED COMPILER SPACE

003100

		SUBROUTINE INVCPIX(A,N,NN,ZERO,IERR,NI,TEMP)	MHG00001
		REM	MHG00005
		REM THE ROUTINE INVERTS THE MATRIX BY A PIVOTAL METHOD	MHG00010
		REM USING THE LARGEST ELEMENT IN THE NEXT ROW AS A PIVOT. THE	MHG00020
		REM FIRST PART DOES THE INVERSION AND THE SECOND PART	MHG00030
		REM REARRANGES THE ROWS AND COLUMNS TO TAKE ACCOUNT OF THE	MHG00040
		REM FACT THAT THE PIVOTS WERE NOT ON THE DIAGONAL.	MHG00050
		REM THIS SUBROUTINE IS CALLED BY CALL INVMAT(A,N,NN,ZERO,IERR,	MHG00060
		REM NI,TEMP) WHERE A IS AN NN BY NN MATRIX IN AN N BY N ARRAY	MHG00070
		REM ZERO IS A TEST VALUE BELOW WHICH AN ELEMENT IS CONSIDERED	MHG00080
		REM TO BE ZERO, IERR IS AN INDICATOR WHICH IS 0 IF THE INVERSE	MHG00090
		REM IS FOUND AND NON-ZERO OTHERWISE, NI IS A WORKING ARRAY	MHG00100
		REM OF NN ELEMENTS AND TEMP IS ANOTHER	MHG00110
		REM	MHG00120
		DIMENSION A(N,1),NI(NN),TEMP(NN)	MHG00140
		COMPLEX A,AA,TEMP	MHG00145
		DO 10 I=1,NN	MHG00150
000012		NI(I)=0	MHG00160
000012	10	DO 20 K=1,NN	MHG00170
000013		AAA=0.	MHG00180
000016		DO 30 J=1,NN	MHG00190
000017		IF (NI(J).NE.0) GO TO 30	MHG00200
000020		C=CAUS(A(K,J))	MHG00215
000022		IF (C.LT.AAA) GO TO 30	MHG00220
000024		AAA=C	MHG00230
000036		JJ=J	MHG00240
000041	30	CONTINUE	MHG00250
000041		IF (AAA.GT.ZERO) GO TO 70	MHG00260
000043		IERR=K	MHG00270
000046		RETURN	MHG00280
000051	70	NI(JJ)=K	MHG00290
000052		AA=A(K,JJ)	MHG00300
000055		A(K,JJ)=1.	MHG00310
000062		DO 50 J=1,NN	MHG00320
000070	50	A(K,J)=A(K,J)/AA	MHG00330
000072		DO 60 I=1,NN	MHG00340
000107		IF (I.EQ.K) GO TO 60	MHG00350
000110		AA=-A(I,JJ)	MHG00360
000112		A(I,JJ)=0.	MHG00370
000120		DO 55 J=1,NN	MHG00380
000125	55	A(I,J)=A(I,J)+AA*A(K,J)	MHG00390
000126	60	CONTINUE	MHG00400
000145	20	CONTINUE	MHG00410
000150		IERR=0	MHG00420
000152		DO 98 I=1,NN	MHG00430
000152	98	IF (NI(I).NE.1) GO TO 95	MHG00440
000154		RETURN	MHG00450
000160	95	DO 80 I=1,NN	MHG00460
000160		DO 79 J=1,NN	MHG00470
000162		K=NI(J)	MHG00480
000163	79	TEMP(K)=A(I,J)	MHG00490
000165		DO 81 J=1,NN	MHG00500
000200	81	A(I,J)=TEMP(J)	MHG00510
000201	80	CONTINUE	MHG00520
000214		DO 90 J=1,NN	MHG00530
000216		DO 89 I=1,NN	MHG00540
000217		K=NI(I)	
000220			

```

000222      89  TEMP(I)=A(K,J)
000235      DO 91 I=1,NN
000236      91  A(I,J)=TEMP(I)
000253      90  CONTINUE
000255      RETURN
000256      END

```

```

MHG00550
MHG00560
MHG00570
MHG00580
MHG00590
MHG00600

```

```

UNUSED COMPILER SPACE
005000

```

```

C
C
C
C
SUBROUTINE FLOOR(P,F,N,T,PFA,PFB,PF)
MULTIPLICATION OF MATRICES ACCORDING TO NO OF FLOOR
P, F ARE THE INPUT MATRICES
N IS THE NO OF FLOORS
T IS THE OUTPUT MATRIX
PFA, PFB, PF ARE THE WORKING MATRICES OF SIZE 8*8
DIMENSION P(8,8),F(8,8),T(8,8),PFA(8,8),PFB(8,8),PF(8,8)
CALL MULT(P,F,PF,8)
KOUNT=1
IF (N.GT.1) GO TO 4
DO 10 I=1,8
DO 10 J=1,8
10 T(I,J)=PF(I,J)
GO TO 30
4 DO 5 I=1,8
DO 5 J=1,8
5 PFA(I,J)=PF(I,J)
6 CALL MULT(PFA,PF,PFB,8)
KOUNT=KOUNT+1
DO 7 I=1,8
DO 7 J=1,8
7 PFA(I,J)=PFB(I,J)
IF (KOUNT.EQ.N) GO TO 9
GO TO 6
9 DO 12 I=1,8
DO 12 J=1,8
12 T(I,J)=PFB(I,J)
30 RETURN
END

```

```

UNUSED COMPILER SPACE
000200

```

```

000007 C SUBROUTINE MATMULT (A,B,C,N)
000007 MULTIPLICATION OF COMPLEX MATRICES
000007 DIMENSION A(N,N),B(N,N),C(N,N)
000010 COMPLEX A,B,C
000011 DO 1 I=1,N
000016 DO 1 J=1,N
000011 C(I,J)=0.
000016 DO 1 K=1,N
000020 C(I,J)=C(I,J)+A(I,K)*B(K,J)
000042 1 CONTINUE
000050 RETURN
000050 END

```

UNUSED COMPILER SPACE

000400

```

000005 C SUBROUTINE BCFXFR(T,TT)
000005 BOUNDARY CONDITION ONE END FIXED OTHER END FREE
000005 DIMENSION T(8,8),TT(4,4)
000006 DO 30 I=1,4
000006 DO 30 J=1,4
000007 K=I+2
000011 N=J+2
000012 KK=I+4
000013 NN=J+4
000014 IF (I.LE.2.AND.J.LE.2) TT(I,J)=T(K,N)
000030 IF (I.LE.2.AND.J.GT.2) TT(I,J)=T(K,NN)
000046 IF (I.GT.2.AND.J.LE.2) TT(I,J)=T(KK,N)
000064 IF (I.GT.2.AND.J.GT.2) TT(I,J)=T(KK,NN)
000101 30 CONTINUE
000105 RETURN
000106 END

```

UNUSED COMPILER SPACE

000300


```

C      SUBROUTINE DETVAL(A,DET,M)
      EVALUATION OF DETERMINANT
      DIMENSION A(M,M)
      NE=0
      MM=M-1
      DO 8 J=1,MM
      JJ=J+1
      DO 30 JB=JJ,M
      IF (ABS(A(J,J))_GE.ABS(A(JB,J))) GO TO 30
      NE=NE+1
      IF (NE.EQ.2) NE=0
      DO 28 KK=1,M
      HOLD=A(J,KK)
      A(J,KK)=A(JB,KK)
28    A(JB,KK)=HOLD
30    CONTINUE
      DO 4 NN=JJ,M
      C=A(NN,J)/A(J,J)
      DO 4 N=1,M
      4    A(NN,N)=A(NN,N)-A(J,N)*C
      8    CONTINUE
      DET=1.
      DO 5 J=1,M
      5    DET=DET*A(J,J)
      IF (NE.EQ.1) DET=-DET
      RETURN
      END

```

UNUSED COMPILER SPACE
006200

```

C      SUBROUTINE MULT(A,B,C,N)
      MULTIPLICATION OF SQUARE MARRICES
      DIMENSION A(N,N),B(N,N),C(N,N)
      DO 1 I=1,N
      DO 1 J=1,N
      C(I,J)=0.
      DO 1 K=1,N
      1    C(I,J)=C(I,J)+A(I,K)*B(K,J)
      CONTINUE
      RETURN
      END

```

UNUSED COMPILER SPACE
006500

```

PROGRAM TST (INPUT,OUTPUT,TAPE5=INPUT,TAPE6=OUTPUT)
NON PLANAR COUPLED SHEAR WALL
M. ENG. THESIS BISWAS
COMPUTER NOTATIONS USED

```

```

SX1,SX2,      MOMENT OF INERTIA OF PIERS ABOUT PARALLEL GLOBAL X TH C.G.
SY1,SY2      MOMENT OF INERTIA OF PIERS ABOUT PARALLEL GLOBAL Y TH C.G.
SXY1,SXY2    PRODUCT MOMENT OF INERTIA OF PIERS
SW1,SW2      SECTORIAL MOMENT OF INERTIA OF PIERS W.R.T SHEAR CENTRE
SJ1,SJ2      TORSION FACTORS OF PIERS
EX1,EY1,EX2,EY2 COORDINATES OF C.G. OF PIERS W.R.T GLOBAL AXES
CX1,CY1,CX2,CY2 COORDINATES OF S.C. OF PIERS W.R.T GLOBAL AXES
A1,A2        AREA OF THE PIERS
W1,W2        SECTORIAL COORDINATES
E, G, PSN    MODULUS OF ELASTICITY, MODULUS OF RIGIDITY, POISONS RATIO
L            CLEAR SPAN OF THE CONNECTING BEAM
SB           MOMENT OF INERTIA PF CONNECTING BEAM
AST          EFFECTIVE SHEAR AREA OF CONNECTING BEAM
H            SPACING OF CONNECTING BEAM
NF           NUMBER OF FLOORS
PX           CONCENTRATED LOAD AT TOP IN X DIRECTION
PY           CONCENTRATED LOAD AT TOP IN Y DIRECTION
QT           CONCENTRATED TORQUE AT TOP
HT           TOTAL HEIGHT OF STRUCTURE
DIMENSION AAA(5,5),BBB(5)

```

```

INPUT DATA
SX1= 9.377
SX2= 9.377
SY1= 2.289
SY2= 2.289
SXY1= -2.594
SXY2= 2.594
SW1= 0.
SW2= 0.
SJ1= .159
SJ2= .159
EX1= -4.313
EX2= 4.313
EY1= 1.687
EY2= 1.687
CX1= -4.804
CX2= 4.804
CY1= 0.
CY2= 0.
A1= 3.128
A2= 3.128
W1= 0.
W2= 0.
E= .4E+06
PSN= .35
G= E/(2.*(1.+PSN))
L= 4.
SB= .10997
AST= .5865
H= 6.
NF= 8

```

```

000056      C      PX= 25.
000060      PY= 0.
000061      QT= 0.
000062      C      COMPUTATION STARTS
000064      HT= FLOAT(NF)*H
000074      D= CX2*EY2- CY2*EX2+ CY1*EX1- CX1*EY1
000075      A= EX2- EX1
000076      B= EY2- EY1
000100      W= W2- W1
000110      RFACT= 1.+ (12.*E*SB)/(L*L*G*AST)
000112      SJB= SB/(H*RFACT)
000114      SJ= SJ1+ SJ2
000116      SX= SX1+ SX2
000120      SY= SY1+ SY2
000122      SXY= SXY1+ SXY2
000132      SYC= CY1*SY1+ CY2*SY2- CX1*SXY1- CX2*SXY2
000141      SXC= CX1*SX1+ CX2*SX2- CY1*SXY1- CY2*SXY2
      SW= CX1*CX1*SX1+ CX2*CX2*SX2+ CY1*CY1*SY1+ CY2*CY2*SY2
      1-2.*CX1*CY1*SXY1- 2.*CX2*CY2*SXY2+ SW1+ SW2
000163      AA= SY*SXC+ SXY*SYC
000167      BB= SXY*SXC+ SX*SYC
000172      CC= A*SXC+ B*SYC
000174      SXP= SX1+ SX2
000176      SYP= SY1+ SY2
000200      SWP= SW- (SXC*SXC)/SX- (SYC*SYC)/SY
000206      IF (AA.EQ.0.) GO TO 40
000207      SX= SX- SXY*(BB/AA)
000212      SY= SY- SXY*(AA/BB)
000214      AP= A- SXY*(CC/BB)
000217      BP= B- SXY*(CC/AA)
000222      P= (SXP*SYC*SYC)/(SY*BB)+ (SXY*SXC*SXC)/(SX*AA)
000233      Q= (SYP*SXC*SXC)/(SX*AA)+ (SXY*SYC*SYC)/(SY*BB)
000244      R= (SXP*SYC*A)/(SY*BB)- (SXY*SXC*B)/(SX*AA)
000256      S= (SYP*SXC*B)/(SX*AA)- (SXY*SYC*A)/(SY*BB)
000270      GO TO 42
000271      40 SX= SXP
000273      SY= SYP
000274      AP= A
000276      BP= B
000277      P= SYC/SY
000301      Q= 0.
000302      R= A/SY
000303      S= 0.
000304      42 WP= W+ D+ (SYC*A)/SY+ (SXC*B)/SX
000315      WB= -W+ (BP*SXC)/SX- (AP*SYC)/SY -D
000326      FF= (A*AP)/SY+ (B*BP)/SX+ 1./A1+ 1./A2
000336      F= 1./FF
000337      ALPH= (L*L*L)/(12.*SJB)
000345      CLTH5= ALPH*E*SWP
000347      CLTH3= -((E*SWP)/F+ ALPH*G*SJ- E*WP*WB)
000356      CLTH1= (G*SJ)/F
000360      CRMYP1= -(R*WB+P/F)
000363      CRMXP1= -(S*WB- Q/F)
000366      CRMYP3= ALPH*P
000370      CRMXP3= -ALPH*Q
000371      CRQTP2= -ALPH
000372      CRQT= 1./F

```

```

000374      WRITE (6,2) SX1,SX2,SY1,SY2,SXY1,SXY2,SW1,SW2,SJ1,SJ2,EX1,EX2,EY1,
000471      1EY2,CX1,CX2,CY1,CY2,A1,A2,W1,W2,E,G,PSN,L,SB,AST,H
000475      WRITE (6,3)
000475      3 FORMAT (10X)
      WRITE (6,2) SXC,SYC,SW,AA,BB,CC,SX,SY,SXP,SYP,AP,BP,SWP,WB,P,Q,
      1WP,F,R,S,ALPH,CLTH5,CLTH3,CLTH1,CRMYP1,CRMXP1,CRMYP3,CRMXP3,
      2CRQT,CRQTP2
000575      2 FORMAT(1X,6E15.6)
000575      RAS= (-CLTH3- SQRT(CLTH3*CLTH3- 4.*CLTH5*CLTH1))/(2.*CLTH5)
000607      RBS= (-CLTH3+ SQRT(CLTH3*CLTH3- 4.*CLTH5*CLTH1))/(2.*CLTH5)
000621      RA= SQRT(RAS)
000623      RB= SQRT(RBS)
000625      CP= -(CRMYP1*PX+ CRMXP1*PY)/CLTH1 + (CRQT*QT)/CLTH1
000635      ALP= ALPH*E*SWP
000637      BET= -(ALPH*G*SJ- E*WP*WB)
000645      RA5= RA*RA*RA*RA*RA
000647      RA4= RA*RA*RA*RA
000650      RA3= RA*RA*RA
000651      RA2= RA*RA
000652      RB5= RB*RB*RB*RB*RB
000654      RB4= RB*RB*RB*RB
000655      RB3= RB*RB*RB
000656      RB2= RB*RB
000657      RAT= RA*HT
000660      RBT= RB*HT
000661      WRITE (6,5)
000664      5 FORMAT (10X)
000664      WRITE(6,2) RAS,RBS,RA,RB,CP,BET,ALP,RAT,RBT
000712      DO 4 I= 1,5
000714      DO 4 J= 1,5
000715      4 AAA(I,J)= 0
000724      AAA(1,1)= 1.
000725      AAA(1,2)= 1.
000726      AAA(1,4)= 1.
000727      AAA(2,3)= RA
000730      AAA(2,5)= RB
000732      AAA(3,3)= RA*RA*RA*E*SWP- G*SJ*RA
000737      AAA(3,5)= RB*RB*RB*E*SWP- G*SJ*RB
000745      AAA(4,2)= (ALP*RA4+ BET*RA2)*COSH(RAT)
000753      AAA(4,3)= (ALP*RA4+ BET*RA2)*SINH(RAT)
000762      AAA(4,4)= (ALP*RB4+ BET*RB2)*COSH(RBT)
000771      AAA(4,5)= (ALP*RB4+ BET*RB2)*SINH(RBT)
001000      AAA(5,2)= RA2*COSH(RAT)
001004      AAA(5,3)= RA2*SINH(RAT)
001007      AAA(5,4)= RB2*COSH(RBT)
001012      AAA(5,5)= RB2*SINH(RBT)
001015      BBB(1)= 0.
001016      BBB(2)= -CP
001017      BBB(3)= G*SJ*CP-P*PX+ Q*PY -QT
001027      BBB(4)= 0.
001030      BBB(5)= 0.
001031      WRITE (6,7) AAA
001036      7 FORMAT (3X,5E18.6)
001036      WRITE (6,7) BBB
001044      CALL SIMQ(AAA,BBB,5,KS)
001047      WRITE (6,9) KS
001055      9 FORMAT (20X,I3)
001055      WRITE (6,7) BBB

```

```

001063      C1= BBB(1)
001065      C2= BBB(2)
001066      C3= BBB(3)
001070      C4= BBB(4)
001071      C5= BBB(5)
001073      Z= 0.
001074      WRITE (6,11)
001077      11 FORMAT(10X)
001077      10 RAZ= RA*Z
001101      RBZ= RB*Z
001103      THET= C1+ C2*COSH(RAZ)+ C3*SINH(RAZ)+ C4*COSH(RBZ)+ C5*SINH(RBZ)+
1CP*Z
001126      THETP1= C2*RA*SINH(RAZ)+ C3*RA*COSH(RAZ)+ C4*RB*SINH(RBZ)+ C5*
1RB*COSH(RBZ)+ CP
001150      THETP2= C2*RA2*COSH(RAZ)+ C3*RA2*SINH(RAZ)+ C4*RB2*COSH(RBZ)+
1C5*RB2*SINH(RBZ)
001172      THETP3= C2*RA3*SINH(RAZ)+ C3*RA3*COSH(RAZ)+ C4*RB3*SINH(RBZ)+
1C5*RB3*COSH(RBZ)
001214      WRITE (6,12) THET,THETP1,THETP2,THETP3
001230      12 FORMAT (1X,6E20.6)
001230      IF (Z.EQ.HT) GO TO 13
001232      Z= Z+H
001234      GO TO 10
001234      13 B1= (C2*(-E*SWP*RA3+ G*SJ*RA))/WB
001244      B2= (C3*(-E*SWP*RA3+ G*SJ*RA))/WB
001253      B3= (C4*(-E*SWP*RB3+ G*SJ*RB))/WB
001262      B4= (C5*(-E*SWP*RB3+ G*SJ*RB))/WB
001271      B5= (G*SJ*CP-P*PX+ Q*PY -QT)/WB
001302      D1= -B1/RA
001304      D2= -B2/RA
001305      D3= -B3/RB
001307      D4= -B4/RB
001311      D5= -B5
001312      D6= (B1*COSH(RAT)+ B2*SINH(RAT))/RA+ (B3*COSH(RBT)+ B4*SINH(RBT))
1/RB+ B5*HT
001337      Z= 0.
001340      WRITE (6,19)
001343      19 FORMAT (10X)
001343      15 RAZ= RA*Z
001345      RBZ= RB*Z
001347      FQ= B1*SINH(RAZ)+ B2*COSH(RAZ)+ B3*SINH(RBZ)+ B4*COSH(RBZ)+ B5
001367      FT= D1*COSH(RAZ)+ D2*SINH(RAZ)+ D3*COSH(RBZ)+ D4*SINH(RBZ)+
1 D5*Z+ D6
001412      THET= C1+ C2*COSH(RAZ)+ C3*SINH(RAZ)+ C4*COSH(RBZ)+ C5*SINH(RBZ)+
1CP*Z
001436      THETP1= C2*RA*SINH(RAZ)+ C3*RA*COSH(RAZ)+ C4*RB*SINH(RBZ)+ C5*
1RB*COSH(RBZ)+ CP
001460      THETP2= C2*RA2*COSH(RAZ)+ C3*RA2*SINH(RAZ)+ C4*RB2*COSH(RBZ)+
1C5*RB2*SINH(RBZ)
001502      THETP3= C2*RA3*SINH(RAZ)+ C3*RA3*COSH(RAZ)+ C4*RB3*SINH(RBZ)+
1C5*RB3*COSH(RBZ)
001524      BIMNT1= -E*SW1*THETP2
001527      BIMNT2= -E*SW2*THETP2
001531      TRQ1= -E*SW1*THETP3+ G*SJ1*THETP1
001536      TRQ2= -E*SW2*THETP3+ G*SJ2*THETP1
001544      WRITE (6,17) FQ,FT,BIMNT1,BIMNT2,TRQ1,TRQ2
001563      17 FORMAT(1X,6E16.6)
001563      IF (Z.EQ.HT) GO TO 29

```

```

001565      Z= Z+H
001567      GO TO 15
001567 29 Z1= (SYC*C2*RA2)/SY- (AP*D1)/(E*SY)
001600      Z2= (SYC*C3*RA2)/SY- (AP*D2)/(E*SY)
001611      Z3= (SYC*C4*RB2)/SY- (AP*D3)/(E*SY)
001622      Z4= (SYC*C5*RB2)/SY- (AP*D4)/(E*SY)
001633      Z5= -(AP*D5)/(E*SY)- (SXP*SYC*PX)/(E*SY*BB)+ (SXY*SYC*PY)/(E*SY*BB)
001652      1 )
      Z6= -(AP*D6)/(E*SY)+ (SXP*SYC*HT*PX)/(E*SY*BB)- (SXY*SYC*HT*PY)
001672      1/(E*SY*BB)
001676      Z7= -Z2/RA- Z4/RB
001702      Z8= -Z1/RA2- Z3/RB2
001713      E1=-(SXC*C2*RA2)/SX- (BP*D1)/(E*SX)
001713      E2=-(SXC*C3*RA2)/SX- (BP*D2)/(E*SX)
001724      E3=-(SXC*C4*RB2)/SX- (BP*D3)/(E*SX)
001735      E4=-(SXC*C5*RB2)/SX- (BP*D4)/(E*SX)
001746      IF (AA.EQ.0.) GO TO 50
001747      E5= -(BP*D5)/(E*SX)+ (SXY*SXC*PX)/(E*SX*AA)- (SYP*SXC*PY)/
001765      1 (E*SX*AA)
      E6= -(BP*D6)/(E*SX)- (SXY*SXC*HT*PX)/(E*SX*AA)+ (SYP*SXC*HT*PY)
002005      1/(E*SX*AA)
002005      GO TO 52
002010 50 E5= -PY/(E*SX)
002014      E6= (PY*HT)/(E*SX)
002020 52 E7= -E2/RA- E4/RB
002024      E8= -E1/RA2- E3/RB2
002025      Z= 0.
002031      WRITE(6,22)
002031      22 FORMAT(10X)
002031 20 RAZ= RA*Z
002033      RBZ= RB*Z
002035      ZETP2= Z1*COSH(RAZ)+ Z2*SINH(RAZ)+ Z3*COSH(RBZ)+ Z4*SINH(RBZ)+
002057      1 Z5*Z+ Z6
      ZETP1= (Z1/RA)*SINH(RAZ)+ (Z2/RA)*COSH(RAZ)+ (Z3/RB)*SINH(RBZ)+
002113      1 (Z4/RB)*COSH(RBZ)+ (Z5/2.)*Z*Z+ Z6*Z+ Z7
      ZET= (Z1/RA2)*COSH(RAZ)+ (Z2/RA2)*SINH(RAZ)+ (Z3/RB2)*COSH(RBZ)+
002153      1 (Z4/RB2)*SINH(RBZ)+ (Z5/6.)*Z*Z*Z+ (Z6/2.)*Z*Z+ Z7*Z+ Z8
      ETAP2= E1*COSH(RAZ)+ E2*SINH(RAZ)+ E3*COSH(RBZ)+ E4*SINH(RBZ)+
002176      1 E5*Z+ E6
      ETAP1= (E1/RA)*SINH(RAZ)+ (E2/RA)*COSH(RAZ)+ (E3/RB)*SINH(RBZ)
002232      1 + (E4/RB)*COSH(RBZ)+ (E5/2.)*Z*Z+ E6*Z+ E7
      ETA= (E1/RA2)*COSH(RAZ)+ (E2/RA2)*SINH(RAZ)+ (E3/RB2)*COSH(RBZ)+
002272      1 (E4/RB2)*SINH(RBZ)+ (E5/6.)*Z*Z*Z+ (E6/2.)*Z*Z+ E7*Z+ E8
002312      WRITE (6,37) ZET,ZETP1,ZETP2, ETA,ETAP1,ETAP2
002312 37 FORMAT (1X, 6E13.4)
002314      IF (Z.EQ.HT) GO TO 26
002316      Z= Z+H
002316      GO TO 20
002316 26 STOP
002320      END

```

UNUSED COMPILER SPACE
010700

APPENDIX-C
EXPERIMENTAL DATA

TABLE C1

STRAIN DATA FOR MODEL WITH FLOORS

strain in μ in/in for load at 8th floor

Strain Gauge No.	Loading				
	2.5 lbs.	5 Lbs.	7.5 lbs.	10 lbs.	12.5 lbs.
1	-30	-56	-92	-116	-140
2	+24	+48	+68	+86	+106
3	-28	-56	-84	-110	-130
4	+14	+38	+64	+94	+118
5	-30	-56	-84	-114	-138
6	+18	+50	+80	+112	+142
7	+15	+30	+48	+60	+76
8	0	-10	-22	-34	-46
9	*				
10	-22	-48	-74	-100	-126
11	+5	+5	+7	+8	+9
12	-12	-25	-40	-55	-68
13	-20	-44	-72	-96	-120
14	+14	+36	+48	+60	+70
15	-14	-36	-52	-64	-78
16	-3	-7	-11	-15	-20
17	-14	-28	-38	-42	-54
18	+6	+16	+32	+48	+60
19	-8	-24	-42	-56	-70
20	+25	+58	+98	+134	+170
21	-28	-58	-90	-118	-148
22	+16	+38	+60	+84	+108
23	+6	+12	+20	+28	+34
24	+14	+36	+56	+78	+98

* Strain Gauge out of order

TABLE C1 (continued)

STRAIN DATA FOR MODEL WITH FLOOR

Strain in μ in/in for load at 8th floor

Strain Gauge No.	Loading				
	10 lbs.	7.5 lbs.	5 lbs.	2.5 lbs.	0
1	-114	-84	-50	-25	-2
2	+90	+70	+50	+28	+2
3	-114	-94	-74	-52	-26
4	+94	+70	+40	+14	0
5	-114	-90	-62	-34	-6
6	+116	+84	+62	+28	+8
7	+62	+48	+36	+20	+2
8	-38	-24	-10	-2	0
10	-104	-80	-54	-26	0
11	+8	+8	+6	+4	0
12	+56	+44	+30	+16	+2
13	-100	-74	-50	-26	-2
14	+60	+52	+42	+26	+10
15	-64	-54	-40	-26	-12
16	-15	-12	-8	-5	-2
17	-44	-36	-23	-13	-2
18	+50	+46	+24	+12	0
19	-60	-46	-32	-22	-11
20	+138	+106	+72	+36	+6
21	-124	-94	-66	-38	-10
22	+86	+62	+40	+14	0
23	+28	+18	+12	+4	+2
24	+80	+58	+40	+16	+3

* Strain Gauge out of order

TABLE C2

DEFLECTION DATA FOR MODEL WITH FLOOR

Deflection in inch for load at 8th floor

Dial Gauge No.	Loading				
	2.5 lbs.	5 lbs.	7.5lbs	10 lbs.	12.5 lbs.
1	.011	.032	.055	.079	.097
2	.0285	.0725	.118	.1695	.2125
3	.009	.026	.043	.062	.084
4	.025	.066	.108	.151	.187
5	.007	.022	.041	.061	.083
6	.020	.053	.086	.122	.152
7	.003	.010	.022	.024	.030
8	.010	.030	.048	.069	.085
9	.002	.004	.006	.009	.011
10	.004	.011	.0175	.025	.031

Dial Gauge No.					
	10 lbs.	7.5 lbs.	5 lbs.	2.5 lbs.	0
1	.082	.063	.040	.016	.005
2	.1715	.1325	.0885	.040	.010
3	.066	.049	.032	.014	.005
4	.153	.120	.080	.036	.006
5	.064	.047	.028	.011	.003
6	.124	.097	.064	.028	.003
7	.025	.020	.013	.005	.001
8	.070	.054	.036	.015	.003
9	.008	.007	.005	.002	.001
10	.025	.020	.013	.090	.002

TABLE C3

STRAIN DATA FOR MODEL WITH FLOORS

Strain in μ in/in for load at 6th floor

Strain Gauge No.	2.5 lbs.	5 lbs.	7.5 lbs.	10 lbs.	12.5 lbs.
1	-24	-44	-62	-84	-98
2	+24	+44	+66	+86	+106
3	-24	-46	-70	-92	-116
4	+4	+18	+36	+56	+78
5	-18	-40	-60	-80	-100
6	+4	+18	+40	60	+80
7	+8	+20	+30	+42	+52
8	0	-4	-12	-24	-36
9	*				
10	-----no reading-----				
11	-4	-8	-12	-17	-22
12	+6	+10	+16	+26	+30
13	-4	-12	-22	-34	-44
14	+10	+18	+28	+36	+42
15	+2	+4	+5	+6	+8
16	-10	-22	-36	-50	-66
17	+5	+10	+18	+26	+34
18	-6	-16	-24	-30	-38
19	0	-2	-12	-24	-42
20	+14	+32	+50	+72	+92
21	-14	-32	-48	-68	-86
22	+5	+14	+24	+34	+44
23	+12	+32	+46	+62	+78
24	-2	-4	-5	-6	-6

* Strain Gauge out of order

TABLE C3 (continued)

STRAIN DATA FOR MODEL WITH FLOORS

Strain in μ in/in for load at 6th floor

Strain Gauge No.	Loading				
	10 lbs.	7.5 lbs.	5 lbs.	2.5 lbs.	0
1	-88	-72	-56	-22	-6
2	+88	+68	+46	+26	0
3	-96	-74	-52	-28	-6
4	+60	+38	+22	+10	+2
5	-82	-64	-44	-26	-6
6	-58	-44	-16	-8	0
7	+42	+34	+24	+12	0
8	-28	-18	-12	-2	0
9	*				
10	-----no reading-----				
11	-18	-13	-9	-4	-2
12	+26	+20	+14	+8	+4
13	-36	-24	-18	-6	-2
14	+34	+28	+20	+10	0
15	+6	+4	+2	0	0
16	-58	-42	+32	-14	+4
17	+30	+20	+14	+6	0
18	-30	-25	-18	-10	-2
19	-42	-22	-12	-2	0
20	+76	+54	+38	+14	+2
21	-72	-54	-36	-18	-4
22	+34	+24	+15	+6	0
23	+62	+48	+32	+14	0
24	-4	-4	-4	-2	0

* Strain Gauge out of order

TABLE C4

DEFLECTION DATA FOR MODEL WITH FLOORS

Deflection in inch for load at 6th floor

Dial Gauge No.	Loading				
	2.5 lbs.	5 lbs.	7.5 lbs.	10lbs.	12.5 lbs.
1	.005	.024	.030	.043	.058
2	.015	.0525	.072	.100	.129
3	.004	.020	.0255	.036	.047
4	.014	.052	.060	.097	.1245
5	.004	.018	.023	.033	.046
6	.012	.043	.059	.081	.105
7	.002	.010	.012	.017	.022
8	.007	.026	.036	.050	.064
9	.001	.004	.005	.006	.008
10	.003	.011	.014	.019	.025

Dial Gauge No.					
	10 lbs.	7.5 lbs.	5 lbs.	2.5 lbs.	0
1	.044	.032	.019	.008	0
2	.105	.076	.047	.020	.001
3	.038	.027	.016	.006	0
4	.101	.074	.046	.019	.001
5	.036	.024	.014	.005	.001
6	.085	.062	.038	.016	.001
7	.018	.014	.008	.003	0
8	.052	.038	.023	.009	0
9	.007	.005	.003	.001	0
10	.020	.015	.009	.004	0

TABLE C5

STRAIN DATA FOR MODEL WITH FLOORS

Strain in μ in/in for load at 4th floor

Strain Gauge No.	Loading				
	2.5 lbs.	5 lbs.	7.5 lbs.	10 lbs.	12.5 lbs.
1	-21	-36	-54	-68	-84
2	+18	+40	+60	+82	+100
3	-14	-34	-50	-68	-84
4	0	+6	+18	+34	+50
5	-12	-26	-28	-50	-62
6	0	+6	+14	+26	+40
7	+8	+20	+30	+42	+52
8	-----no reading-----				
9	*				
10	-7	-18	-30	-40	-52
11	-4	-8	-12	-17	-22
12	+6	+10	+16	+22	+26
13	-4	-12	-22	-34	-44
14	+10	+18	+28	+36	+42
15	-2	-4	-5	-6	-8
16	-10	-22	-36	-50	-66
17	+5	+10	+18	+26	+34
18	-6	-16	-24	-30	-38
19	0	-2	-12	-24	-42
20	+14	+32	+50	+72	+92
21	-14	-32	-48	-68	-86
22	+5	+14	+24	+34	+44
23	+12	+32	+46	+62	+78
24	-2	-4	-5	-6	-6

* Strain Gauge out of order

TABLE C5 (continued)

STRAIN DATA FOR MODEL WITH FLOORS
 Strain in μ in/in for load at 4th floor

Strain Gauge No.	10 lbs.	7.5 lbs.	5 lbs.	2.5 lbs.	0
1	-72	-56	-44	-28	-10
2	+78	+58	+34	+14	+2
3	-70	-54	-36	-18	-4
4	+40	+34	+24	+10	+6
5	-52	-40	-30	-16	-6
6	+32	+20	+12	+4	0
7	+42	+34	+24	+12	0
8	-----no reading-----				
9	*				
10	-43	-32	-22	-10	-4
11	-18	-13	-10	-5	0
12	+22	+16	+10	+4	0
13	-32	-20	-14	-4	0
14	+34	+28	+20	+10	0
15	-6	-4	-2	-2	0
16	-58	-42	-32	-14	-4
17	+30	+20	+14	+8	0
18	-30	-22	-17	-10	-2
19	-42	-22	-12	-2	0
20	+76	+54	+38	+14	+2
21	-72	-54	-36	-18	-4
22	+34	+25	+15	+6	0
23	+62	+48	+32	+14	0
24	-4	-4	-4	-2	0

TABLE C6

DEFLECTION DATA FOR MODEL WITH FLOOR

Deflection in inch for load at 4th floor

Dial Gauge No.	Loading				
	2.5 lbs.	5 lbs.	7.5 lbs.	10 lb.s	12.5 lbs.
1	.001	.008	.012	.020	.029
2	.008	.025	.037	.054	.073
3	.001	.008	.012	.018	.026
4	.007	.024	.037	.055	.074
5	.001	.007	.011	.017	.025
6	.006	.020	.032	.047	.064
7	.001	.005	.008	.011	.015
8	.005	.015	.023	.034	.046
9	0	.002	.003	.005	.006
10	.002	.007	.010	.015	.020

Dial Gauge No.	Loading				
	10 lbs.	7.5 lbs.	5 lbs.	2.5 lbs.	0
1	.023	.017	.012	.007	0
2	.060	.046	.031	.020	.003
3	.021	.016	.011	.007	0
4	.060	.048	.031	.019	.003
5	.019	.014	.010	.006	.001
6	.052	.040	.026	.016	.003
7	.012	.0095	.007	.004	.001
8	.037	.029	.019	.011	.001
9	.005	.004	.003	.0015	0
10	.016	.012	.008	.006	.001

TABLE C7

STRAIN DATA FOR MODEL WITH BEAMS

Strain in μ in/in for lateral load at top

Strain Gauge No.	5 lbs.	10 lbs.	15 lbs.	20 lbs.	25 lbs.
1	-20	-60	-82	-120	-152
2	-18	-42	-72	-104	-132
3	-16	-40	-76	-112	-140
4	no strain				
5	-12	-24	-46	-52	-64
6	-22	-46	-74	-98	-124
11	+8	+24	+38	+52	+64
12	+22	+42	+62	+82	+110
13	+35	+63	+83	+103	+123
14	+16	+38	+62	+88	+114
15	+14	+24	+34	+44	+54
16	no strain				
17	-14	-29	-44	-59	-73
18	+18	+42	+76	+114	+160
19	-16	-33	+50	-66	-84
20	+10	+30	+40	+54	+70
21	-54	-138	-230	-326	+410

TABLE C7 (continued)

STRAIN DATA FOR MODEL WITH BEAMS

strain in μ in/in for lateral load at top

Strain Gauge No.	Loading				
	20 lbs.	15 lbs.	10 lbs.	5 lbs.	0
1	-124	-92	-64	-34	-6
2	-106	-76	-46	-22	0
3	-120	-86	-56	-26	-4
4	no strain				
5	-54	-42	-26	-14	-2
6	-104	-80	-56	-30	-5
11	+56	+44	+30	+14	0
12	+90	+72	+52	+28	+2
13	+99	+79	+65	+41	+6
14	+94	+76	+50	+26	+4
15	+46	+36	+26	+16	+2
16	no strain				
17	-62	-46	-30	-13	-2
18	+130	+98	+68	+36	+10
19	-70	-52	-34	-14	0
20	-60	-50	-34	-20	-2
21	-336	-252	-182	-76	-8

TABLE C8

DEFLECTION DATA FOR MODEL WITH FLOORS

Deflection in inch for lateral load at top

Dial Gauge No.	Loading				
	5 lbs.	10 lbs.	15 lbs.	20 lbs.	25 lbs.
1	.015	.030	.052	.064	.079
2	.013	.024	.046	.056	.061
3	.013	.020	.037	.042	.053
4	.007	.012	.019	.026	.030
5	.0036	.0069	.0095	.0119	.0127
6	.024	.060	.100	.132	.167
7	.019	.050	.086	.113	.144
8	.016	.042	.071	.100	.121
9	.009	.026	.045	.060	.077
10	.0051	.0115	.0206	.0270	.0336

Dial Gauge No.	Loading				
	20 lbs.	15 lbs.	10 lbs.	5 lbs.	0
1	.066	.054	.035	.020	.004
2	.057	.046	.029	.016	.006
3	.045	.039	.033	.014	.004
4	.029	.021	.015	.010	.003
5	.0124	.0099	.0073	.0048	.0013
6	.147	.115	.085	.046	.011
7	.127	.099	.074	.038	.009
8	.107	.083	.063	.032	.007
9	.068	.054	.039	.019	.006
10	.0289	.0230	.0178	.0121	.0043

TABLE C9

STRAIN DATA FOR MODEL WITH BEAMS

Strain in μ in/in for torque at top

Strain Gauge No.	Loading				
	25 lb-in	50 lb-in	75 lb-in	100 lb-in	125lb-in
1	-72	-164	-248	-356	-446
2	-30	-66	-116	-158	-200
3	no strain				
4	-34	-76	-124	-168	-200
5	-12	-24	-38	-50	-60
6	+6	+22	+36	+52	+60
11	-10	-28	-48	-70	-88
12	+15	+38	+58	+79	+101
13	+50	+113	+173	+233	+301
14	-11	-24	-36	-49	-62
15	+10	+23	+35	+52	+65
16	+25	+55	+83	+114	+142
17	+30	+70	+108	+154	+184
18	+26	+54	+89	+117	+145
19	+10	+23	+36	+48	+64
20	-14	-22	-34	-47	-62
21	-20	-32	-42	-56	-70

TABLE C9 (continued)

STRAIN DATA FOR MODEL WITH BEAMS

Strain in μ in/in for torque at top

Strain Gauge No.	Loading				
	100lb-in	75lb-in	50lb-in	25lb-in	0
1	-380	-308	-212	-92	-15
2	-172	-134	-96	-56	-6
3	no strain				
4	-174	-144	-108	-74	-4
5	-50	-38	-28	-14	0
6	+54	+38	+25	+16	+2
11	-76	-64	-46	-28	-8
12	+82	+62	+42	+22	+2
13	+236	+176	+118	+60	+4
14	-51	-39	-27	-15	-2
15	-50	-38	-26	-14	-3
16	+116	+89	+60	+30	+2
17	+160	+128	+100	+62	+6
18	+119	+92	+56	+30	+3
19	+52	+39	+25	+14	+2
20	-54	-38	-26	-10	0
21	-50	-43	-33	-22	-2

TABLE C10

DEFLECTION DATA FOR MODEL WITH FLOORS

Deflection in inch for torque at top

Dial Gauge No.	Loading				
	25lb-in	50lb-in	75lb-in	100lb-in	125lb-in
1	.004	.008	.018	.023	.031
2	.003	.007	.016	.022	.028
3	.003	.0005	.013	.018	.025
4	.002	.004	.008	.011	.015
5	.0006	.0018	.0041	.0054	.0071
6	.017	.059	.101	.139	.180
7	.013	.043	.081	.110	.146
8	.009	.035	.066	.091	.121
9	.006	.023	.043	.059	.078
10	.0023	.0099	.019	.0262	.0354

Dial Gauge No.	Loading				
	100lbs-in	75lb-in	50lb-in	25lb-in	0
1	.026	.018	.011	.008	.003
2	.023	.016	.010	.007	.004
3	.021	.014	.008	.005	.002
4	.013	.009	.005	.003	.001
5	.006	.0042	.0024	.0018	.0009
6	.149	.111	.072	.045	.008
7	.119	.094	.056	.035	.007
8	.099	.073	.046	.029	.004
9	.064	.047	.030	.018	.003
10	.0284	.0208	.0125	.0071	.001

BIBLIOGRAPHY

1. AFSAR, M.
"Lateral Loading of Shear Wall Models", M. Eng. Thesis, August 1967.
2. QURESHI, A.A.
"Behaviour of Shear Wall Models with Circular Wall Openings", M. Eng. Thesis, February 1968.
3. SPEIRS, J.W..
"Lateral Loading of Small Scale Shear Wall Building with Floor Slab", M. Eng. Thesis, March 1969.
4. RAINA, R.K.
"Dynamic Response of Thin Wall Shear Wall Model", M. Eng. Thesis, January 1970.
5. SWIFT, R.D.
"Analysis of Asymmetrical Coupled Shear Wall Structures", M. Eng. Thesis, April 1970.
6. COULL, A. and SMITH, S.B.
"Analysis of Shear Wall Structures", A review of previous research, Symposium of Tall Buildings, pp. 138-58. April 1966.
7. WINOKUR, A. and GLUCK, J.
"Lateral Loads in Asymmetric Multi-storey Structures", Journal of the Structural Division, Proc. of the Am. Soc. of Civil Eng., V. 94, n. ST. 3, March 1968.
8. VLASOV, V.Z.
"Thin Walled Elastic Beams", Published for the National Science Foundation, Washington D.C. by the Israel Program for Scientific Translations, Jerusalem 1961 (First edition published in original Russian 1940).
9. ZBIROHOWSKI-KOSCIA
"Thin Walled Beams--from Theory to Practice", Crosby-Lockwood, London 1969.

10. QADEER, A.
"The Interaction of Floor Slabs and Shear Walls", Ph.D. Thesis at University of Southampton, October 1968.
11. QADEER, A. and SMITH, S.B.
"The Bending Stiffness of Slabs Connecting Shear Walls", ACI Journal, June 1969.
12. TARANATH, B.S.
"The Torsional Behaviour of Open Section Shear Wall Structures", Ph.D. Thesis at University of Southampton, October 1968.
13. BECK, H.
"Contribution to the Analysis of Shear Walls", Journal of the American Concrete Institute, V.59, pp 1055-69 1962.
14. ROSMAN, R.
"Approximate Analysis of Shear Wall Subjected to Lateral Loads". Journal of the American Concrete Institute, V.61, pp 717-32, 1964.
15. COULL, A. and CHOUDHURY, J.R.
"Stresses and Deflections in Coupled Shear Walls", Journal of the American Concrete Insitute, pp. 65-72, February 1967.
16. COULL, A. and CHOUDHURY, J.R.
"Analysis of Coupled Shear Walls", Journal of the American Concrete Institute, pp. 587-93, September 1967.
17. CHOUDHURY, J.R.
"Analysis of Plane Spatial System of Interconnected Shear Walls", Ph.D. Thesis. University of Southampton, October 1968.
18. MICHAEL, D.
"Torsional Coupling of Core Walls in Tall Buildings", The Structural Engineer, N.2, V 47, February 1969.

19. JENKINS, W.M. and HARRISON, T.
" Analysis of Tall Buildings with Shear Wall Under Bending and Torsion", Tall Buildings, pp. 413-49.
20. HOLMES, M. And ASTILL, A.W.
"Experimental Stresses and Deflections of a Model Shear Wall Structure", Journal of the American Concrete Institute, pp. 667-76, August 1969.
21. ROSMAN, R.
"Analysis of Pierced Torsion Boxes", Acta Technica Scientiarum Hungaricae, Tomus 65 (3-4), pp.365-97, 1969.
22. GLUCK, J.
"Lateral Load Analysis of Asymmetric Multistorey Structures", Journal of the Structural Division, Proc. of the Am. Soc. of Civil Eng., ST2, pp. 317-33, February 1970.
23. MACLEOD, I.A.
"Connected Shear Wall of Unequal Width", Journal of American Concrete Institute, pp. 408-12, May 1970.
24. BISWAS, J.K. and TSO, W.K.
Discussion of Ref. (22), Journal of Structural Division Proceedings Am. Soc. of Civil Engineers, ST 11, pp. 2540-44, November 1970.
25. HOLTZER, H.
Die Berechnung der Drehshwingungen, Springer-Verlag, Berlin 1921.
26. MYKLESTAD, N.O.
"New Methods of Calculating Natural Modes of Coupled Bending Torsion Vibration of Beam", Tran. ASME, Vol. 67 pp. 61-67, 1945.
27. THOMSON, W.T.
"Matrix Solution of Vibration of Nonuniform Beam", Jour. Appl. Mech. Vol. 17, No. 3, pp.337-339, 1950.
28. COULL, A. and IRWIN, A.W.
"Analysis of Load Distribution in Multistorey Shear Wall Structures", The Structural Engineer, N.8, V.48, August 1970.

29. TSO, W.K.

"Coupled Vibration of Thin-Walled Elastic Bars",
J. Engineering Mechanics Division, A.S.C.E. Vol. 91
EM3 1965 pp.33-52.